

## Dependence of the Jahn-Teller distortion in $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$ on the isovalent Mn-site substitution

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**Abstract.** We investigated the relative importance of removing the  $\text{Mn}^{3+}$  Jahn-Teller distortion in driving ferromagnetism in  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  combining x-ray powder diffraction and x-ray absorption spectroscopy at the Mn and Sc K-edges. By increasing the Sc content, the orthorhombic distortion of the *Pbnm* cell in  $\text{LaMnO}_3$  decreases but the unit-cell remains slightly distorted in  $\text{LaScO}_3$ . Besides, the nearly tetragonal-distorted  $\text{MO}_6$  in  $\text{LaMnO}_3$  continuously evolves into a nearly regular one in  $\text{LaScO}_3$ . On the other hand, x-ray absorption spectra show that the  $\text{MnO}_6$  octahedron remains Jahn-Teller distorted and the  $\text{ScO}_6$  octahedron is nearly regular along the whole series. Moreover, the ordering of the  $\text{Mn}^{3+}$  Jahn-Teller distortion is not disrupted in the *ab* plane for any Sc concentration. This contrasts with the Ga-substituted compounds, where a regular  $\text{MnO}_6$  is found for  $x > 0.5$ . However, both  $\text{LaMn}_{0.5}\text{Sc}_{0.5}\text{O}_3$  and  $\text{LaMn}_{0.5}\text{Ga}_{0.5}\text{O}_3$  show ferromagnetic behavior independently of the presence (or not) of Jahn-Teller distorted  $\text{Mn}^{3+}$ . Thus, our results point to the Mn-sublattice dilution as the main effect in driving ferromagnetism in these manganites over local structure effects previously proposed by the spin flipping or the vibronic superexchange models.

### 1. Introduction

Doped manganites have been extensively studied owing to their colossal magnetoresistance property and its potential applications to magnetic devices. In particular,  $\text{LaMnO}_3$  is considered as the prototype of the cooperative Jahn-Teller (JT) systems and orbital-ordered state [1]. It also develops long-range antiferromagnetic (AFM) ordering of type A below  $T_N = 140$  K. The substitution at the Mn dramatically modifies the magnetic and electronic properties. It is known that the A-type AFM orbitally ordered state of  $\text{LaMnO}_3$  is easily destroyed to give way to ferromagnetism (FM) by introducing various cations on the Mn site [2-5]. However, this ferromagnetic behavior derives from a complex interplay between several mechanisms: double-exchange and super-exchange interactions, orbital order and JT distortions. The substitution of Mn with non-magnetic isovalent ions such as  $\text{Ga}^{3+}$  ( $3d^{10}$ ) or  $\text{Sc}^{3+}$  ( $3d^0$ ) weakens the  $\text{Mn}^{3+}$ - $\text{Mn}^{3+}$  super-exchange interactions but surprisingly favors the appearance of a FM ground state for  $\text{LaMn}_{0.5}\text{Ga}_{0.5}\text{O}_3$  [3,4] and  $\text{LaMn}_{0.5}\text{Sc}_{0.5}\text{O}_3$  [5]. Two models have been reported to account for this property: the spin flipping of Mn  $e_g$ -orbital in JT-distorted  $\text{Mn}^{3+}$  atoms surrounded by



Ga (Sc) atoms [6] and the FM vibronic superexchange interaction coming from disordered and fluctuating JT-distorted  $\text{Mn}^{3+}$  atoms [7]. Both models assume that the change from A-type AFM to FM ordering is originated by the orbital direction adopted by the  $\text{Mn}^{3+}$  ion and it is similar for Ga and Sc non-magnetic substitutions, so the atomic  $e_g$  orbital would be perturbed in the same way for the two series.

These models contrast with previous x-ray absorption spectroscopy (XAS) studies [8-10]. A XAS study of the  $\text{LaMn}_{1-x}\text{Ga}_x\text{O}_3$  series have revealed that the tetragonal distortion of  $\text{MnO}_6$  octahedron continuously decreases with the gallium content in such a way that for  $x > 0.5$   $\text{MnO}_6$  distortion disappears and becomes regular [8,9], simultaneously to the appearance of long-range FM. Moreover, a XAS study of the isostructural  $\text{TbMn}_{1-x}\text{Sc}_x\text{O}_3$  series showed that the  $\text{MnO}_6$  octahedra remain tetragonally distorted along the whole dilution range, independently of the Sc content [10].

Here, we have studied the crystallographic structure by x-ray powder diffraction (XRD) and the local structure by XAS at the Mn and Sc K-edges in the  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  series in order to clarify the importance of removing the Jahn-Teller distortion in driving the FM behavior.

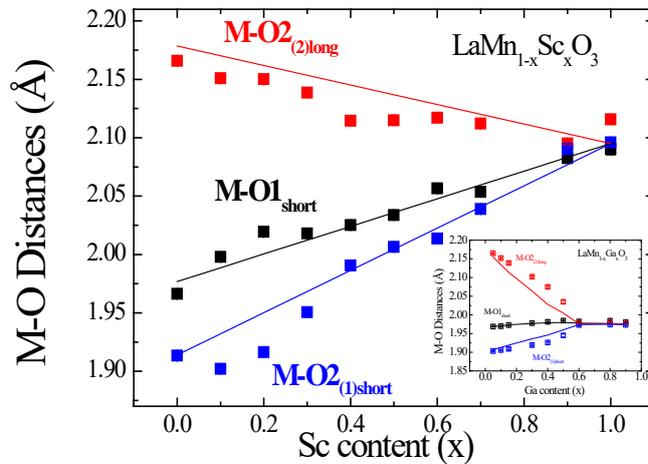
## 2. Experimental details

Powder  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  ( $x = 0, 0.1, 0.2, 0.3, 0.4, 0.5, 0.6, 0.7, 0.9, 1$ ) samples were synthesized following the same ceramic procedure as previously outlined for  $\text{LaMn}_{1-x}\text{Ga}_x\text{O}_3$  [3]. XRD patterns were collected for the whole series in the 2-theta range between  $18^\circ$  and  $135^\circ$  at room temperature by using a D-max Rigaku system with a rotating anode and selecting the Cu  $K\alpha$  radiation. The crystal structures were refined by the Rietveld method using the program FULLPROF [11]

The Mn K-edge XAS spectra were recorded at the BM23 beamline of the European Synchrotron Radiation Facility. Data were recorded at a fixed  $T \sim 80$  K in transmission mode for Sc concentration up to  $x = 0.6$  while for higher dilutions fluorescence detection was used. The Sc K-edge XAS measurements were carried out at fixed temperature at the CLAESS beamline of the ALBA synchrotron for Sc concentration  $x \geq 0.3$  whereas temperature-dependent XAS spectra of the  $\text{LaMn}_{0.4}\text{Sc}_{0.6}\text{O}_3$  sample were collected at the XAFS beamline of the Elettra Synchrotron. Data for  $x > 0.4$  were recorded in transmission mode while for lower Sc concentration fluorescence detection was used. XANES spectra were normalized to unity edge jump using the Athena software from the Demeter package [12]. The extraction of the  $\chi(k)$  signals were also performed using Athena program, and R-space EXAFS spectra were obtained by calculating the Fourier Transform (FT) of the  $k^2\chi(k)$  signals in the  $(2.5-14) \text{ \AA}^{-1}$  k-range. The EXAFS structural analysis was performed using theoretical phases and amplitudes calculated by the FEFF-6 code [13] and fits to the experimental data were carried out in R-space with the Artemis program of the Demeter package [12].

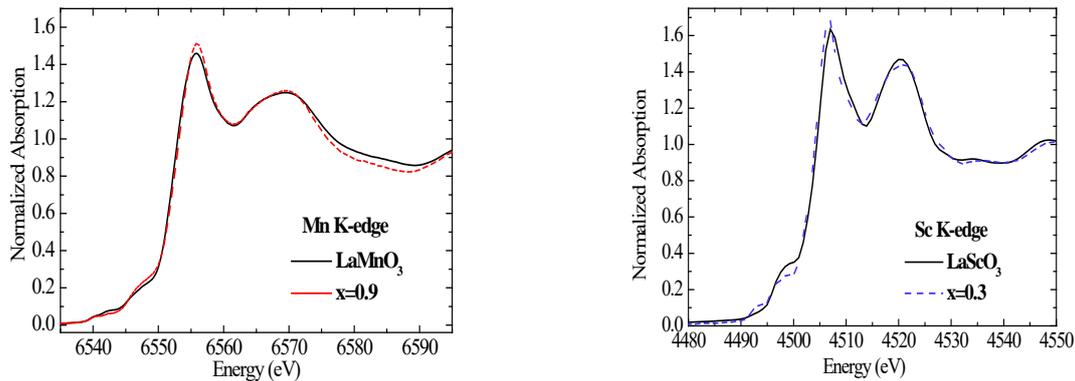
## 3. Results and conclusions

The XRD results show that the tetragonal distortion of the  $\text{MO}_6$  octahedron continuously decreases with the Sc content (figure 1). There are two possibilities that could explain this fact: either both  $\text{MnO}_6$  and  $\text{ScO}_6$  octahedron are equally distorted, and the magnitude of the correlated distortion decreases with  $x$ , or the observed distortion of  $\text{MO}_6$  by XRD corresponds to the average of different  $\text{MnO}_6$  and  $\text{ScO}_6$  octahedra. In figure 1, we have compared the M-O distances from XRD with the weighted addition of the Mn-O and Sc-O distances taken from XRD for the end-member compounds of the series,  $\text{LaMnO}_3$  and  $\text{LaScO}_3$ . The same comparison was also made with the weighted addition of the Mn-O and Ga-O distances for the  $\text{LaMn}_{1-x}\text{Ga}_x\text{O}_3$  series as reported in the inset of figure 1. The agreement is very good, which means that the geometrical orientation of the  $\text{MnO}_6$  tetragonal distortion is preserved in the  $ab$  plane in both, Sc-substituted and Ga-substituted compounds. In the Ga-substituted samples, different  $\text{MnO}_6$  and  $\text{GaO}_6$  octahedra are deduced by XAS where the JT-distortion of the  $\text{Mn}^{3+}\text{O}_6$  decreases with the Ga substitution and the appearing of FM concurs with the presence of regular  $\text{Mn}^{3+}\text{O}_6$  octahedron ( $x > 0.5$ ). The occurrence of Mn tetragonal distortions, which remain ordered even up to high Sc content, has been also reported for the  $\text{TbMn}_{1-x}\text{Sc}_x\text{O}_3$  series [10].



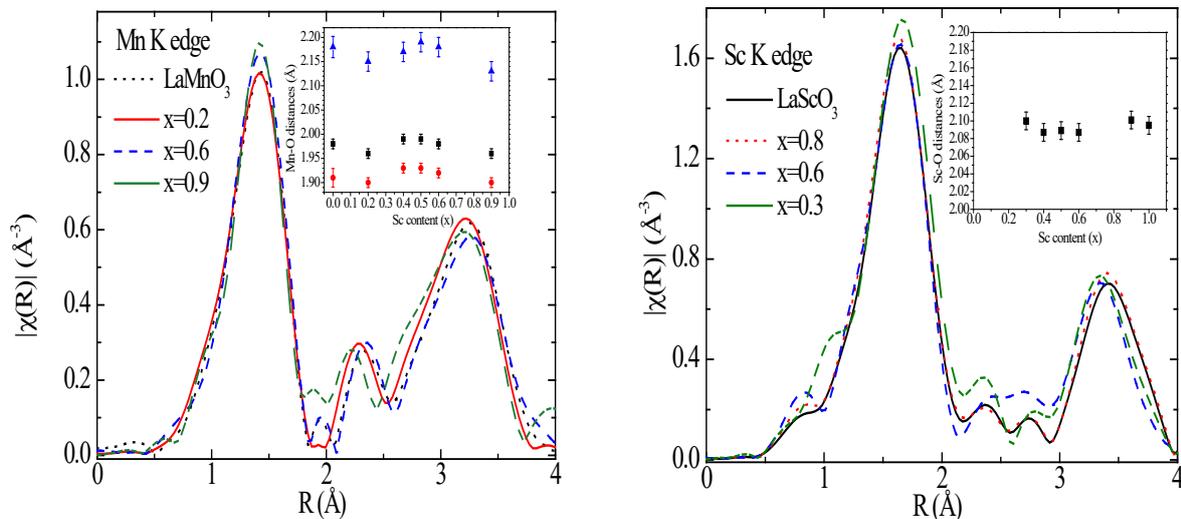
**Figure 1.** Evolution with the Sc content of the  $\text{MO}_6$  octahedron (closed symbols) compared to the weighted addition of the M-O distances of the two end-member  $\text{LaMnO}_3$  and  $\text{LaScO}_3$  (lines) from XRD. **Inset:** same comparison for the  $\text{LaMn}_{1-x}\text{Ga}_x\text{O}_3$  series.

No edge energy shift is observed between the  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  samples and either the  $\text{LaMnO}_3$  sample at the Mn K-edge or the  $\text{LaScO}_3$  sample at the Sc K-edge (figure 2) within the precision of the experiments (0.4 eV). Thus,  $\text{Mn}^{3+}$  formal valence keeps constant along the dilution.



**Figure 2. Left.** Normalized Mn K-edge XANES spectra of  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  ( $x=0, 0.9$ ). **Right.** Normalized Sc K-edge XANES spectra of  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  ( $x=0.3, 1$ ).

Figure 3 shows the modulus of the FT of the  $k^2$ -weighted EXAFS signals of selected concentrations of Sc. All samples show a main peak at  $\sim 1.5 \text{ \AA}$  (Mn K-edge) and  $\sim 1.6 \text{ \AA}$  (Sc K-edge) corresponding to the first (Mn-O or Sc-O) coordination shell without phase shift correction. The intensity and shape of this peak is practically independent of the dilution indicating that the oxygen coordination geometry around both, Mn and Sc atoms, remains almost unaltered along the whole series. The structural analysis at the Mn K-edge was performed considering the same shift on the equatorial plane ( $\text{Mn-O}_{1\text{short}}$  and  $\text{Mn-O}_{2\text{short}}$ ,  $\Delta r_1$ ) and a different one for the apical distances ( $\text{Mn-O}_{2\text{long}}$ ,  $\Delta r_2$ ) of the  $\text{MnO}_6$  octahedron. The relevant structural parameters for the first oxygen coordination shell of all samples are shown in the inset of figure 3(left). The same tetragonal distorted (JT)  $\text{MnO}_6$  as in  $\text{LaMnO}_3$  is found for all Sc substitutions. At the Sc K-edge, we have performed the fits considering the three pairs of Sc-O distances slightly distorted of  $\text{LaScO}_3$  but they converge to a unique distance. As shown in the inset of figure 3 (right), the same nearly regular  $\text{ScO}_6$  octahedron as in  $\text{LaScO}_3$  is found for all Sc concentrations. The temperature-dependent EXAFS spectra of  $\text{LaMn}_{0.4}\text{Sc}_{0.6}\text{O}_3$  at the Mn and Sc K-edges between 80 and 300 K confirm the robustness of  $\text{MnO}_6$  tetragonal distortion and the lack of significant static distortion of  $\text{ScO}_6$  for  $\text{LaMn}_{0.4}\text{Sc}_{0.6}\text{O}_3$  along this temperature range.



**Figure 3. Left.** Modulus of the FT at the Mn K-edge for selected  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  samples. **Inset:** Evolution of the fitted Mn-O distances with the Sc content. **Right.** Modulus of the FT at the Sc K-edge for selected  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$  samples. **Inset:** Evolution of the fitted Sc-O distances with the Sc content.

The present results show a marked difference in the local structure for  $\text{LaMn}_{1-x}\text{Ga}_x\text{O}_3$  and  $\text{LaMn}_{1-x}\text{Sc}_x\text{O}_3$ .  $\text{MnO}_6$  is hardly distorted in an environment composed by regular  $\text{GaO}_6$  octahedra whereas this distortion results more feasible when  $\text{MnO}_6$  is surrounded by  $\text{ScO}_6$  octahedra because of the larger tilts of the  $\text{ScO}_6$  octahedron in the orthorhombic  $Pbnm$  cell. Therefore, the occurrence of the tetragonal distortion of the  $\text{MnO}_6$  octahedron strongly depends on the lattice in which Mn atom is allocated. Despite the different structural behaviour found for Ga and Sc substitutions, the magnetic response for  $x \geq 0.5$  is similar in both cases. The two proposed models [6,7] to account for the change from A-type AFM to FM cannot be supported by either the present structural study or that in the Ga-substituted samples [8,9] so the Mn-sublattice dilution effect may be the main responsible for the magnetic behaviour of these manganites.

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