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# Tuning the out-of-plane magnetic textures of electrodeposited $Ni_{so}Fe_{10}$ thin films

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# Tuning the out-of-plane magnetic textures of electrodeposited Ni<sub>90</sub>Fe<sub>10</sub> thin films

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# ABSTRACT

This study investigates the out-of-plane magnetization component of electrodeposited  $Ni_{90}Fe_{10}$  thin films grown under different applied magnetic field conditions. The formation of stripe domains is gradual, as there is a thickness range in which the transcritical shape appears in the hysteresis loops, while only magnetic ripples are measured in the magnetic force microscopy images. For instance, samples deposited under the residual magnetic field generated by the switched-off magnetic stirrer exhibit the transcritical shape in the in-plane hysteresis loops at a thickness of 400 nm, even though corresponding magnetic force microscopy images do not reveal the presence of stripe domains. When a perpendicular magnetic field of 100 Oe is applied during growth, stripe domains become visible in microscopy images, along with the transcritical shape in the hysteresis loop at 400 nm. This implies that the critical thickness for stripe formation can be reduced by applying a perpendicular magnetic field during glowth, the smaller the critical thickness. These results underscore the importance of controlling the external magnetic field during electrodeposition for more effective tuning of the magnetic textures in electrodeposited  $Ni_{90}Fe_{10}$  films.

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# I. INTRODUCTION

Spin textures are gaining increasing interest in the field of spintronics due to their reconfigurability and stability.<sup>1–3</sup> The periodic modulation of a ferromagnet can be achieved through artificial arrangements of dots and stripes, or ion implantation.<sup>4</sup> In materials with moderate perpendicular magnetic anisotropy (PMA), the periodic magnetic configuration known as stripe domains appear due to competing interactions, arising from the alternating up and down out-of-plane (OOP) orientation of the magnetization.<sup>5,6</sup> These stripe domains can be utilized, for example, to generate spin waves.<sup>7–10</sup> In ferromagnets with moderate or low PMA, stripe domains appear above a critical thickness ( $t_{cr}$ ),

$$t_{cr} = 2\pi \sqrt{A_{ex}/K_{OOP}},\tag{1}$$

with  $A_{ex}$  being the exchange energy per unit length and  $K_{OOP}$  the OOP magnetic anisotropy energy constant. When PMA is not sufficiently large to induce the formation of stripe domains, a magnetic contrast known as ripple, characterized by corrugation, becomes observable.<sup>11–15</sup> The origin of magnetic ripple lies in the inhomogeneous magnetization caused by magnetic anisotropy fluctuations. Therefore, reducing these fluctuations through the enhancement of the PMA will promote the formation of stripe domains.

Ni–Fe alloys have been extensively studied due to their low coercivity and large saturation magnetization  $(M_{sat})$ .<sup>16</sup> Ni<sub>90</sub>Fe<sub>10</sub> is particularly interesting because of its magnetoelastic properties,<sup>17–19</sup> exhibiting a magnetostriction constant ( $\lambda$ ) of approximately –20 ppm. Evaporated Ni<sub>90</sub>Fe<sub>10</sub> thin layers deposited on Cu can show a significant OOP magnetic anisotropy contribution due to mechanical

strain.<sup>18</sup> In our investigation of electrodeposited Ni<sub>90</sub>Fe<sub>10</sub> thin films, we have successfully induced the formation of stripe domains. However, this phenomenon occurs for thicknesses considerably larger (800 nm) than the theoretically inferred  $t_{cr}$  (105 nm) based on expression (1), and only when the electrolyte is not magnetically stirred during growth.<sup>17</sup> Typically, in-plane (IP) hysteresis loops exhibit the characteristic shape denoted as "transcritical" when stripe domains are present.<sup>5,20</sup> However, our observations show that electrodeposited Ni<sub>90</sub>Fe<sub>10</sub> layers display the transcritical shape for smaller thicknesses, where stripe domains are visible in magnetic force microscopy (MFM) images.<sup>17</sup> One of the aims of this study is to gain deeper insights into this thickness range where mixed features are observed in the experimental measurements. Such insights can provide valuable hints about the process of stripe domain formation.

Here, we have systematically investigated the evolution of the out-of-plane magnetic textures in electrodeposited Ni<sub>90</sub>Fe<sub>10</sub> layers with thickness ranging from 250 to 800 nm, where the formation of stripe domains has not been observed previously, but rather a magnetic ripple. This study focusses on examining the impact of various applied magnetic fields during the growth process on the magnetic behavior of these layers to gain a better understanding of stripe domain formation. To achieve this, we have not magnetically stirred during growth and compared layers deposited under three conditions: (i) zero applied field, (ii) the residual magnetic field generated by a switched-off magnetic stirrer, and (iii) a magnetic field of 100 Oe perpendicular to the sample plane (see Fig. 1). These represent completely new growth conditions in comparison to our previous work.<sup>17</sup> Our experimental findings reveal that the higher the perpendicular magnetic field applied during growth, the lower the critical thickness for the formation of stripe domains. Due to the gradual evolution, the observed features in hysteresis loops can vary depending on growth conditions and thickness; transcritical shape may be observed in hysteresis loops, while magnetic ripple appears in MFM images.

# **II. EXPERIMENTAL SECTION**

 $Ni_{90}Fe_{10}$  layers were grown by electrodeposition on glass substrates covered with a gold layer to increase the electrical conductivity. Samples were deposited at room temperature and not stirred during growth. We have used a three-electrode cell with a platinum mesh as counter electrode and an Ag/AgCl (3M NaCl) reference electrode. A PalmSens EmStat3 + Blue potentiostat was used to perform the depositions at room temperature in non-rotating substrates in the horizontal position at different growth conditions: (i) at zero applied magnetic field, (ii) under the residual magnetic field ( $H_{res}$ ) that has an IP component of 30 Oe and OOP of 24 Oe, respectively, and (iii) under an applied external magnetic field of 100 Oe perpendicular to the sample plane ( $H_{per}$ ) produced by a Helmholtz coil (Fig. 1),

The growth time of NiFe was adjusted to reach the expected thicknesses (t) by means of the Faraday's law:

$$C = \frac{nFd}{M_{atom}} St,$$
 (2)

where *C* is the electric charge measured in the cathode, *n* is the number of electrons involved in the reduction reaction, *F* is Faraday's constant (96 485.34 C/mol), *d* is the density of the electrodeposited material, *S* is the area of the sample, and  $M_{atom}$  is the molecular mass. First, a set of calibration samples was deposited in which the layer thickness was measured by atomic force microscopy (AFM). These samples were used to adjust the formula of Faraday's law to take into account hydrogen evolution.

The electrolyte for  $Ni_{90}Fe_{10}$  was only slightly modified taken into account previous literature.<sup>21,22</sup> It is water-based with a fixed pH of 2.2, and a composition of 0.4M H<sub>3</sub>BO<sub>3</sub>, 0.017M saccharine, g 0.7M NiSO<sub>4</sub>, and 0.02M FeSO<sub>4</sub>. It was applied a growth potential of -1.2 V in all cases.

We have used x-ray diffractometry (XRD) in the Bragg–Brentano  $\frac{2}{50}$  configuration to study the structural properties. Measurements were  $\frac{2}{50}$  performed in a D8 Brucker equipment using the Cu K<sub> $\alpha$ </sub> wavelength



FIG. 1. Scheme of the different growth conditions studied: (a) no external applied magnetic field, (b) residual magnetic field generated by the stirrer when it is switched off, and (c) under an applied field of 100 Oe.



**FIG. 2.** IP hysteresis loops for Ni<sub>90</sub>Fe<sub>10</sub> layers with different thicknesses deposited under the magnetic field generated by the magnetic stirrer when it is switched off ( $H_{res}$ ) ( $\bullet$ ) 250 nm, ( $\blacksquare$ ) 300 nm, ( $\Delta$ ) 350 nm, ( $\bigcirc$ ) 400 nm, ( $\blacktriangle$ ) 500 nm, and ( $\nabla$ ) 800 nm.

(1.540 56 Å). IP and OOP hysteresis loops were performed in a vibrating sample magnetometer (VSM) at room temperature. AFM and MFM measurements were performed under standard conditions using a customized and modified system from Nanotec Electrónica S.L. The data analysis process was carried out using WSxM software.<sup>23</sup> The amplitude modulation (AM) method was employed keeping the cantilever oscillation amplitude at about 20 nm for MFM imaging. The resulting magnetic signal was recorded in the phase shift channel acquired at a typical distance of 60 nm. Commercial Budget Sensors MagneticMulti75 probes with CoCr coating were used for the measurements. The magnetic signal of the probes was always checked before and after each experiment. For the roughness study, highly doped silicon probes with a typical tip radius of curvature of less than 10 nm from Nanosensors (PPP-FMR) have been used.

#### **III. RESULTS AND DISCUSSION**

First, we measured the roughness of the studied samples using AFM. Regardless of the growth conditions, a root mean square of approximately 7 nm was inferred. In Fig. 2, we present the IP hysteresis loops for layers with thicknesses ranging from 250 to 800 nm deposited under  $H_{res}$  generated by the magnetic stirrer when it is switched off. It is important to note that the magnetic bar is not placed inside the electrolyte cell for this growth condition [Fig. 1(b)]. Although none of the studied layers were magnetically stirred during growth, it is important to investigate this growth condition because switching off the magnetic stirrer does not imply that the layers are deposited under a zero magnetic field. This is due to the existence of a residual magnetic field when the stirrer is switched off.

The first noteworthy point is the decrease in the coercive field  $(H_{\rm C})$  from 61 Oe for a thickness of 250–23 Oe when the thickness is equal to or above 400 nm (Fig. 2). This reduction in  $H_C$  is accompanied by the emergence of a somewhat linear behavior in the hysteresis loop, which can be associated with the onset of the transcritical shape in the IP hysteresis loop for a thickness of 400 nm. From XRD diffraction patterns (Fig. 3), we determined the grain size (D) using Scherrer's formula. The observed changes in both  $H_C$  and hysteresis loop shape cannot be attributed to variations in the grain size, as D remains between 20.0 and 22.6 nm in all these samples. While the transcritical shape is somewhat present in the hysteresis loop for a thickness of 400 nm but not for 300 nm (Fig. 2), MFM images reveal only an enhancement of the magnetic ripple when comparing 300-400 nm, with no evidence of magnetic a stripes (Fig. 4). According to the Stoner-Wohlfarth model, hard E magnetic axes are known to exhibit a smaller coercivity compared to easy ones.<sup>16,24</sup> Therefore, the decrease in  $H_C$  together with the  $\frac{1}{2}$ presence of a linear behavior in the IP hysteresis loop can be asso- $\frac{1}{2}$ ciated with an enhancement of the PMA for a thickness of 400 nm. However, the PMA for this thickness (400 nm) is not significant



**FIG. 3.** (a) XRD diffraction patterns for Ni<sub>90</sub>Fe<sub>10</sub> layers deposited under the magnetic field generated by the magnetic stirrer when it is switched off ( $H_{res}$ ) with different thicknesses ranging from 300 to 800 nm. (b) XRD diffraction patterns for 500 nm thick Ni<sub>90</sub>Fe<sub>10</sub> layers deposited under the different magnetic field conditions studied in this work: H = 0 Oe,  $H_{res}$ , and  $H_{per}$ . Curves in (a) and (b) have been shifted for clarity.



FIG. 4. MFM images for Ni<sub>90</sub>Fe<sub>10</sub> samples deposited under the magnetic field generated by the magnetic stirrer when it is switched off ( $H_{res}$ ) with a thickness of (a) 300 nm and (b) 400 nm. Images size:  $5 \times 5 \mu m^2$ .



**FIG. 5.** (a) IP hysteresis loops for Ni<sub>90</sub>Fe<sub>10</sub> layers with a thickness of 300 nm deposited under different growth conditions: ( $\blacksquare$ )  $H_{res}$ , ( $\square$ )  $H_{per}$  = 100 Oe, and ( $\bullet$ ) zero field. MFM images for samples with 300 nm deposited under different growth conditions: (b)  $H_{res}$ , (c)  $H_{per}$  = 100 Oe, and (d) zero field. Images size: 10 × 10  $\mu$ m<sup>2</sup>.

**TABLE I.** Summary of the OOP magnetic anisotropy constant energy for layers deposited under different growth conditions in terms of applied magnetic field:  $H_{res}$  refers to the residual magnetic field generated by the magnetic stirrer when it is switched off, and  $H_{per}$  is the perpendicular to the sample plane of 100 Oe generated by the Helmholtz coil.

	$H_{res}$	<i>H</i> <sub>per</sub> =100 Oe	H = 0 Oe
Thickness (nm)	$K_{OOP}$ (kerg/cm <sup>3</sup> )		
300	17	27	
400	18	28	
500	18	28	

enough to promote the formation of stripe domains. Instead, an increase in the magnetic ripple is observed in the MFM image compared to 300 nm (Fig. 4), despite the transcritical shape being somewhat present in the IP hysteresis loop for 400 nm (Fig. 2).

In our investigation into the potential effect of magnetic fields during growth, we conducted electrodeposition of layers both under an applied perpendicular magnetic field of 100 Oe and without any magnetic field. The XRD diffraction patterns show no apparent differences regardless of the magnetic field condition during growth [Fig. 3(b)]. Lattice parameter and grain size are neither significantly modified. As an example, a comparison is presented among 500 nm thick layers deposited under zero field,  $H_{res}$ , and  $H_{per}$ .

In the hysteresis loops we observed that when the layer is deposited under  $H_{per}$ , there is a notable decrease in coercivity from 53 to 22 Oe for a smaller thickness of 300 nm, compared to the 400 nm thickness required to observe the same behavior when depositing under the residual field  $H_{res}$  [as depicted in Fig. 5(a)]. In contrast, layers deposited under zero magnetic field exhibit an IP hysteresis loop characteristic of a soft magnetic thin film. This observation further supports the notion that the perpendicular magnetic field applied during growth enhances the PMA. This conclusion is supported by the calculation of  $K_{OOP}$  (refer to Table I), which was quantitatively inferred from both IP and OOP hysteresis loops, following the same procedure as outlined in prior studies.<sup>17,25</sup> Specifically, for the 300 nm thickness,  $K_{OOP}$  is 17 kerg/cm<sup>3</sup> for the layer deposited under  $H_{res}$ , while it increases to 27 kerg/cm<sup>3</sup> for the layer deposited under  $H_{per}$  (as shown in Table I).



**FIG. 6.** (a) IP hysteresis loops for Ni<sub>90</sub>Fe<sub>10</sub> layers with a thickness of 400 nm deposited under different growth conditions: (b)  $H_{res}$ , (c)  $H_{per}$  = 100 Oe, and (e) zero field. MFM images for samples with 400 nm deposited under different growth conditions: (b)  $H_{res}$ , (c)  $H_{per}$  = 100 Oe, and (d) zero field. Images size: 10 × 10  $\mu$ m<sup>2</sup>.



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**FIG. 7.** (a) IP hysteresis loops for Ni<sub>90</sub>Fe<sub>10</sub> layers with a thickness of 500 nm deposited under different growth conditions: (**A**)  $H_{res}$ , (**C**)  $H_{per} = 100$  Oe, and (**O**) zero field. MFM images for samples with 500 nm deposited under different growth conditions: (b)  $H_{res}$ , (c)  $H_{per} = 100$  Oe, and (d) zero field. Images size:  $10 \times 10 \,\mu\text{m}^2$ .

Notably, it is not feasible to calculate  $K_{OOP}$  when no magnetic field is applied during growth. Furthermore, all 300 nm-thick layers exhibit a similar MFM magnetic contrast, suggesting that the PMA, even in the sample deposited under  $H_{per}$ , is not significant enough to promote stripe domains [as depicted in Figs. 5(b)–5(d)].

If we compare the hysteresis loops of the layers deposited under  $H_{per}$  with thicknesses of 300 and 400 nm [Figs. 5(a) and 6(a)], a more distinct transcritical shape is observed in the layer of 400 nm. For both  $H_{per}$  and  $H_{res}$  conditions,  $K_{OOP}$  only slightly increases when the thickness is increased from 300 to 400 nm (Table I). Nevertheless, there is a much greater enhancement of  $K_{OOP}$  when evaluating samples with the same thickness but deposited under  $H_{per}$  compared to  $H_{res}$  (Table I). Therefore, applying a perpendicular magnetic field during growth is more efficient than increasing the thickness to enhance the PMA. Layers deposited under zero field with 300 and 400 nm do not exhibit any evolution [Figs. 5(a) and 6(a)]; coercivity remains at 14 Oe, and the shape of the hysteresis remains consistent. When comparing MFM images for 400 nm layers [Figs. 6(b)-6(d)], stripe domains can be observed in the layer deposited under  $H_{per}$ , whereas the layer deposited under  $H_{res}$  only exhibits a magnetic ripple that is somewhat more intense than in the layer deposited under zero field.

The impact of the applied magnetic field in Ni<sub>90</sub>Fe<sub>10</sub> layers is further confirmed for a thickness of 500 nm (Fig. 7). The comparison between IP and OOP hysteresis loops is shown in Fig. 8. When the transcritical shape is present in the IP loops, the remanence in the OOP magnetization curves is almost zero due to the presence of stripe domains, in agreement with previously reported works.<sup>5,26</sup> The layer deposited under  $H_{per}$  clearly displays both the transcritical shape in the hysteresis loop and the presence of stripe domains in the MFM image. In contrast, the sample deposited under  $H_{res}$ exhibits a hysteresis loop quite similar to that of 400 nm, although the magnetic ripple is more pronounced in the MFM image. The layer deposited under zero fields continues to exhibit a magnetic ripple in the MFM, and the IP hysteresis loop still lacks features associated with the transcritical shape.

The formation of stripe domains is attributed to the presence of significant PMA, which can counterbalance demagnetizing energy.<sup>5,6,25</sup> While the transcritical shape in IP hysteresis loops is typically considered an indication of stripe domain formation, we have





experimentally observed a progressive evolution of hysteresis loops shape toward the transcritical shape without clear evidence of stripes in the MFM images. In the study of Cao and co-workers on electrodeposited  $Ni_{50}Fe_{50}$  layers,<sup>27</sup> a thickness range was identified where it was challenging to infer the presence of stripe domains from hysteresis loops alone. Similar observations were reported for FePt, where the transition is gradual, and layers close to the critical thickness can show mixed features.<sup>25</sup> These findings align with our results, as layers can exhibit the transcritical shape in IP hysteresis loops without clear evidence of stripe domains in MFM. Moreover, in this transition regime, the magnetic contrast is characterized by a ripple that is due to magnetic fluctuations, as also observed in FeGa films.<sup>14,17</sup>

Typically, the quality of factor (*Q*) is used to quantitatively classify materials with PMA, defined  $as^6$ 

$$Q = K_{OOP} / 2\pi M_{sat}^2, \tag{3}$$

where  $M_{sat}$  is the saturation magnetization. Q < 1 indicates materials with moderate or low PMA. With  $M_{sat} = 590 \text{ emu/cm}^3$  for Ni<sub>90</sub>Fe<sub>10</sub>,<sup>17</sup> Q ranges from 0.008 for the layers deposited under the residual field to 0.012–0.013 for those deposited under  $H_{per}$ . These quality factors are smaller compared to systems with higher PMA, where stripe domains are visible in the MFM images and Q values are around 0.3. For example, Fe–N reports values of 0.28–0.38,<sup>5</sup> and the heterostructure Cr<sub>2</sub>O<sub>3</sub>/FeGa obtains 0.3.<sup>28</sup> However, Q values for Ni<sub>90</sub>Fe<sub>10</sub> are comparable to works on other soft magnetic alloys where stripes are present, with  $Q \approx 0.02 - \approx 0.005$ , even lower than in our study.<sup>29</sup>

The critical thickness for the formation of stripe domains in  $Ni_{50}Fe_{50}$  layers<sup>27</sup> closely aligns with previously published data, where columnar growth also appears to play a pivotal role in enhancing PMA.<sup>30</sup> However, various parameters such as substrate temperature or deposition pressure can alter  $t_{cr}$  in sputtered layers,<sup>25,31</sup> underscoring the sensitivity of the PMA to growth conditions. In our electrodeposited  $Ni_{90}Fe_{10}$  layers, we have observed that  $t_{cr}$  for the formation of stripe domains depends on the applied magnetic field during growth. Higher perpendicular field applied during electrodeposition leads to a larger PMA, enabling to obtain stripe domains at a smaller thickness than previously reported.<sup>17</sup> Indeed, applying a perpendicular magnetic field

during growth proves to be more effective than increasing the thickness to enhance the PMA.

#### **IV. CONCLUSIONS**

The combined use of magnetic magnetometry and MFM characterization has provided valuable insights into the formation of stripe domains in electrodeposited Ni<sub>90</sub>Fe<sub>10</sub> thin films. When no magnetic field is applied during growth, neither stripe domains nor the transcritical shape are observed, even for a thickness exceeding 500 nm. However, applying a perpendicular magnetic field of 100 Oe during growth results in the presence of stripes for a thick-  $\frac{1}{2}$ ness of 400 nm. Layers deposited under a magnetic field with IP ≦ and OOP components  $(H_{res})$  exhibit a mixed behavior. The transcritical shape is observed in the hysteresis loop for a thickness of 8 400 nm, but MFM only detects magnetic ripple, even a thickness of  $\frac{4}{2}$ 500 nm. In summary, our experimental findings on  $Ni_{90}Fe_{10}$  elecmagnetic texture is more efficiently promoted by an external perpendicular magnetic field applied during growth than by the increase in thickness. Therefore, this work opens new routes to tailor magnetic textures in Ni<sub>90</sub>Fe<sub>10</sub> films.

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#### AUTHOR DECLARATIONS

#### **Conflict of Interest**

The authors have no conflicts to disclose.

# **Author Contributions**

**N. Cotón:** Data curation (equal); Investigation (lead); Methodology (equal); Writing – review & editing (equal). **J. P. Andrés:** Funding acquisition (equal); Investigation (equal); Methodology (equal); Writing – review & editing (equal). **M. Jaafar:** Funding acquisition (equal); Investigation (equal); Methodology (equal); Writing – review & editing (equal). **A. Begué:** Investigation (equal); Methodology (equal); Writing – review & editing (equal). **R. Ranchal:** Conceptualization (lead); Funding acquisition (equal); Supervision (lead); Writing – original draft (lead); Writing – review & editing (equal).

## DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

#### REFERENCES

<sup>1</sup>R. Wiesendanger, "Nanoscale magnetic skyrmions in metallic films and multilayers: A new twist for spintronics," Nat. Rev. Mater. **1**, 16044 (2016).

<sup>2</sup>D. Petti, S. Tacchi, and E. Albisetti, "Review on magnonics with engineered spin textures," J. Phys. D: Appl. Phys. **55**, 293003 (2022).

<sup>3</sup>S. Luo, M. Song, X. Li, Y. Zhang, J. Hong, X. Yang, X. Zou, N. Xu, and L. You, "Reconfigurable skyrmion logic gates," Nano Lett. **18**, 1180–1184 (2018).

<sup>4</sup>M. Krawczyk and D. Grundler, "Review and prospects of magnonic crystals and devices with reprogrammable band structure," J. Phys.: Condens. Matter **26**, 123202 (2014).

<sup>5</sup>L.-C. Garnier, M. Marangolo, M. Eddrief, D. Bisero, S. Fin, F. Casoli, M. G. Pini, A. Rettori, and S. Tacchi, "Stripe domains reorientation in ferromagnetic films with perpendicular magnetic anisotropy," J. Phys. Mater. 3, 024001 (2020).

<sup>6</sup>A. Hubert and R. Schäfer, *Magnetic Domains: The Analysis of Magnetic Microstructures* (Springer Science & Business Media, 2008).

<sup>7</sup>C. Liu, S. Wu, J. Zhang, J. Chen, J. Ding, J. Ma, Y. Zhang, Y. Sun, S. Tu, H. Wang, P. Liu, C. Li, Y. Jiang, P. Gao, D. Yu, J. Xiao, R. Duine, M. Wu, C.-W. Nan, J. Zhang, and H. Yu, "Current-controlled propagation of spin waves in antiparallel, coupled domains," Nat. Nanotechnol. 14, 691 (2019).

<sup>8</sup>J. Ben Youssef, N. Vukadinovic, D. Billet, and M. Labrune, "Thickness-dependent magnetic excitations in permalloy films with nonuniform magnetization," Phys. Rev. B **69**, 174402 (2004).

<sup>9</sup>M. Liu, S. Du, F. Wang, R. Adam, Q. Li, X. Ma, X. Guo, X. Chen, J. Yu, Y. Song, J. Xu, S. Li, and D. Cao, "Influence of surface pinning in the domain on the magnetization dynamics in permalloy striped domain films," J. Alloys Compd. **869**, 159327 (2021).

<sup>10</sup>M. Liu, Q. Li, C. Song, H. Feng, Y. Song, L. Zhong, L. Pan, C. Zhao, Q. Li, J. Xu, S. Li, J. Wang, Q. Liu, and D. Cao, "Microwave excitations and hysteretic magnetization dynamics of stripe domain films," J. Magn. Magn. Mater. 547, 168939 (2022).

<sup>11</sup>P. Bartolomé, M. Maicas, and R. Ranchal, "Out of plane component of the magnetization of sputtered Fe<sub>72</sub>Ga<sub>28</sub> layers," J. Magn. Magn. Mater. **514**, 167183 (2020).

<sup>12</sup>K. J. Harte, "Theory of large angle ripple in magnetic films," J. Appl. Phys. 37, 1295–1296 (1966).

<sup>13</sup>H. W. Fuller and M. E. Hale, "Determination of magnetization distribution in thin films using electron microscopy," J. Appl. Phys. **31**, 238–248 (1960).

<sup>14</sup>A. Begué, M. G. Proietti, J. I. Arnaudas, and M. Ciria, "Magnetic ripple domain structure in FeGa/MgO thin films," J. Magn. Magn. Mater. **498**, 166135 (2020).

<sup>15</sup>R. Ranchal, S. Fin, and D. Bisero, "Magnetic microstructures in electrodeposited  $\text{Fe}_{1-x}\text{Ga}_x$  thin films ( $15 \le x \le 22 \text{ at.}\%$ )," J. Phys. D: Appl. Phys. **48**, 075001 (2015).

 $^{17}$ N. Cotón, J. P. Andrés, E. Molina, M. Jaafar, and R. Ranchal, "Stripe domains in electrodeposited Ni\_{90}Fe\_{10} thin films," J. Magn. Magn. Mater. **565**, 170246 (2023).

 $^{18}$  M. Ciria, K. Ha, D. Bono, and R. C. O'Handley, "Magnetoelastic coupling in epitaxial Cu/Ni\_{90}Fe\_{10}/Cu/Si(001) thin films," J. Appl. Phys. **91**, 8150–8152 (2002).

<sup>19</sup>M. Ciria and R. C. O'Handley, "Surface- and strain-induced anisotropy in epitaxial Cu/Ni<sub>90</sub>Fe<sub>10</sub>/Cu/Si(001) films," J. Magn. Magn. Mater. **240**, 464–466 (2002).

<sup>20</sup>N. Saito, H. Fujiwara, and Y. Sugita, "A new type of magnetic domain structure in negative magnetostriction Ni-Fe films," J. Phys. Soc. Jpn. **19**, 1116–1125 (1964).

<sup>21</sup>L. J. Gao, P. Ma, K. M. Novogradecz, and P. R. Norton, "Characterization of Permalloy thin films electrodeposited on Si(111) surfaces," J. Appl. Phys. 81, 7595 (1997).

<sup>22</sup>W. Bang, Y.-D. Ko, H. Lee, K. Hong, J.-S. Chung, and H. Lee, "Effects of saccharine N-propane sulfonate on the microstructures, magnetic properties, and magnetoimpedance of electroplated Ni–Fe Permalloy thin films," J. Electrochem. Soc 155, D429 (2008).

<sup>23</sup>I. Horcas, R. Fernández, J. M. Gómez-Rodríguez, J. Colchero, J. Gómez-Herrero, and A. M. Baró, "WSXM: A software for scanning probe microscopy and a tool for nanotechnology," Rev. Sci. Instrum. 78, 013705 (2007).

<sup>24</sup>M. Maicas, R. Ranchal, C. Aroca, P. Sánchez, and E. López, "Magnetic properties of Permalloy multilayers with alternating perpendicular anisotropies," Eur. *B* Phys. J. B **62**, 267–270 (2008).

<sup>26</sup>D. Wu, T. Jin, Y. Lou, and F. Wei, "Understanding the dense stripe domains in soft magnetic film," Appl. Surf. Sci. 346, 567–573 (2015).

27 D. Cao, Z. Wang, E. Feng, J. Wei, J. Wang, and Q. Liu, "Magnetic properties and microstructure investigation of electrodeposited FeNi/ITO films with different thickness," J. Alloys Compd. 581, 66–70 (2013).

<sup>28</sup>I. Hontecillas, M. Maicas, J. P. Andres, and R. Ranchal, "Interfacial coupling effect of  $Cr_2O_3$  on the magnetic properties of  $Fe_{72}Ga_{28}$  thin films," Sci. Rep. 11, 13429 (2021).

<sup>29</sup>M. Coïsson, G. Barrera, F. Celegato, and P. Tiberto, "Rotatable magnetic anisotropy in  $Fe_{78}Si_9B_{13}$  thin films displaying stripe domains," Appl. Surf. Sci. **476**, 402–411 (2019).

<sup>30</sup>M. Romera, R. Ranchal, D. Ciudad, M. Maicas, and C. Aroca, "Magnetic properties of sputtered Permalloy/molybdenum multilayers," J. Appl. Phys. 110, 083910 (2011).

<sup>31</sup>N. Amos, R. Fernandez, R. Ikkawi, B. Lee, A. Lavrenov, A. Krichevsky, D. Litvinov, and S. Khizroev, "Magnetic force microscopy study of magnetic stripe domains in sputter deposited Permalloy thin films," J. Appl. Phys. 103, 07E732 (2008).