

Daily modulations and broadband strategy in axion searches: An application with CAST-CAPP detector

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It has been previously advocated that the presence of the daily and annual modulations of the axion flux on the Earth's surface may dramatically change the strategy of the axion searches. The arguments were based on the so-called Axion Quark Nugget (AQN) dark matter model which was originally put forward to explain the similarity of the dark and visible cosmological matter densities $\Omega_{\text{dark}} \sim \Omega_{\text{visible}}$. In this framework, the population of galactic axions with mass $10^{-6}\text{eV} \lesssim m_a \lesssim 10^{-3}\text{eV}$ and velocity $\langle v_a \rangle \sim 10^{-3}c$ will be accompanied by axions with typical velocities $\langle v_a \rangle \sim 0.6c$ emitted by AQNs. Furthermore, in this framework, it has also been argued that the AQN-induced axion daily modulation (in contrast with the conventional WIMP paradigm) could be as large as (10–20)%, representing the main motivation for the present investigation. We argue that the daily

modulations along with the broadband detection strategy can be very useful tools for the discovery of such relativistic axions. The data from the CAST-CAPP detector have been used following such arguments. Unfortunately, due to the dependence of the amplifier chain on temperature-dependent gain drifts and other factors, we could not conclusively show the presence or absence of a dark sector-originated daily modulation. However, this proof of principle analysis procedure can serve as a reference for future studies.

I. INTRODUCTION

The Peccei-Quinn mechanism, accompanied by axions, offers the most compelling solution to the strong \mathcal{CP} problem; see foundational works [1–7] and recent reviews [8–18]. The conventional approach for producing dark matter (DM) axions relies on either the misalignment mechanism—where the cosmological field $\theta(t)$ oscillates and emits cold axions before reaching a minimum or the decay of topological defects, as outlined in recent reviews [8–18]. In both cases, the resulting axions are nonrelativistic, with velocities of approximately $\langle v_a \rangle \sim 10^{-3}c$.

Beyond these established mechanisms, recent studies [19–23] have proposed a novel mechanism for axion production based on the axion quark nugget (AQN) DM model [24], briefly reviewed in Sec. II. This model shares similarities with Witten’s quark-nugget model [25], but is “cosmologically dark” due to its low cross section-to-mass ratio, which reduces observable consequences of what is otherwise a strongly interacting DM candidate. An essential feature of the AQN model is the production of relativistic axions with velocities of $\langle v_a \rangle \sim 0.6c$, in contrast to conventional mechanisms where axion velocities are typical of galactic DM ($\langle v_a \rangle \sim 10^{-3}c$), as detailed in Sec. III.

This unique characteristic necessitates a broadband search for axion detection, markedly changing the detection strategy. Such broadband searches could incorporate additional elements, including a specific analysis of daily modulation effects (detailed in Sec. IV).

Relativistic axion production has also been discussed outside of the AQN framework, as seen in Ref. [26], where several production mechanisms were suggested and collectively termed the cosmic axion background (CaB). Thus, our proposed broadband search strategy could apply more broadly and is not limited solely to the AQN model.

Despite these parallels, several distinctions exist between AQN-induced axions and CaB axions as considered in Ref. [26]. AQN-induced axions arise from interactions between AQNs and the Earth’s material, meaning they are largely localized near the Earth’s (or other celestial bodies’) surface. In contrast, the CaB axions in Ref. [26] are cosmologically distributed throughout the Universe.

Consequently, the dominant portion of the relativistic axions (which is the topic of the present work) should be related to the AQN-induced events with a well-defined estimate for the axion density on the Earth’s surface as reviewed in Sec. III.

Another key difference concerns daily modulation. The CaB-induced daily modulation observed in Ref. [26] requires directional sensitivity in detectors, as the modulation arises from the shifting orientation of Earth’s axis relative to the galactic wind. Although current instruments were not specifically designed for directional detection, recent studies like [27] have utilized Axion Dark Matter Experiment (ADMX) data to examine daily modulation due to CaB. For AQN-induced axions, an additional source of daily modulation arises from the changing mass (and size) of AQNs entering or exiting Earth’s atmosphere, as discussed in Sec. III. This additional modulation allows the study of daily modulation effects even with spherically symmetric, nondirectionally sensitive instruments.

Furthermore, the AQN-induced axion density near the Earth’s surface is expected to surpass the CaB axion density, which is constrained by the cosmic microwave background (CMB) photon density (see Sec. III for estimates). Therefore, time modulation effects, which are the topic of the present study, will be dominated by AQN-induced axions.

One additional feature of AQN-induced axions that differs from CaB mechanisms in [26] is that AQNs, being massive objects, produce a highly nonuniform flux on Earth. This nonuniformity, which is in contrast with the conventional DM models, and the CaB flux, can create “local flashes” of axions. This can temporarily increase their count by several orders of magnitude when an AQN impacts the Earth’s surface within about 100 km of the detector (see Sec. II for further estimates).

This paper is organized as follows: Section II provides an overview of the AQN framework, while Secs. III and IV describe the main properties of AQN-induced axions relevant to this study, including their spectrum and intensity at Earth’s surface. Section V presents our estimates of daily modulation effects and suggests methods to distinguish genuine signals from noise. In Sec. VI, we apply these analytical features to explore daily modulation in previously recorded CAST-CAPP data. We conclude in Sec. VII, with recommendations for future studies.

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II. AQN DARK MATTER MODEL: THE BASICS

The AQN model shares several key features with Witten’s quark nugget construction [25, 28, 29]. Like quark nuggets, AQNs are “cosmologically dark” due to the very low cross-section-to-mass ratio relevant for cosmological interactions. This small ratio effectively limits the observable impacts of these otherwise strongly interacting DM candidates.

The AQN model, however, incorporates additional elements beyond the original quark nugget framework [25, 28, 29]. First, AQNs are stabilized by axion domain walls, which form abundantly during the QCD phase transition and mitigate several issues associated with Witten’s original model. Second, during this QCD transition, AQNs can form as either *matter* or *antimatter* nuggets, adding further diversity to their composition.

The presence of antimatter nuggets in the AQN framework is an inevitable and direct consequence of the \mathcal{CP} violating axion field which is present in the system during the QCD time. As a result of this feature, the DM density Ω_{DM} , and the visible density Ω_{visible} , will automatically assume the same order of magnitude densities $\Omega_{\text{DM}} \sim \Omega_{\text{visible}}$ irrespectively to the parameters of the model, such as the axion mass m_a . This feature represents a generic property of the construction [24] as both components, the visible, and the dark are proportional to the same fundamental dimensional constant of the theory, the Λ_{QCD} .

For detailed discussions on the formation of AQNs, the generation of baryon asymmetry, and the survival mechanisms of these nuggets in the hostile environment of the early Universe, we refer to the original studies [36–39]. Additionally, a recent review article [30] provides insights into the formation and survival pattern of AQNs during the early stages of the evolution, including the CMB Big Bang Nucleosynthesis (BBN), and recombination epochs.

In this work, we take an agnostic stance and assume that *antimatter*-based nuggets exist in our Universe today and have existed since earlier cosmic times, independent of any specific formation mechanism. This assumption aligns with all current cosmological, astrophysical, and terrestrial constraints, provided that the average baryon charge of these nuggets is sufficiently large, as we review below.

The most stringent direct detection limit¹ comes from the IceCube Observatory; see Appendix A in Ref. [21]:

$$\langle B \rangle > 3 \cdot 10^{24} \quad [\text{direct (non)detection constraint}]. \quad (1)$$

Similar limits can also be obtained from ANITA and geothermal constraints, which are consistent with Eq. 1, as estimated in Ref. [41].

¹ Nondetection of etching tracks in ancient mica offers an indirect constraint on the flux of DM nuggets with masses above $M > 55g$ [40]. This constraint assumes that all nuggets have identical mass, which does not apply to the AQN model.

Assuming the conventional halo model with a local DM density of $\rho_{\text{DM}} \simeq 0.3 \text{ GeV cm}^{-3}$, the rate at which AQNs hit the Earth [21] is

$$\begin{aligned} \frac{\langle \dot{N} \rangle}{4\pi R_{\oplus}^2} &\simeq \frac{4 \cdot 10^{-2}}{\text{km}^2 \text{ yr}} \left(\frac{\rho_{\text{DM}}}{0.3 \frac{\text{GeV}}{\text{cm}^3}} \right) \left(\frac{\langle v_{\text{AQN}} \rangle}{220 \frac{\text{km}}{\text{s}}} \right) \left(\frac{10^{25}}{\langle B \rangle} \right), \\ \langle \dot{N} \rangle &\simeq 0.67 \text{ s}^{-1} \left(\frac{10^{25}}{\langle B \rangle} \right) \simeq 2.1 \cdot 10^7 \text{ yr}^{-1} \left(\frac{10^{25}}{\langle B \rangle} \right) \quad (2) \end{aligned}$$

In obtaining Eq. 2 it was assumed that $\langle \dot{N} \rangle \approx 4\pi b^2 n_{\text{AQN}} v_{\text{AQN}}$, where $b \approx 1.0013 R_{\oplus}$ is the impact parameter for Earth, and we averaged over incident angles of impacting AQNs; see Appendix C in Ref. [21] with details. We also assume that the AQNs represent the dominant portion of the DM. Numerical expressions for Eq. 2 were derived in Ref. [21] using full-scale Monte Carlo simulations which account for all types of AQN trajectories with different AQN masses $M_N \simeq m_p |B|$, different incident angles, and different initial velocities and size distributions. It was shown that none of these factors significantly affects our estimate. The result in Eq. 2 suggests that AQNs with a typical size $B \approx 10^{25}$ hit the Earth’s surface with a frequency on the order of once a day (i.e., hundreds per year) per $100 \times 100 \text{ km}^2$.

We conclude this brief review subsection with Table I, which summarizes the fundamental features and parameters of the AQNs. The parameter κ in Table I accounts for the fact that not all matter interacting with the nugget will annihilate, and not all of the energy released during the annihilation will be thermalized within the nuggets. The ratio $\Delta B/B \ll 1$ in the Table I indicates that only a small portion, ΔB , of the total (anti)baryon charge B hidden in the form of AQNs gets annihilated during BBN, CMB, or post-recombination epochs (including galaxy and star formation), while the majority of the baryon charge survives to the present day. Independent analyses [34] and [33] further support our original claims, as cited in Table I, that the antiquark nuggets survive the BBN and CMB epochs.

III. AQN-INDUCED AXIONS

As previously noted, AQNs traveling through the Earth’s atmosphere and interior lose baryon charge due to annihilation processes, which in turn lead to axion emission as proposed in Ref. [19]. This novel axion production mechanism can be outlined as follows: the total energy of an AQN reaches an equilibrium minimum when the axion domain wall contributes roughly one-third of the AQN’s total mass [39]. In this equilibrium state, the configuration does not emit axions due to a purely kinematic constraint; the static axion domain wall can be considered a superposition of off-shell, nonpropagating axions. However, this static balance is disrupted when a baryon charge is annihilated through interactions with the environment, causing time-dependent perturbations

Property	Typical value or feature
AQN's mass $[M_N]$	$M_N \approx 16 g (B/10^{25})$ [30]
Baryon charge constraints $[B]$	$B \geq 3 \cdot 10^{24}$ [30]
Annihilation cross section $[\sigma]$	$\sigma \approx \kappa \pi R^2 \simeq 1.5 \cdot 10^{-9} \text{cm}^2 \cdot \kappa (R/2.2 \cdot 10^{-5} \text{cm})^2$
Density of AQNs $[n_{\text{AQN}}]$	$n_{\text{AQN}} \sim 0.3 \cdot 10^{-25} \text{cm}^{-3} (10^{25}/B)$ [30]
Survival pattern during BBN	$\Delta B/B \ll 1$ [31–34]
Survival pattern during CMB	$\Delta B/B \ll 1$ [31, 33, 35]
Survival pattern during post-recombination	$\Delta B/B \ll 1$ [36]

TABLE I. Basic properties of the AQNs

that shift the AQN out of equilibrium. In essence, the AQN configuration becomes unstable and prone to axion emission because it no longer resides at its minimum energy with fewer baryon charges in the quark nugget core. As annihilation proceeds, the AQN gradually loses mass, leading the axion domain wall around the quark matter core to oscillate and contract. These domain wall oscillations generate excitation modes that eventually result in the emission of propagating axions. The spectrum for these AQN-induced axions has been calculated in Ref. [20].

Several key features of this spectrum merit mention. First, the typical velocities $\langle v_a^{\text{AQN}} \rangle \simeq 0.6c$ are significantly higher than those of conventional galactic nonrelativistic axions, which have $\langle v_a \rangle \sim 10^{-3}c$. Monte Carlo simulations presented in Ref. [22] provide the following estimate for the density of these relativistic axions at the Earth's surface:

$$\langle \rho_a^{\text{AQN}}(R_\oplus) \rangle \sim 5 \cdot 10^{-6} A(t) \frac{\text{GeV}}{\text{cm}^3}, \quad \langle v_a^{\text{AQN}}(R_\oplus) \rangle \simeq 0.6c, \quad (3)$$

where $A(t)$ is the time-dependent modulation/amplification factor. The factor A for the daily and annual modulations is discussed in Sec. IV below, specifically in Eq. 6 and 7. In both cases, A does not deviate from the average value $\langle A \rangle = 1$ by more than 10%.

However, this normalization assumes the Standard Halo Model (SHM) distribution for DM, with a density of approximately $\rho_{\text{DM}} \approx 0.3 \text{ GeV}/\text{cm}^3$. The actual DM density on the Earth's surface could differ significantly from this standard value due to the gravitational capture of DM particles by stars and planets; see, for example, Ref. [42] and references therein. Indeed, one of the primary motivations in Ref. [43] was to critically examine the SHM assumption within the AQN framework through specific computations.

The AQN-induced axions are predominantly generated in the Earth's deep interior, where the density of surrounding matter is highest. Notably, the energy density of the axions [Eq. 3] and the energy flux [Eq. 4] are relatively insensitive to the axion mass m_a , as detailed in Ref. [22]. This behavior contrasts with that of galactic nonrelativistic axions, under the assumption that DM is primarily composed of these galactic axions.

It is useful to compare the energy density of the relativistic AQN-induced axions [Eq. 3] with that of CMB

photons, as both are relativistic components. The energy density of CMB photons is well established at $\rho_\gamma \approx 0.26 \text{ eV}/\text{cm}^3$, nearly 4 orders of magnitude lower than the AQN-induced axion density on Earth's surface (Eq. 3). By contrast, the density of CaB axions discussed in Ref. [26] and referenced in Sec. I must satisfy $\rho_{\text{CaB}} \ll \rho_\gamma$. This constraint implies that the CaB axion density is at least four orders of magnitude lower than the density of AQN-induced axions at Earth's surface [Eq. 3].

There is a significant difference in the distribution of CMB photons compared to the AQN-induced axion energy density. CMB photons are nearly uniformly distributed across the Universe, while the axion density [Eq. 3] is concentrated around the Earth and rapidly decreases at distances much greater than R_\oplus . Similar localized distributions can be estimated for other stars and planets. Unfortunately, axions interact far more weakly with visible matter than CMB photons, which poses a substantial challenge for their detection despite their higher density near the Earth (Eq. 3) relative to CMB photons.

The resulting number density $n_a^{\text{AQN}} \simeq \rho_a^{\text{AQN}}/m_a$ is approximately 5 orders of magnitude smaller than the conventional galactic axion number density assuming that the galactic nonrelativistic axions saturate the DM today. However, the flux ($v_a^{\text{AQN}} \cdot n_a^{\text{AQN}}$) related to relativistic axions (Eq. 3) is only 2 orders of magnitude smaller than the conventional flux of nonrelativistic axion. Indeed, the AQN-induced axion energy flux on the Earth's surface can be estimated as follows [22]:

$$\langle E_a \cdot \Phi_a^{\text{AQN}} \rangle \simeq 10^{14} A(t) \left[\frac{\text{eV}}{\text{cm}^2 \cdot \text{s}} \right], \quad \langle E_a \rangle \simeq 1.3 m_a. \quad (4)$$

It is instructive to compare the AQN-related flux [Eq. 4] with the flux computed from the assumption that the galactic axions saturate the DM density $\rho_{\text{DM}} \sim 0.3 \text{ GeV} \cdot \text{cm}^{-3}$ today. In this case, the numerical value for the flux [Eq. 4] is approximately 2 orders of magnitude below the value $10^{16} \text{ eV}/(\text{cm}^2 \cdot \text{s})$ computed for the conventional galactic nonrelativistic axions assuming these axions saturate the DM density. It is emphasized once more that the density [Eq. 3] as well as the flux [Eq. 4] do not strongly depend on the axion mass m_a nor the initial misalignment angle θ_0 in contrast with the conventional expression for the galactic axions.

A key characteristic of AQN-induced axions is their daily modulation, a property not typically found in

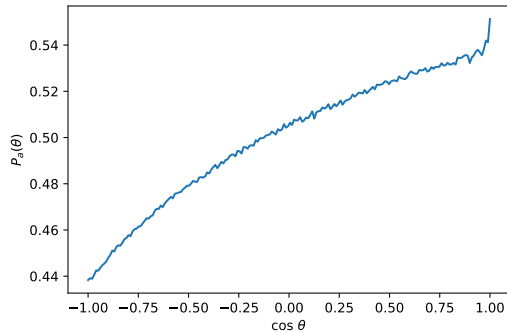


FIG. 1. The result of the Monte Carlo numerical simulations for the azimuthal distribution of axion flux $P_a(\theta)$ on the surface of the Earth, adopted from [22]. A relative orientation of the Earth's axis of rotation and DM galactic wind generates the flux asymmetry sensitive to the latitude θ , which consequently leads to the daily modulation effect.

generic DM models, especially when detectors lack sensitivity to directionality, as is common with most modern DM detectors. This daily modulation is a distinctive feature of the AQN-induced production mechanism and is central to our analysis.

The origin of this modulation lies in the difference in AQN size between its entry and exit points, resulting from annihilation processes within the Earth's interior. Such effects are absent for fundamental particles like weakly interacting massive particles (WIMPs), whose mass remains unchanged between entry and exit. In contrast, the average size difference for AQNs as they enter and exit Earth's surface can produce a significant daily modulation effect, approximately 10%, as calculated in Ref. [22] and discussed below.

The key factor driving the daily modulation here is the relative orientation between Earth's rotational axis and the galactic DM wind. Unlike WIMPs, where this orientation is irrelevant without directional detectors, AQN fluxes are sensitive to it, generating a flux asymmetry that varies with latitude, θ , and leads to daily modulation. This feature of the axion flux is captured by the parameter $P_a(\theta)$, which is normalized such that $\int P_a(\theta) \sin \theta d\theta = 1$. In the absence of daily modulation, $P_a(\theta)$ would be uniform at 0.5, independent of latitude. However, our numerical Monte Carlo simulations, shown in Fig. 1, clearly demonstrate a θ dependence that produces daily modulation.

It is worth noting that the daily modulation discussed here differs substantially from that in Ref. [27], where the directional sensitivity of the instrument plays a central role in detecting daily modulations from CaB. In our case, even a spherically symmetric detector, which lacks directional sensitivity, can still measure this daily modulation effect.

The production rate of the low-energy AQN-induced axions with $v_a \ll c$ is strongly suppressed as explained above. However, the axions that are produced with ex-

tremely low velocities $v_a \lesssim 11$ km/s will be trapped by the Earth's gravitational field. These axions will be orbiting the Earth indefinitely; therefore, they will be accumulated around the Earth during its lifetime of 4.5 billion years.

The corresponding Monte Carlo simulations have been performed in Ref. [21] with the following estimate for the gravitationally bound axions:

$$\rho_a^{\text{bound}}(R_\oplus) \sim 10^{-4} \frac{\text{GeV}}{\text{cm}^3}, \quad \langle v_a(R_\oplus) \rangle \simeq 8 \frac{\text{km}}{\text{s}}. \quad (5)$$

The number density of the bound axions is at least 2 orders of magnitude smaller than conventional axion number density assuming that the galactic nonrelativistic axions saturate the DM density. However, the corresponding wavelength $\lambda_a \sim \hbar/(m_a v_a)$ of the gravitationally bound axions is approximately 30 times greater than for galactic axions, which have a typical velocity of about ~ 250 km/s. Therefore, coherent effects can be maintained for a longer time period compared to those for conventional galactic axion searches. One may hope that the feature of having a large coherence length, $\lambda_a \sim v_a^{-1}$, could play a key role in the design of instruments, capable of discovering such gravitationally trapped axions.

These AQN-induced axions (relativistic as well as gravitationally bound) can be treated as a classical field because the number of the AQN-induced axions (Eq. 3) accommodated by a single de Broglie volume is very large even though the de Broglie wavelength λ for relativistic AQN-induced axions is much shorter than for galactic axions,

$$n_a^{\text{AQN}} \lambda^3 \sim \frac{\langle \rho_a^{\text{AQN}} \rangle}{m_a} \left(\frac{\hbar}{m_a v_a} \right)^3 \sim 10^6 \left(\frac{10^{-4} \text{eV}}{m_a} \right)^4 \gg 1,$$

where we use $m_a \sim 10^{-4} \text{eV}$ as a benchmark axion mass to be used for all numerical estimates in the present work.

IV. NEW STRATEGY: DETECTION OF BROADBAND AXIONS

The large average velocities $\langle v_a \rangle \simeq 0.6c$ of the emitted axions by AQNs change the entire strategy of axion searches since they are characterized by a broad distribution with $m_a \lesssim \omega_a \lesssim 1.8 m_a$ as discussed in the previous section. Therefore, the corresponding axion detectors must be some kind of broadband instrument. The cavity-type experiments such as ADMX [44] and CAPP [45] are to date the only ones to probe the parameter space of the conventional QCD axions with $\langle v_a \rangle \sim 10^{-3}c$, while we are interested in the detection of the relativistic axions with $\langle v_a \rangle \sim 0.6c$. This requires different types of instruments and drastically different search strategies. We assume that some kind of broadband instruments can be designed and built, see reviews [14–18] for a description of possible detectors.

With this assumption in mind, a new strategy to probe the QCD axion can be formulated as follows [23]. It has

been known since Ref. [46] that the DM flux shows annual modulation due to the differences in relative orientations of the DM wind and the direction of the Earth's motion around the Sun. The corresponding effect for AQN-induced axions was computed in Ref. [22]. The daily modulation which is a very specific feature of the AQN model, as explained in the previous section, was also computed in the same paper [22]. The broadband strategy is to separate a large frequency band into a number of smaller frequency bins with the width $\Delta\nu \sim \nu$ according to the axion dispersion relation as discussed above.

The time-dependent signal in each frequency bin $\Delta\nu_i$ has to be fitted according to the expected modulation pattern, daily or annual. For example, the daily modulation should be fitted according to the formula

$$\begin{aligned} \rho_a^{\text{AQN}}(t) &\equiv A_{(d)}(t) \langle \rho_a^{\text{AQN}}(R_\oplus) \rangle \\ A_{(d)}(t) &\equiv [1 + \kappa_{(d)} \cos(\Omega_d t - \phi_0)], \end{aligned} \quad (6)$$

where the density $\langle \rho_a^{\text{AQN}}(R_\oplus) \rangle$ is estimated according to Eq. 3. In this formula $\Omega_d = 2\pi \text{ day}^{-1}$ is the angular frequency of the daily modulation, while ϕ_0 is the phase shift. It can be assumed to be constant on the scale of days. However, it slowly changes during the annual seasons due to the variation of the direction of DM wind with respect to the Earth. The modulation factor $A_{(d)}(t)$ is normalized to unity, $\langle A_{(d)}(t) \rangle = 1$ when averaged over a large number of daily cycles. The same modulation factor $A_{(d)}(t)$ with the same $\kappa_{(d)}$ also enters the estimation for the flux [Eq. ??].

A similar formula holds for the annual modulation defined as:

$$\begin{aligned} \rho_a^{\text{AQN}}(t) &\equiv A_{(a)}(t) \langle \rho_a^{\text{AQN}}(R_\oplus) \rangle \\ A_{(a)}(t) &\equiv [1 + \kappa_{(a)} \cos \Omega_a(t - t_0)], \end{aligned} \quad (7)$$

where $\Omega_a = 2\pi \text{ yr}^{-1}$ is the angular frequency of the annual modulation and label “a” in Ω_a stands for annual. The $\Omega_a t_0$ is the phase shift corresponding to the maximum on June 1 and minimum on December 1 for the standard galactic DM distribution, see Refs. [46, 47].

Several tests can unambiguously discriminate a true genuine signal from spurious signals and noise. These tests will be discussed in detail in Sec. VB after some simplified estimates for the expected signal are performed in Sec. VA.

V. EXPECTED SIGNALS FROM THE AQN-INDUCED AXIONS

The goal of this section is twofold. First, in Sec. VA we want to highlight the basic differences in the conversion rate between the AQN-induced relativistic axions and the conventional galactic axions which normally enter the well-known formulae. Second, in Sec. VB we formulate the key element of this work: despite the very low rate of axion-photon conversion the daily modulation

effect reviewed in the previous Sec. IV effectively allows us to discriminate the genuine signal from the spurious signals.

A. Estimates for the axion to photon conversion

The starting point for our analysis is the computation of the conversion rate which is determined by the well-known formula (e.g. see Ref. [18] for a review):

$$\begin{aligned} P_{a \rightarrow \gamma} &= \frac{1}{v_a} \left(\frac{g\mathcal{B}}{q} \right)^2 \sin^2 \left(\frac{qL}{2} \right), \\ q &= \omega - \sqrt{\omega^2 - m_a^2} \end{aligned} \quad (8)$$

where L is a typical distance where the magnetic field \mathcal{B} is present and $g \equiv g_{a\gamma\gamma}$ is defined as usual

$$g \simeq \frac{g_\gamma \alpha}{f_a \pi}, \quad g_\gamma = \frac{1}{2} \left(\frac{N_e}{N} - \frac{5}{3} - \frac{m_d - m_u}{m_u + m_d} \right) \quad (9)$$

where the parameter $N_e/N = 0$ for KSVZ model, and $N_e/N = 8/3$ for DFSZ model such that $g_\gamma \approx -0.97$ for the Kim-Shifman-Vainshtein-Zakharov (KSVZ) [4, 5] model and $g_\gamma \approx 0.36$ for the Dine-Fischler-Srednicki-Zhitnitsky (DFSZ) [6, 7] model. The parameters f_a and m_a are not independent, but tightly linked in conventional QCD axion models:

$$m_a \approx 6 \cdot 10^{-4} \left(\frac{10^{10} \text{ GeV}}{f_a} \right) \text{ eV}. \quad (10)$$

In case of the resonant conversion when combination $(qL) \ll 1$ could be very small the conversion rate assumes the usual form $P_{a \rightarrow \gamma} = v_a^{-1} (g\mathcal{B}L/2)^2$. The dependence $P_{a \rightarrow \gamma} \propto L^2$ should be interpreted as the presence of a coherent phenomenon when the amplitude of the transition is linearly proportional to the coherence length $\propto L$, while the probability $P_{a \rightarrow \gamma}$ shows the quadratic dependence $P_{a \rightarrow \gamma} \propto L^2$ as stated above.

For nonrelativistic axions, which are normally assumed to be the case for haloscope studies, we can approximate $q \simeq m_a$. If Eq. 8 is multiplied with the axion flux $(\rho_a v_a)/m_a$ we arrive at the following expression for the flux Φ_γ which counts the number of microwave photons emerging from the magnetic field region per unit time and unit area:

$$\Phi_\gamma = \frac{(\rho_a v_a)}{m_a} P_{a \rightarrow \gamma} \approx \frac{\rho_a}{m_a} \left[\frac{g\mathcal{B}}{m_a} \sin \left(\frac{m_a L}{2} \right) \right]^2. \quad (11)$$

This expression agrees with the formula of ref. [48] which was derived using a very different approach. It is important to note that Eq. 11 does not depend on the axion's velocity v_a in the nonrelativistic limit $v_a \rightarrow 0$. This is because Eq. 8 for the conversion is proportional to v_a^{-1} while computation of the Φ_γ requires multiplication of the conversion $P_{a \rightarrow \gamma}$ to the axion flux which itself is proportional to v_a such that the dependence on v_a cancels out at small $v_a/c \ll 1$.

We now turn to our relativistic case for the AQN-induced axions. In this case, the axion density and the flux are determined by Eq. 3 and Eq. 4 correspondingly. Furthermore, the time formation of a photon is very short $\sim m_a^{-1}$. The corresponding length scale is determined by the de Broglie length $\lambda \sim m_a^{-1}$ which is also much shorter than the size of the system L . Therefore, the conversion process can be treated as it happens in an infinite volume system.

As a result, the flux Φ_γ^{AQN} can be estimated, which counts the number of microwave photons emerging from the magnetic field region per unit time and unit area as a result of a *single* elementary process as:

$$\Phi_\gamma^{\text{AQN}} \approx \Phi_a^{\text{AQN}} P_{a \rightarrow \gamma} \approx \Phi_a^{\text{AQN}} \left[\frac{g\mathcal{B}}{m_a} \sin\left(\frac{m_a L}{2}\right) \right]^2, \quad (12)$$

Here, we neglect numerical factors of order of 1 and assume $\langle E_a \rangle \simeq m_a$, $\langle v_a^{\text{AQN}} \rangle \simeq 1$ for simplicity. We also disregard complications related to the relative orientation of the magnetic field and axion momentum, as well as differences between L_x, L_y, L_z that characterize the magnetic field geometry. All these effects are negligible for our qualitative analysis, given that the detector L in all directions is much larger than the Compton wavelength m_a^{-1} , which determines the typical length scale of an elementary process, ensuring $(m_a L) \gg 1$. This should be contrasted with the canonical treatment of the problem when the de Broglie wavelength is very large, i.e. $\lambda \sim (m_a v_a)^{-1} \gg L$, in which case the geometrical properties of the cavity become important and a single cavity mode gets excited. In contrast, in the relativistic case, many modes get excited simultaneously with approximately the same quality factor of order 1. Therefore, higher-order modes, if excited by AQNs, can contribute to the total received power. However, in practice, the coherent addition of the power of the higher-order modes must be carefully accounted for. This complexity underscores the challenges in detecting broadband signals in resonant cavities. Further comments on this issue are provided in item 5 of Sec. VII.

In principle, we should integrate over the axion spectrum computed previously with a specific geometry of the instrument when E_a approximately varies in the range $E_a \in (m_a - 1.8m_a)$. The non-resonant Primakoff conversion generates a complicated photon's broadband spectrum with $\nu \in (m_a - 1.8m_a)$. However, for our qualitative estimates when we want to study the parametrical dependence of the rate as a function of the key parameters such as the axion mass m_a , magnetic field \mathcal{B} , the basic size of the detector L , all these numerical factors of the order of one are irrelevant and will be ignored in what follows.

Several key elements lead to the different treatments of the nonrelativistic galactic axions in comparison with relativistic AQN-induced axions. First of all, the conventional treatment of the problem assumes the coherent conversion such that $P_{a \rightarrow \gamma} \propto L^2$ as explained above.

This L^2 scaling holds as long as L is smaller than the De Broglie wavelength λ where conversion effectively occurs. If coherence persists up to distance $l < L$ the enhancement is much smaller, $L^2 \rightarrow Ll$ [18]. It should be contrasted with the AQN induced axions when $\Delta\nu/\nu \sim 1$ and one should average over many cycles such that $\sin^2(\frac{qL}{2}) \rightarrow 1/2$ in Eq. 12.

The second important element is as follows. Equation 12 describes the conversion as a result of a *single* elementary process when a typical formation length is of order the Compton length m_a^{-1} . One should sum over the number of effective layers along the axion's path which is of order $(Lm_a) \gg 1$. This factor represents a strong enhancement factor. It is noted that this factor $(Lm_a) \gg 1$ agrees with formula Ll presented in Ref. [18] when the coherence length l should be replaced to the Compton length, i.e., $l \rightarrow m_a^{-1}$.

Therefore, based on the above, we arrive at the following estimate for the total flux Φ_γ^{AQN} which counts the number of microwave photons emerging from the magnetic field region per unit time and unit area:

$$\Phi_\gamma^{\text{AQN}}(\text{tot.}) \approx \frac{A}{2} \Phi_a^{\text{AQN}} \cdot \left[\frac{g\mathcal{B}}{m_a} \right]^2 \cdot (m_a L). \quad (13)$$

Numerically, this estimate can be presented as:²

$$\Phi_\gamma^{\text{AQN}}(\text{tot.}) \approx 0.1 A \left(\frac{g_\gamma}{0.97} \right)^2 \left(\frac{\mathcal{B}}{10T} \right)^2 \left(\frac{L}{m} \right) \frac{\text{ph.}}{\text{m}^2 \text{day}}, \quad (14)$$

which implies that a detector with area $\sim 10 \text{ m}^2$ could see a few microwave photons a day. This rate, of course, is quite low. However, recent technological advances suggest that such photons can be, in principle, detected using such instruments as Broadband Reflector Experiment for Axion Detection (BREAD). We refer to Ref. [49] on broadband technology in axion searches.

Another point here is that the average rate is low, however, no detector tuning is required in contrast with conventional resonant-type experiments (see, for example, Ref. [50] for a review). Furthermore, as we already mentioned, the AQN-induced axion DM flux [Eq. 4] does not depend on the axion mass m_a in contrast with the conventional axion production mechanisms when $\rho_{\text{DM}} \sim m_a^{-7/6}$. As a result, the rate in Eq. 14 also does not depend on m_a and this holds as long as our assumption on relatively large mass is satisfied, i.e. $(m_a L) \gg 1$.

The conversion rate in Eq. 8 allows us to estimate the excess of power due to the emitted microwave photons. First, we have to multiply the axion density [Eq. 3] by the conversion rate in Eq. 8. Second, an average over many cycles should be performed such that $\sin^2(\frac{qL}{2}) \rightarrow 1/2$ and insert the enhancement factor $(Lm_a) \gg 1$ as explained

² We use the conversion factor for the magnetic field when Gauss $\approx 1.95 \cdot 10^{-2} \text{ eV}^2$ in Heaviside-Lorentz units when $\alpha^2 = e^2/(4\pi)$, see e.g. Appendix A in [18].

after Eq. 12. Third, to estimate the released energy per unit of time, the obtained expression should be divided to the typical formation time of the elementary process, $\Delta t \sim m_a^{-1}$. Collecting all these factors together we arrive at the following estimate for the power excess $P_\gamma^{\text{AQN}}(\text{tot.})$ due to the emitted microwave photons in unit volume:

$$P_\gamma^{\text{AQN}}(\text{tot.}) \approx A(t) \left(\frac{\langle \rho_a^{\text{AQN}} \rangle m_a}{2} \right) \left(\frac{g\mathcal{B}}{m_a} \right)^2 (m_a L). \quad (15)$$

This estimate can also be represented in a conventional way as an excess of power in unit volume measured in W/m^3 :

$$P_\gamma^{\text{AQN}}(\text{tot}) \approx 10^{-26} A(t) \left(\frac{\text{Watt}}{\text{m}^3} \right) \left(\frac{g_\gamma}{0.97} \right)^2 \times \left(\frac{\mathcal{B}}{10 \text{ Tesla}} \right)^2 \left(\frac{L}{m} \right) \left(\frac{m_a}{10^{-4} \text{eV}} \right)^2. \quad (16)$$

The theoretical estimate (16) does not include additional factors that may considerably modify the numerical value of 10^{-26} . First of all, as we already mentioned we do not distinguish the geometrical parameters L_x, L_y, L_z which we assume are similar in values. Another factor that may potentially modify a numerical value of (16) is the different sensitivity of a cavity to specific frequency bands, the “filters”. We do not include such factors in the estimate (16) as these modifications could be very different for different instruments, while (16) represents a very generic parametrical behavior for an excess of power as a function of m_a describing the conversion of the relativistic axions. In particular, it includes the most important features of the system such as the typical formation time of the elementary process $\sim m_a^{-1}$ and the presence of a large number of layers along the path of relativistic axions. It is noted, that the value of $m_a \sim 10^{-4} \text{eV}$ serves as a benchmark value for all our numerical estimates in the present work as mentioned at the end of Sec. III. As long as the condition $(m_a L) \gg 1$ is satisfied, our assumptions remain valid, and our parametric dependence on m_a holds.

It is instructive to compare this estimate with the conventional computations for the power excess P_a in cavity resonant-designed detectors when $P_a \propto 10^{-26} Q_\alpha (W/m^3)$, see [18] for a review. Numerically, the value presented in [18] is similar to our estimate (Eq. 16) though the conversion rates and the magnitudes for the axion densities of the AQN-induced and the galactic axions are very different. This numerical similarity will be used in what follows for illustrative purposes in our comparisons of the broadband strategy with conventional resonant type searches.

The key element in resonant searches is the so-called quality factor $Q_\alpha \sim 10^6$, which increases the likelihood of the detection of the axion signal. Indeed, the power excess P_a in the conventional analysis is normally compared with the noise power P_{noise} such that the signal-to-noise

ratio is determined by the formula

$$\frac{S}{N} \approx \frac{\sqrt{N} P_a}{P_{\text{noise}}} \approx \frac{P_a}{T} \sqrt{\frac{t}{\Delta\nu}}, \quad P_{\text{noise}} = T \Delta\nu, \quad N \approx t \Delta\nu$$

where T is the total system noise temperature, t is the integration time and the $\Delta\nu$ is the bandwidth of the axion signal. The “quality factor” $Q_\alpha \sim 10^6$ plays a crucial role in this approach. The basic idea here is to detect the narrow resonance line with $\Delta\nu \simeq 10^{-6} \nu$ with sufficiently long integration time by scanning each narrow frequency band to increase the signal-to-noise ratio.

However, the formula for the power excess $P_\gamma^{\text{AQN}}(\text{tot.})$, as given by Eq. 16, does not carry the enhancement factor $Q_\alpha \sim 10^6$ for the case of non-resonant conversion of the relativistic axions with $\Delta\nu \sim \nu$. Therefore, the entire logic of collecting the signal and discriminating it from the noise is very different from the conventional treatment of the cavity-type experiments briefly reviewed above. The key point for our studies is that there is a well-defined modulation factor $A(t)$ which allows us to discriminate the signal from the spurious events and noise by collecting the signal over entire frequency bands instead of scanning in a single narrow frequency band in cavity-type experiments. In particular, a daily modulation is expected, as discussed in the previous section.

The absence of the quality factor $Q_\alpha \sim 10^6$ in Eq. 16 implies that the broadband search is not particularly sensitive to the axion mass, except for a basic normalization factor $\sim m_a^2$ resulted from a large number of “layers” along the path of the relativistic axion, as explained above. This should be contrasted with conventional resonant type searches when the cavity is designed to be highly sensitive to a very specific narrow frequency band expressed in terms of Q_α . This feature (absence of the resonant sensitivity to m_a) is of course a direct manifestation of the non-resonant axion-photon conversion in the magnetic field when the resonant features of the cavity do not manifest themselves in the expression of the power excess as given by Eq. 16.

Therefore, the only option we are aware of to determine the axion mass m_a in the broadband setting is to determine the frequency of a single photon recorded by a non-resonant axion photon converter. As mentioned after Eq. 14 such technologies that could detect a single photon in this frequency band may indeed become available soon.

We would like to illustrate the arguments presented above with the following numerical estimates. In conventional resonance-type searches the factor $N \approx t \Delta\nu \approx 10^5$ for $\Delta\nu \simeq 10^{-6} \nu$ and $\nu \sim 1 \text{ GHz}$ such that $\delta P_{\text{noise}} \approx P_{\text{noise}} / \sqrt{N} \approx 0.3 \cdot 10^{-2} P_{\text{noise}}$ for $t \approx 1 \text{ min}$ which is a typical integration time in resonant searches.

It should be contrasted with broadband strategy when a signal is collected over the entire frequency band $\Delta\nu \sim \nu \sim 1 \text{ GHz}$ during the entire year (separately for each hour during the day). In this case the integration time $t \approx 10^6 \text{ s}$ and $N \approx t \Delta\nu \approx 10^{15}$ which is more than

sufficient to overcome the absence of the quality factor Q in the broadband Eq. 13, Eq. 15, and Eq. 16.

Furthermore, the estimates in Eq. 13 and Eq. 15 explicitly show the presence of enhancement factor $(Lm_a) \gg 1$ which counts the number of layers along the relativistic axion's path, which is a very distinct element in comparison with conventional searches for nonrelativistic axions. This feature implies that if an instrument is asymmetric, e.g. $L_x \gg L_y$ it automatically becomes sensitive to the direction of the axion's momentum, i.e. the instrument potentially becomes directionally sensitive to the AQN-induced axion flux. Therefore, potentially such an instrument allows us to study the DM distribution locally, in our galaxy beyond the conventional SHM paradigm. A hope is that these novel elements in data analysis can overcome the absence of large numerical factor $Q_\alpha \sim 10^6$ which is normally present in the cavity resonant type studies.

B. Reasoning behind the expected daily modulation

We now focus on the central aspect of this proposal: the anticipated daily modulation of AQN-induced axions, as described by Eq. 6. This modulation suggests that the count of emitted microwave photons (Eq. 14) or a power excess (Eq. 16) can accumulate over extended periods, potentially years. To identify a signal, we analyze its presence by fitting the daily modulation pattern given by Eq. 6, summing data across all days of the entire season at a fixed time each day. By examining power excess as a function of time over a 24-hour cycle, any axion-related observables—such as induced currents or voltages in the pickup loop (as reviewed in Sec. IV)—should display this same daily modulation, regardless of the specific recording method used, since the expression for $A_{(d)}(t)$ in Eq. 6 remains consistent across all observables.

While the emission rate (or power excess) is relatively low, the accumulation of events over several years could yield a statistically significant detection of daily modulations if present in the data. More importantly, this approach can effectively differentiate a genuine signal from spurious signals and noise.

A straightforward method to validate the signal is to observe whether the phase ϕ_0 in the fitting function of Eq. 6 remains constant throughout the year. Instead, within the AQN model, ϕ_0 should vary slowly with the annual cycle. Specifically, the phase should reach approximately $\phi_0 \simeq \pi$ in opposite seasons (e.g., summer vs. winter), indicating a phase shift where the daily maximum and minimum values invert. This test, which involves comparing frequency patterns in daily and annual modulation data, requires a long observation period (spanning years) but is robust, as systematic errors differ markedly between daily and annual cycles.

Another simple test involves analyzing datasets with $\mathcal{B} = 0$ T. The results for $\mathcal{B} = 0$ should differ significantly

from those obtained with $\mathcal{B} \neq 0$ T, helping to distinguish genuine signals from noise.

Daily modulations can also be analyzed through Fourier transformation, where a single, well-defined peak at $\Omega_d = 2\pi \text{ day}^{-1}$ would indicate a genuine modulation, with smaller, randomly distributed peaks representing noise and other spurious effects. Performing the Fourier transform on $\mathcal{B} = 0$ data, in contrast, should yield a more random noise spectrum rather than a clear peak at $\Omega_d = 2\pi \text{ day}^{-1}$.

An even stronger test for ruling out spurious signals involves using a network of synchronized instruments to detect correlated signals, as discussed in original studies [23, 43].

Various systematic background fluctuations—such as temperature changes, tidal effects, atmospheric phenomena, and human activity—might mimic the daily modulations described here. These spurious signals could be distinguished from genuine axion signals by analyzing the $\mathcal{B} = 0$ data and observing phase shifts in ϕ_0 over different seasons. It would be challenging for spurious signals to replicate all the unique features associated with AQN-induced axions described above.

These considerations suggest that cavity-type experiments—currently the only instruments probing the parameter space for standard QCD axion models with $10^{-6} \text{ eV} \lesssim m_a \lesssim 10^{-3} \text{ eV}$ could be used to investigate the daily modulation pattern. These instruments function as non-resonant axion-photon converters within a strong magnetic field.

In principle, any current axion detection instrument could analyze daily modulation following this method, including all previously collected data. This approach allows haloscopes to search for daily modulation by aggregating data (e.g., within a specific month where ϕ_0 can be assumed constant) for each hour, enabling detection of potential daily modulation patterns.

In this work, we apply this methodology to data collected by the CAST-CAPP detector from September 2019 to June 2021, presenting it as a proof-of-principle demonstration.

VI. DATA ANALYSIS

CAST-CAPP is an axion haloscope that operated from September 2019 to June 2021 at CERN in Geneva, Switzerland. It consists of four tuneable $23 \times 25 \times 390 \text{ mm}$ rectangular cavities, each with a volume of $V = 224 \text{ cm}^3$ which were placed inside one of the two bores of CAST's superconducting dipole magnet of 8.8 T. CAST-CAPP introduced two novel mechanisms to increase its sensitivity to DM axions but also to transient events:

- The phase-matching technique combines the power outputs from the four identical cavities coherently to increase the Signal-to-noise ratio (SNR) linearly with the number of cavities.

- The state-of-the-art fast-tuning mechanism with a speed of 10 MHz/min and a resolution of 100 Hz, together with wideband electronics allow for a search for transient events such as axion streams and axion miniclusters.

CAST-CAPP acquired 172 days of data using both single and phase-matched cavities in data-taking conditions. The total covered frequency range extended from 4.77 GHz to 5.43 GHz covering a parameter phase space of ~ 660 MHz corresponding to axion masses between $19.74 \mu\text{eV}$ to $22.47 \mu\text{eV}$. This allowed CAST-CAPP to exclude axion-photon couplings for virialized galactic axions down to $g_{a\gamma\gamma} = 8 \times 10^{-14} \text{ GeV}^{-1}$ at 90% confidence level [51].

The design and the data-taking scheme of the CAST-CAPP experiment were not originally optimized for a search of daily modulated signals. However, the large acquired bandwidth of 5 MHz compared to the resonance width of about 200 kHz, together with the wide frequency range of 660 MHz, allows a first proof of principle demonstration of such a novel analysis procedure.

Finally, it is mentioned that for the probed axion mass range $19.74 \mu\text{eV}$ to $22.47 \mu\text{eV}$, the corresponding Compton wavelength $\lambda_C = h/m_a c$ ranges from approximately 6.3 cm to 5.5 cm. Given the length of each single rectangular cavity $L = 39$ cm, the condition $m_a L$ yields:

$$m_a L \approx \frac{1}{\lambda_C} \cdot L \approx 6.2 \text{ to } 7.1. \quad (17)$$

Thus, the condition $m_a L \gg 1$ being used in our estimates is marginally satisfied for the entire frequency range of CAST-CAPP, indicating that the cavity length is sufficiently long to justify the requirement for relativistic axion detection.

A. Simulation

The AQN signal search is simulated using artificially modulated signals incorporated into simulated spectra. These spectra are composed of randomly generated noise superimposed onto authentic spectral baselines acquired from CAST-CAPP. The CAST-CAPP data are recorded in 1 min bins which are then converted into frequency-domain using a fast Fourier transformation. Each spectra carry a unique shape originating from the TE_{101} resonance mode of the cavities and the amplifier gain profile. To preserve this unique shape in the simulated spectra, we produce spectral baselines using a low-pass filter. Each simulated spectrum is comprised of artificial Nyquist noise multiplied by a randomly selected spectral baseline. The artificial noise has a Gaussian distribution of $\mu = 1$ and $\sigma = 1/\sqrt{\text{resolution bandwidth} \cdot \text{spectrum time length}} = 0.018$. Then, the artificial daily modulated signal is simulated using a sine function of tuned amplitude added on top of the artificial noise spectrum.

We performed multiple simulations with the data-taking time ranging from 1 day to 20 days at a single frequency and an AQN signal amplitude ranging from 0.01% to 3.5% of the mean spectral power. Figure 2 shows the combinations of simulation parameters i.e. number of data taking days and AQN amplitude percentage with respect to the mean spectral power on a two-dimensional (2D) space, while the color denotes if that specific combination yielded a significant signal that can be detected on a Lomb Scargle periodogram. The confidence level for the significance is set to be 90%. In Fig. 3, a periodogram for a significant AQN signal is shown for 5 days of data taking. In this case, the power of the AQN signal is adjusted to be 0.8% of the mean spectrum power.

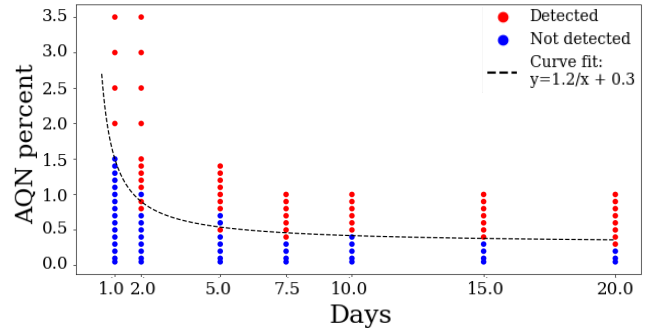


FIG. 2. Simulation of AQN signals with different power percentages as a function of data taking time. The red points indicate discoveries while the blue ones denote insignificant signals. A decision boundary fit is denoted by the black dashed line.

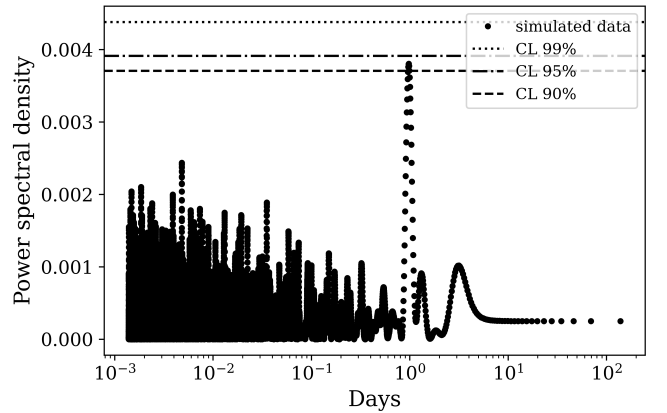


FIG. 3. Periodogram created using the Lomb Scargle method with 5 days of simulated data taking time (5 days / 1 min = 7200 simulated spectra). The observed signal is above 90% CL.

B. Data selection

During its ~ 2 years of operation, CAST-CAPP recorded 64.3 days of data using phase-matched cavities and 107.5 days of data using a single cavity. At the same time, a total of 16.4 days of data were recorded with $B = 0$ T. The total frequency range of CAST-CAPP is 644 MHz between 4.79 MHz to 5.43 MHz.

The strategy for quick identification of parasitic environmental background electromagnetic signals involved the usage of an external omnidirectional antenna that was placed outside the CAST magnet and was connected to a spectrum analyzer recording data simultaneously with the cavities. This Electromagnetic Interference / Electromagnetic Compatibility (EMI/EMC) identification system was introduced in November 2020 and collected 89.2 d of data.

For the search for a daily modulation with the CAST-CAPP data, only the data acquired from single cavities are used. The reason is that for the phase-matching procedure, a series of programmable phase-independent attenuators were used to adjust the amplitude of the resonance of the individual cavities with 0.25 dB to match the resonance peaks of the phase-matched cavities. This introduced additional attenuation although it is not expected to follow a daily rhythm, it could introduce unnecessary artifacts in the whole procedure.

Additionally, the data that were recorded during the usage of the fast-frequency scanning search approach are omitted from the analysis due to the inherent frequency dependency of the amplitude of the resonance peak of the TE_{101} mode.

Therefore, as a first approximation only data taken from single cavities at fixed frequencies are used. Because of the high mechanical Q -factor of the sapphire strips of the tuning mechanisms, which were accidentally acting as a mechanical tuning fork, a frequency modulation of the electromagnetic mode of interest was observed in certain frequencies. The data that were dominated by such mechanical vibrations as well as other undesired non-systematic effects or possible EMI/EMC interferences were excluded from further consideration using a series of selection criteria.

C. Data treatment

Each CAST-CAPP spectrum that fulfills the aforementioned cuts is averaged within itself to get the mean spectral power, then a timestamp is attached to it. This way, a 2D array of mean spectral powers vs corresponding acquisition timestamps is created. This array is then rebinned hourly to have mean spectral powers for each hour of data acquisition.

The baseline on which the daily modulation in the CAST-CAPP data are investigated is expected to fluctuate due to weekly environmental temperature changes as well as amplifier gain drifts. To remove this effect before

the daily modulation analysis, the data are detrended. For this purpose, a Savitzky Golay filter is used with a window length of two days and a polynomial order zero. Figure 4 shows an example of the mean spectral power time-series belonging to cavity 3 between August 13, 2020 and September 20, 2020, with $B = 8.8$ T before and after the detrending procedure. The parameters of the low-pass filter are selected such that the long-range fluctuations are removed while the investigated daily variations are conserved. After detrending, the mean spectral variations are around a baseline that equals 1.

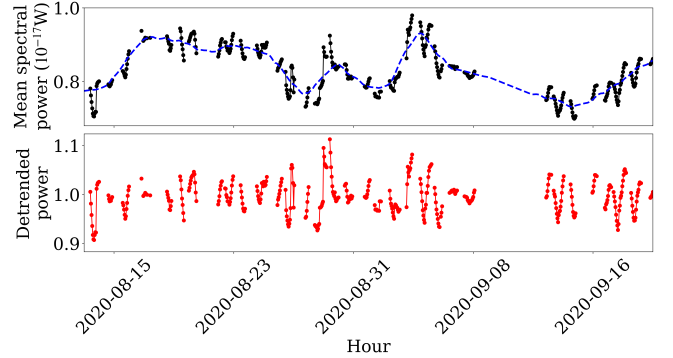


FIG. 4. Mean hourly spectral power (black) of cavity 3 between August 13 and September 20. The blue dashed curve denotes the trend of the data created by the low-pass filter. The detrended power (red) is calculated by dividing data (black) by its trend (blue).

D. Systematics

CAST-CAPP's experimental setup consists of a three-stage signal amplification using High-electron-mobility transistor (HEMT) amplifiers. The first amplification for each cavity was performed close to the critically coupled signal output port at the 4 K stage using a HEMT low-noise amplifier (LNA) from Low Noise Factory (LNF-LNC4_8D) providing about 39 dB amplification with 2 K noise for the frequency of CAST-CAPP at data-taking conditions. Then, an external amplification of about 22 dB for each of the four cavities is performed using four room-temperature ZX60-83LN-S+ LNAs. Finally, a third-stage amplification is performed at the output signal after the signal combination using a Miteq AFD3-0208-40-ST LNA providing about 30 dB gain.

Two main parameters that could influence the performance of these amplifiers are the provided bias voltage as well as the temperature. Regarding the bias voltage of the room-temperature amplifiers, this was provided using a Hameg HM7044 power supply. The accuracy of this power supply is on the level of 10 mV which translates into a tiny gain variation of about 0.007 dB or $\sim 0.2\%$ in terms of power for the second- and third-stage amplification. For the first-stage amplification, a LNF-

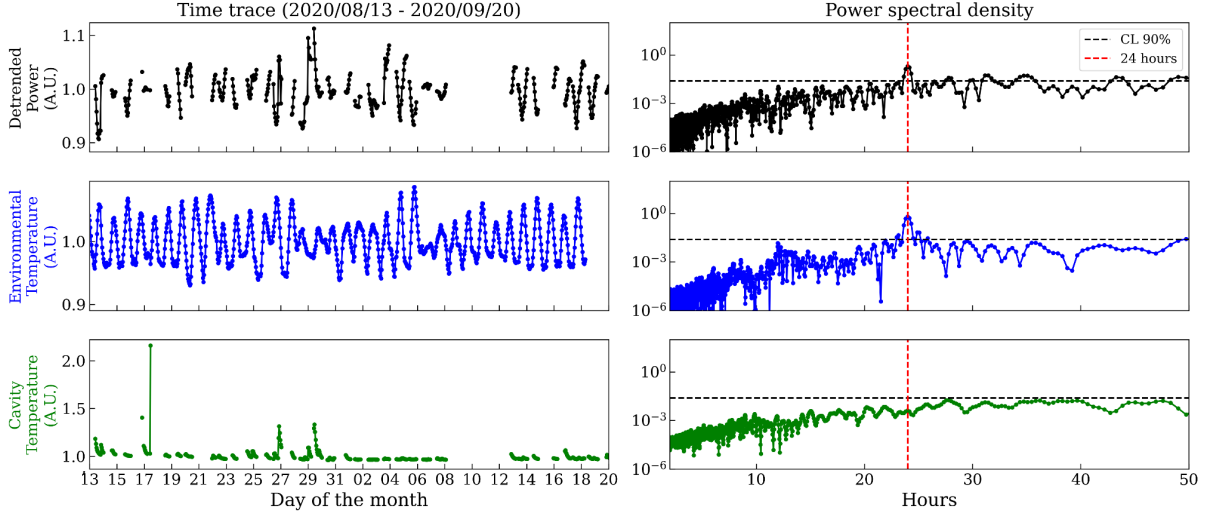


FIG. 5. Time series (left) and Lomb Scargle periodograms (right) for data taken between August 13, 2020 and September 20, 2020, with cavity 3 and $B = 8.8$ T. The results for the ambient environmental temperature as well as the cavity temperature are also shown.

PS.3 was used for the biasing (0.25 V and 7 mA) for each one of the four LNAs corresponding to the four cavities. The biasing conditions of the LNAs were optimized for minimal heat dissipation and minimum noise figure. A current variation of about 0.02 mA was observed which would translate into a variation of the gain on the level of 0.02 dB or $\sim 0.5\%$ in terms of power.

The temperature of the first stage LNA was monitored through Cernox temperature sensors to be varied between 6 K to 10 K. For the specific bias of about 0.25 V and 7 mA at a frequency of ~ 5 GHz the gain of LNF-LNC4.8D has a linear dependency on temperature, increasing the gain of about 0.25 dB and an increase of their noise temperature 0.42 K for the aforementioned temperature range [52]. This translates into a $\sim 5.9\%$ variation of the gain in terms of power.

The temperature and humidity of the experimental area of CAST were monitored and kept stable via a central ventilation and temperature control system. However, small changes of a few degrees were observed. More specifically, the temperature variation observed in the room-temperature electronics of CAST-CAPP during the whole data-acquisition period was measured using PT100 Resistance Temperature Detector (RTD) temperature sensors to be between 19°C to 26°C . These small temperature changes result in a linear gain dependency of the room-temperature amplifiers. For the second stage amplification a temperature change between 19°C to 26°C results in a decrease of the gain of about 0.078 dB or about 1.8% decrease in power [53]. Then, the third stage amplifier is expected also to have a decrease of its gain with increasing temperature of about 0.02 dB/ $^\circ\text{C}$ [54]. Therefore, for 19°C to 26°C , the third-stage room-temperature amplifier would exhibit a decrease of ~ 0.14 dB or about 3.3% decrease in terms of

power.

It is noted that the gain stability of the amplifiers over time is expected to have a much smaller variation than the changes in temperature, voltage, or other ambient conditions. An overview of the aforementioned systematics and their effect on the gain of the various amplifiers of CAST-CAPP is shown in Table II.

	First stage	Second stage	Third stage
Temperature	+5.9%	-1.8%	-3.3%
Bias	-0.5%	+0.2%	+0.2%

TABLE II. Systematic effects on the variation of power from the various amplification stages of CAST-CAPP. The \pm signs indicate a positive or a negative correlation with the corresponding parameter.

Finally, the magnetic field of CAST's dipole magnet was measured to have a stability of 0.1 mT at the level of 8.8 T, translating into an insignificant 0.001% variation.

E. Results

1. Lomb Scargle periodogram

The first step for a search for daily periodicities is performed using a Lomb-Scargle periodogram in the detrended data. This way periodicities around 24 h can be investigated. Figure 5 shows the time series (top left) as well as the periodogram (top right) for $B = 8.8$ T, where the time period from August 13, 2020 to September 20, 2020 is selected for single cavity 3. At the same time, using the same period, the environmental temperature from the PT100 RTD temperature sensors (Fig. 5 middle row) as well as the cavity temperature from the Cernox

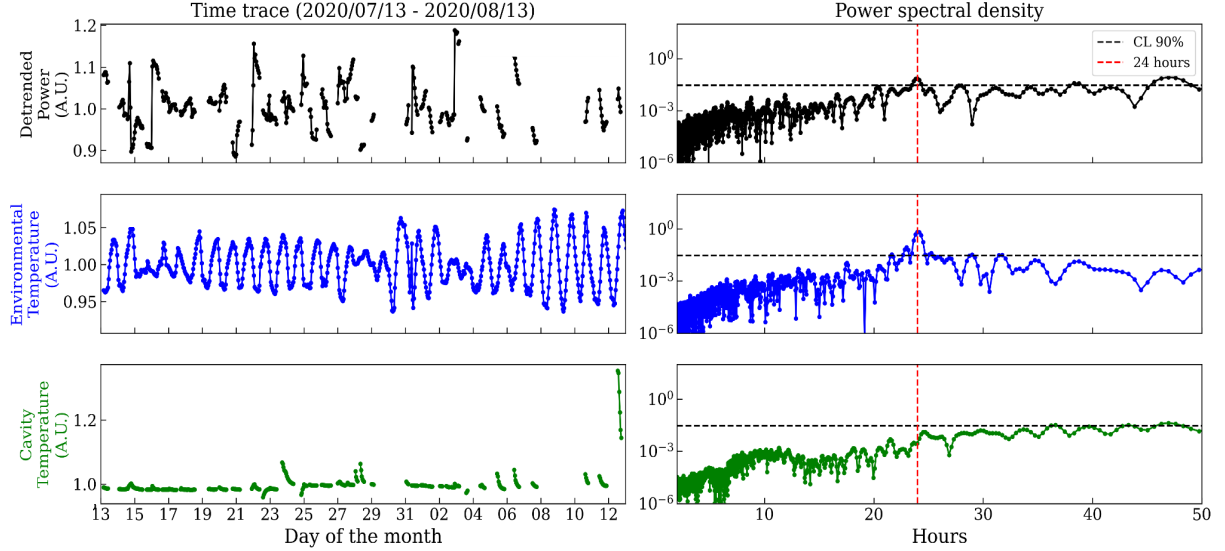


FIG. 6. Time series (left) and Lomb Scargle periodograms (right) for data taken between July 13, 2020 and August 13, 2020 with cavity 2 and $B = 0$ T. The results for the ambient environmental temperature as well as the cavity temperature are also shown.

sensors (Fig. 5 bottom row) are shown. Table III gives a detailed account of the data statistics. The observed gaps in the time series for both the cavity power and the cavity temperature data are due to CAST-CAPP's data acquisition scheme not being optimized for a daily modulation search, i.e. CAST-CAPP acquired single-cavity data for a limited and noncontinuous amount of time. As the time series data were detrended and normalized, the y axes are given in arbitrary units. In the periodograms (Fig. 5 right) that show the power spectral densities of the corresponding time series, the red dashed vertical line denotes the daily periodicity whereas the black dashed horizontal line shows the threshold for a significant signal with a confidence level of 90%.

As expected the environmental temperature follows a daily rhythm while the cavity, and therefore the cryogenic LNA temperature, does not. Interestingly, also the detrended cavity data for $B = 8.8$ T show a 24 h periodicity above 90% confidence level. However, the ambient temperature change can be the leading cause behind the observed daily variation due to its direct effect on the receiver chain, as mentioned in the previous section. This is due to the room-temperature setup not being sufficiently isolated from temperature changes.

For comparison reasons, also, $B = 0$ T data are analyzed. For cavity 2, these are shown in Fig. 6 with the selected period ranging from July 13, 2020 to August 13, 2020. For practical reasons, this period could not be the same as the period for the $B = 8.8$ T, and therefore, a perfect comparison between the two cannot be made. The time difference between the two is about 1 month. As can be seen, these data with $B = 0$ T also indicate a daily modulation being less pronounced than the $B = 8.8$ T data, which may be attributed to the

$B = 0$ T data being more sparse compared to $B = 8.8$ T data. Therefore, we conclude that the appearance of similar daily modulation in the detrended power spectra of Figs. 5 ($B = 8.8$ T) and 6 ($B = 0$ T), does not allow to exclude systematics for the hint in Fig. 5 but strengthens the assumption of a temperature-dependent amplifier gain effect. Of course, there is always the possible dark photon contribution with no magnetic field dependence.

2. Hourly distributions

Next, the hourly distributions of both $B = 8.8$ T and $B = 0$ T data are shown in Fig. 7 which could provide more information on the observed daily modulation. These distributions are derived for the same time intervals as in Fig. 5 and Fig. 6, and are shown in the form of box plots in Fig. 7a, top and bottom, respectively. A sine function is then fitted to the data. The maxima of the fitted sinus shapes are found around 05:58 and 05:19 local hour, for $B = 8.8$ T and $B = 0$ T respectively.

Even assuming the sine-wave fit of $8.1\% \pm 0.6\%$ and $6.0\% \pm 0.5\%$ for the mean daily distributions of $B = 8.8$ T and $B = 0$ T respectively, the CAST-CAPP data do not allow for a significant conclusion in favor of the AQN-related axion signal.

The hourly temperature distributions corresponding to Fig. 7a are shown in Fig. 7b where we observe that the temperature is maximum at around 18:30 Central European Time (CET). This corresponds approximately to half a day phase shift between the mean spectral power and the ambient temperature variations shown in Fig. 5 and can be explained by the argumentation in Sec. VID.

Table IV gives an overview of the sine fit characteris-

	$\mathcal{B} = 8.8 \text{ T}$ (Fig. 5)	$\mathcal{B} = 0 \text{ T}$ (Fig. 6)
Interval	August 13 - September 20	July 13 - August 13
Data amount (1-min files)	14400	11500
Total hours	240	192
Hours per day	8.0	5.2
Frequency range [GHz]	5.20 - 5.33	5.32 - 5.41

TABLE III. Data statistics for Figs. 5 and 6.

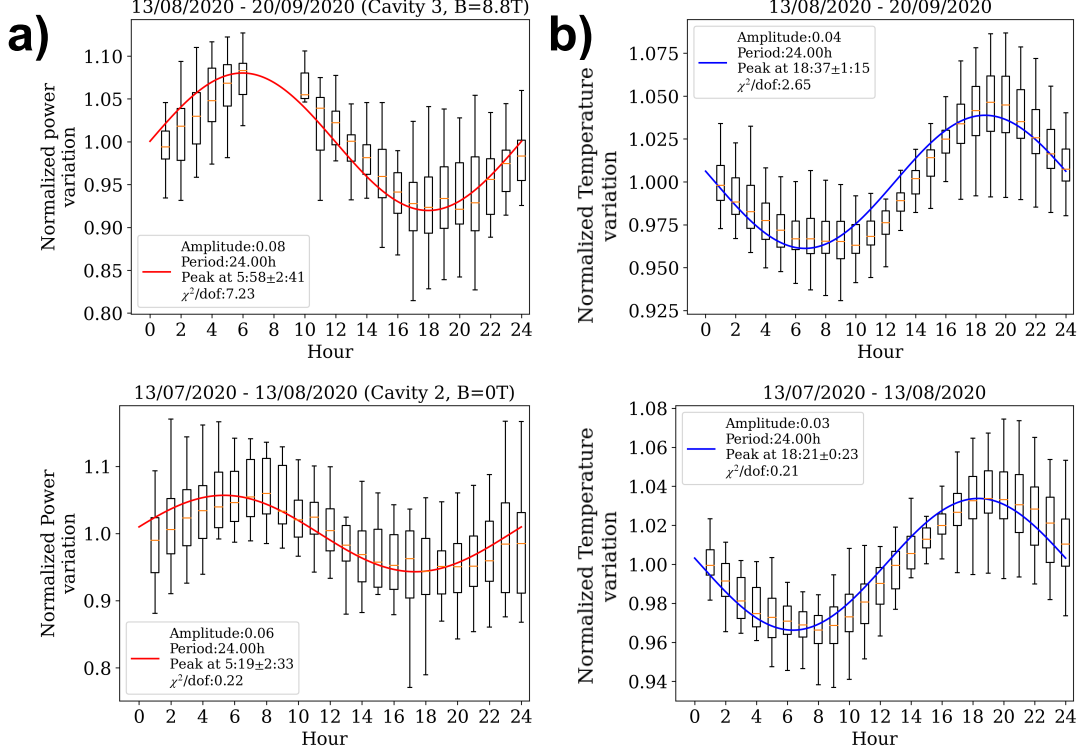


FIG. 7. Box plots showing the hourly distributions of the normalized power variation (left) and normalized environmental temperature variation (right) for $\mathcal{B} = 8.8 \text{ T}$ (top) and $\mathcal{B} = 0 \text{ T}$ (bottom). A sine function with a daily periodicity is fitted to the data. The gap in the upper left plot is due to the CAST-CAPP data acquisition schedule in that time interval of data taken.

tics for the box plots shown in Fig. 7. The sine fits, as shown by chi-squared tests, are significantly more suitable to the data compared to linear fits, confirming the daily periodicity observed. However, it is apparent that the daily change in the environmental temperature of the CAST hall does not allow for a positive conclusion.

		Cavity power	Environmental temperature
B=8.8T	Fluctuation	$8.1\% \pm 0.6\%$	$3.7\% \pm 0.3\%$
	max hour	5:58±3:05 CET	18:37±1:15 CET
	χ^2/dof	7.2(sine), 19.2(line)	2.6
B=0T	Fluctuation	$6.0\% \pm 0.5\%$	$3.4\% \pm 0.3\%$
	max hour	5:19±2:33 CET	18:21±0:23 CET
	χ^2/dof	0.2(sine), 8.3(line)	0.21

TABLE IV. Data statistics for Fig. 7.

Finally, as described in [51], an external quasispherical antenna was placed next to the CAST mag-

net and operated simultaneously with the Spectrum Analyzer (SA) connected to the cavities and at the same frequency band. This was used as a fast veto procedure where a candidate axion signal could be easily discarded if it appeared both in the cavities and in the witness channel at the same frequency. Therefore as an additional comparison, the hourly distribution of the external antenna data is shown in Fig. 8. As expected due to the lack of any temperature isolation, the mean daily spectral variation observed is about 21%, larger compared to both $\mathcal{B} = 8.8 \text{ T}$ and $\mathcal{B} = 0 \text{ T}$ data acquired by a receiver chain that is partially temperature isolated.

3. Seasonal phase shift

Because of the limited amount of data taken with CAST-CAPP, the expected seasonal phase shift (see Sec. VB) has been searched for, but could not be iden-

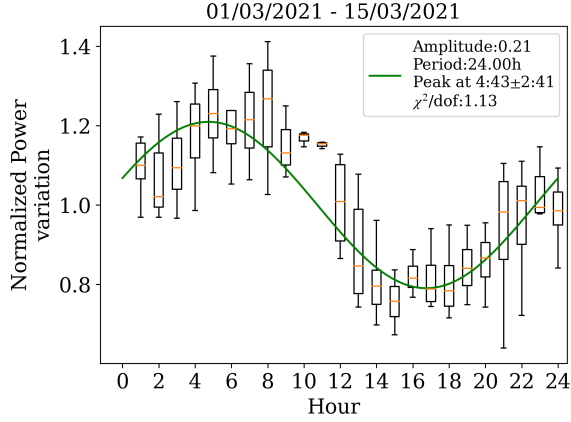


FIG. 8. Box plot of the external antenna spectral power.

tified. The periods August 13, 2020, to September 20, 2020, and November 18, 2020, to December 02, 2020 have been analyzed, where the expected phase shift is only of the order of 0.3π or 3 hours and 40 minutes, however, due to the limited statistics, this number could not be determined.

VII. CONCLUSION AND FUTURE STUDIES

In this work, as a proof of principle study, we have reanalyzed the CAST-CAPP data already published in [51], focusing on the daily modulation of a potential axion signal due to DM axions or from AQNs as it has been presented in Sec. V. We have observed daily variations, but at the same time also for $\mathcal{B} = 0$ T (see Fig. 7 bottom) and with the external antenna (see Fig. 8). Since a similar dependence is expected for the electronic gain of the amplifier chain of the Data Acquisition (DAQ) system, we could not conclusively determine an axion-related signal (whatever its origin: from streams or AQNs), or a signal from the hidden sector. Additional tests to discriminate a true genuine signal from the noise or spurious signals cannot be carried out, as the CAST-CAPP detector is no longer operational. Although the data in Fig. 7 do show some daily oscillation, we conclude that this could be due to temperature effects. However, we can formulate some lessons to be learned, which can be useful for future studies and different instruments, yet to be built:

1. The unavoidable fluctuations of the electronic gain of the amplifier chain, in radiometers, have led to the invention of the Dicke switch to eliminate such gain fluctuations of the entire chain, and not only the front-end like for the noise figure. The inherent switching noise added by this type of radiometer can then be mitigated by the use of the correlation-type radiometer [55, 56]. Virtually every modern microwave remote-sensing satellite is based on this correlation technique, and therefore axion DM ex-

periments can also benefit from this approach and replace the total power radiometer that is currently widely used. This approach can allow for a conclusive search for daily modulations in the future.

2. The expected AQN-related phase shift during different seasons can also serve as an additional check to distinguish a genuine AQN signal from purely environmental or terrestrial influences.
3. An additional unequivocal verification of the existence of an external AQN-related modulation can be derived from the search for sidereal periodicities. These periodicities are anchored to a fixed referenced frame relative to remote stars rather than the Sun thus indicating an exosolar modulation source. The sidereal day of 23 h 56 min and the Moon's sidereal month of 27.32 days are characteristic examples of such periodicities. With high-quality data with enough statistics, this sidereal modulation can become statistically resolvable, thus linking the potential signal to cosmic rather than local environmental sources.
4. To determine the mass and other parameters of the AQN-originated axion a truly broadband instrument should be used. This should be contrasted with the canonical resonant type of experiments when the positive observation of a signal in the form of a narrow peak at ν_0 unambiguously predicts the axion mass $m_a = \nu_0 * (2\pi)$ irrespective of any normalization factors such as the DM density. However, in the case of the nonresonant studies as advocated in this work, the axion mass cannot be extracted with sufficiently high precision as m_a enters only as a normalization factor (16) and is highly sensitive to many uncertainties in the estimates including the general normalization factor such as the local DM density and its possible deviations from the SHM as mentioned above. The local DM density can also vary due to possible sidereal effects as mentioned above.
5. For cavity-type experiments, the sensitivity to higher-frequency photons in broadband detection is limited by the specific coupling and the characteristics of the amplifier chain, whose gain and noise figure are typically constrained to several GHz. Additionally, the bandwidth of the readout system acts as a limiting factor, effectively functioning as a bandpass filter. Future instruments must address these limitations by designing amplifier chains and readout electronics with significantly broader frequency coverage to detect signals from AQN-produced axions.
6. Finally, the best option to discriminate the true genuine signal from the noise or spurious signals is to analyze the correlations in two or more different instruments representing a synchronized network,

as suggested in [23, 43]. The basic idea is to search for some transient events that occur at (almost) the same instant in the different instruments of a network of axion detectors. The typical rate for such correlated events depends on the strength of the spikes as discussed in [23, 43].

If future experiments conclude that the AQN-induced axions are responsible for daily modulation, it would have profound consequences. It would imply that two long-standing puzzles in cosmology, the nature of the DM and the matter-antimatter asymmetry of our Universe are simultaneously resolved. These two puzzles are intimately linked within the AQN framework and an observed signal cannot be explained by any other model since such a signal could be only generated by the relativistic AQN-induced axions (which are mostly localized on the Earth's surface in contrast with CaB³).

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- ³ As we mentioned in Sec. I the local density of the AQN-induced axions on the Earth's surface should be at least 4 orders of magnitude higher than the density of the cosmologically produced CaB axions considered in Ref. [26]; see the paragraph above Eq. 4 for an estimate. This is because the CMB photon density strongly constrains CaB density. The spectral features of the AQN-induced axions are also vastly different from the CaB axions.

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