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Highlights

- A new multi-scale chemo-mechanical model is presented in order to analyze the role of Ca^{2+} in muscle fatigue and to predict fatigue behavior.
- The output results of isometric simulations were fitted with experimental data obtained for rabbit Extensor Digitorum Longus muscle.
- Varying the Ca^{2+} concentration level and its kinetics in the tissue, the model was able to predict the evolution of the active force of an experimental fatigue protocol.

Assessing the role of Ca^{2+} in skeletal muscle fatigue using a multi-scale continuum model.

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Abstract

The Calcium ion Ca^{2+} plays a critical role as an initiator and preserving agent of the cross-bridge cycle in the force generation of skeletal muscle. A new multi-scale chemo-mechanical model is presented in order to analyze the role of Ca^{2+} in muscle fatigue and to predict fatigue behavior. To this end, a cross-bridge kinematic model was incorporated in a continuum based mechanical model, considering a thermodynamic compatible framework. The contractile velocity and the generated active force were directly related to the force-bearing states that were considered for the cross-bridge cycle. In order to determine the values of the model parameters, the output results of an isometric simulation were initially fitted with experimental data obtained for rabbit *Extensor Digitorum Longus* muscle. Furthermore, a simulated force-velocity curve under concentric contractions was compared with reported experimental results. Finally, by

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varying the Ca^{2+} concentration level and its kinetics in the tissue, the model was able to predict the evolution of the active force of an experimental fatigue protocol. The good agreement observed between the simulated results and the experimental outcomes proves the ability of the model to reproduce the fatigue behavior and its applicability for more detailed multidisciplinary investigations related to chemical conditions in muscle performance.

Keywords: Skeletal muscle, Fatigue, Chemo-mechanical model, Calcium

1. Introduction

When skeletal muscles are activated repeatedly with intensity, it is well known that the force output declines. But the force is not the only feature affected by fatigue in this tissue: shortening velocity and relaxation behavior are also affected (Allen et al., 2002; Jones, 2010). Although the central nervous system, motor nerves and the neuromuscular junctions can contribute to this phenomenon, the main mechanisms are located in the muscle itself (Allen et al., 2002). A variety of intracellular processes appear to be responsible for fatigue, and these have been assessed and reported extensively in the literature (Westerblad and Allen, 1991, 1993; Allen and Westerblad, 2001; Allen et al., 2002, 2008). Although all of these experiments provide invaluable insights into these mechanisms, according to Röhrle et al. (2012) computer simulations in conjunction with experimental findings can be a powerful tool for evaluating complex hypotheses and conclusions.

Modeling the behavior of skeletal muscle has typically been focused on sub-cellular processes of a half-sarcomere or on simplified phenomenological relationships to simulate the whole muscle (Röhrle et al., 2012). Regarding the first type of models, the interaction between actin and myosin filaments was initially simulated using a two state model by Huxley (1957). Although this model predicted good results under the mechanical point of view, it was later suggested that the attachment of the cross-bridge occurred at different stages to fulfill heat release outcomes (Huxley, 1973) and transient dynamic responses

23 (Huxley and Simmons, 1971). The complexity of these models was also reduced
24 to achieve more accessible formulation to simulate macroscopic muscle dynam-
25 ics (Zahalak, 1981; Wu et al., 1997). Furthermore, new formulations inspired
26 by contractile processes appeared in order to consider dynamic contractions in
27 a half-sarcomere (Razumova et al., 1999, 2000).

28 Continuous reduction of intracellular Ca^{2+} together with force deficit due
29 to fatigue condition have been proved by experimental investigations on skeletal
30 muscles (Westerblad and Allen, 1991; Allen and Westerblad, 2001). Moreover,
31 the shortening velocity of the tissue is also affected by fatigue, which means a loss
32 of power production and poor performance. These changes in the force-velocity
33 relationship could be caused by Ca^{2+} deficit (Ruiter et al., 2000; Jones, 2010).
34 Since Ca^{2+} is the initiator agent of active behavior in sarcomere, experimental
35 tests suggest a direct relationship between the decay of Ca^{2+} and the functional
36 effects of fatigue (Allen and Westerblad, 2001; Jones, 2010).

37 In this paper, a multi-scale continuum model that fulfills thermodynamic and
38 mechanical requirements is presented to simulate skeletal muscle contraction
39 under fatigue conditions. First, the chemical phase involved in activation at the
40 sarcomere level is described and formulated in Section 2. The thermodynamic
41 basis that allows the derivation of constitutive laws is presented in Section 3
42 and is particularized in Section 4. Finally, the assessment and validation of
43 the model using experimental tests is presented in Section 5. The ability of
44 the model to predict the fatigue response of the tissue is discussed in Section 6
45 together with the formulation proposed and the numerical results obtained.

46 **2. Cross-bridge kinetics**

47 Cross-bridge cycling is responsible for the movement and force production
48 in skeletal muscle cells. In conditions of relaxation, the tropomyosin-troponin
49 complex on the actin filament blocks the actin binding sites and avoids the for-
50 mation of the cross-bridge with the myosin head. When the Ca^{2+} concentration
51 increases above a certain level or resting threshold, Ca^{2+} binds to troponin and

52 tropomyosin exposes the myosin binding sites to cross-bridge formation. Since
 53 cross-bridges within a half-sarcomere are all in a parallel configuration (Razu-
 54 mova et al., 1999), the generated force equals the sum of forces generated by
 55 attached cross-bridges.

56 Although various models of chemical kinetics in cross-bridge cycling have
 57 been suggested (Huxley, 1957; Razumova et al., 2000; Shorten et al., 2007), it
 58 has been accepted that an appropriate model for studying the chemomechanical
 59 force generation mechanism should have at least three contrasting states (Wahr
 60 and Metzger, 1999). The basis for modeling the subcellular level in this work is
 61 the model developed by Razumova et al. (1999), which describes myofilament
 62 regulation and cross-bridge cycling as a four-state process (Fig. 1). The R_{off}
 63 state corresponds to the blocked position of the myosin binding sites by the
 64 tropomyosin-troponin complex. As the Ca^{2+} concentration increases, at the D
 65 state the regulatory unit unbinds the myosin sites but the cross-bridge is not
 66 yet formed. Attachment occurs in A_1 and A_2 states which represent the pre
 67 and post-power stroke, respectively.

68 Transitions through the different states are governed by seven rates k_{on} ,
 69 k_{off} , f , f' , h , h' and g . k_{on} and k_{off} are the rates responsible for the conversion
 70 between the “on” and “off” states of the regulatory unit proportional to the
 71 available Ca^{2+} concentration. The forward power stroke is governed by h and
 72 detachment is governed by g . The reverse reactions are designated with primes
 73 (f' and h') following the same notation as that given by Razumova et al. (1999).

74 In spite of different types of muscle fibers depending on their metabolism, the
 75 behavior of a sarcomere can be generalized to obtain a global muscle response.
 76 Therefore in this model, every state represents a fraction of the cross-bridges or
 77 a probability of being in one of the defined states fulfilling the constrain relation:

$$R_{\text{off}}(t) + D(t) + A_1(t) + A_2(t) = 1 \quad (1)$$

78 Considering this constraint equation and the seven rates, the fraction of
 79 chemical states can be obtained through solving three ordinary differential equa-

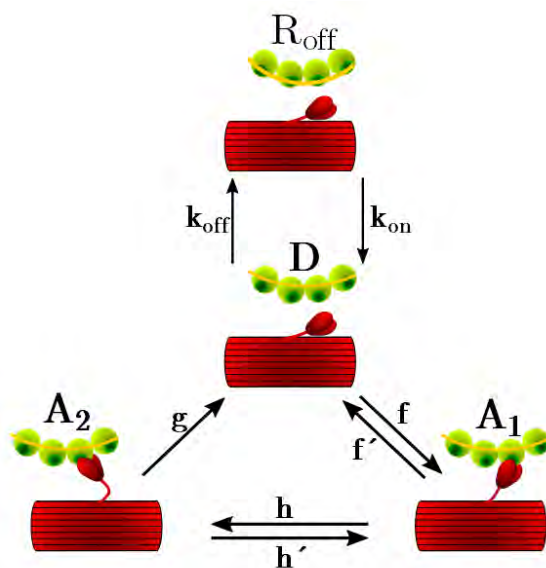


Figure 1: Four-state chemical kinetics (Razumova et al., 1999): R_{off} regulatory unit in “off” situation, D regulatory unit in “on” but detached situation, A_1 attached pre-power stroke and A_2 attached post-power stroke.

80 tions:

$$\begin{aligned}
 \frac{dD(t)}{dt} &= (f' - k_{\text{on}}) A_1(t) + (g - k_{\text{on}}) A_2(t) - (k_{\text{on}} + k_{\text{off}} + f) D(t) + k_{\text{on}} \\
 \frac{dA_1(t)}{dt} &= fD(t) + h'A_2(t) - (f' + h) A_1(t) \\
 \frac{dA_2(t)}{dt} &= hA_1(t) - (h' + g) A_2(t)
 \end{aligned} \tag{2}$$

81 The rates which control switching between pre- and post-power stroke states,
 82 h and h' , and f' have been taken as constants. The rates k_{on} and k_{off} as a
 83 function of Ca^{2+} concentration can be expressed as (Campbell et al., 2001):

$$k_{\text{on}} = k_{\text{on}}^0 + (k_{\text{on}}^{Ca} - k_{\text{on}}^0) \frac{[Ca^{2+}]^p}{[Ca_{50}^{2+}]^p + [Ca^{2+}]^p} \tag{3}$$

$$k_{\text{off}} = k_{\text{off}}^0 + (k_{\text{off}}^{Ca} - k_{\text{off}}^0) \frac{[Ca^{2+}]^p}{[Ca_{50}^{2+}]^p + [Ca^{2+}]^p} \tag{4}$$

84 in which k_{on}^{Ca} and k_{on}^0 are the rate constants of the regulatory unit with and
 85 without activation effect of Ca^{2+} , respectively and k_{off}^{Ca} and k_{off}^0 have an equiv-
 86 alent meaning in the “off” situation. $[Ca_{50}^{2+}]$ is the Ca^{2+} concentration level
 87 needed for generating 50% of maximum force.

88
 89 The f rate in Eqs. (2) depends on myofilament distortions and cooperative
 90 effects as follows (Razumova et al., 1999):

$$f = f_0 \left[1 + A_1(t) \left(\exp \left(\frac{x_1}{x_0} (v - 1) \right) - 1 \right) + A_2(t) \left(\exp \left(\frac{x_2}{x_0} (v - 1) \right) - 1 \right) \right]^2 \tag{5}$$

91 where f_0 and v are constant parameters. $v \geq 1$ accounts for the effect
 92 of nearest-neighbor cross-bridges in the force-bearing state contribution. The
 93 unit value of this parameter represents a zero neighboring effect while a larger
 94 number indicates a greater cooperative impact. x_i ($i = 0, 1, 2$) represents the
 95 dimensionless distortion of myofilaments.

96

97 The irreversible detached rate is also affected by the distortions as:

$$g = g_0 e^{\varrho(x_2 - x_0)^2} \quad (6)$$

98 in which g_0 is a constant coefficient and ϱ is a parameter which grades the
99 impact of x_2 .

100 During contraction, elastic deformations or distortions of cross-bridges oc-
101 cur. This average distortion induced by the power stroke during an isometric
102 contraction of a half-sarcomere is denoted by x_0 . The average elastic deforma-
103 tions in A_1 and A_2 states induced through filament sliding during non-isometric
104 contractions are denoted by x_1 and x_2 , respectively (the reader is referred to the
105 work of Heidlauf and Röhrle (2014) for a graphical interpretation of these deforma-
106 tions). Considering these distortions as dimensionless and averaged for the
107 whole tissue, their evolution can be written modifying those initially proposed
108 by Campbell et al. (2001):

$$\frac{dx_1(t)}{dt} = - \left(f \frac{D(t)}{A_1(t)} + h' \frac{A_2(t)}{A_1(t)} \right) x_1(t) + h' \frac{A_2(t)}{A_1(t)} (x_2(t) - x_0) + \frac{\dot{\lambda}_a}{2} \quad (7)$$

$$\frac{dx_2(t)}{dt} = -h \frac{A_1(t)}{A_2(t)} (x_2(t) - (x_1(t) + x_0)) + \frac{\dot{\lambda}_a}{2} \quad (8)$$

109 where $\dot{\lambda}_a$ is the velocity of the sliding between thick and thin filaments.

110 3. Thermodynamic model

111 The deformation associated with muscle activity can be modeled as two
112 fictitious steps (Stålhand et al., 2008; Hernández-Gascón et al., 2013). The first
113 corresponds to the relative motion of the myosin with respect to actin, while
114 the second relates to the elastic deformation of cross-bridges. Mathematically,
115 it can be expressed as a multiplicative decomposition of the tissue stretch as:

$$\lambda = \lambda_a \lambda_e \quad (9)$$

116 where λ_a defines the deformation associated with the contractile response
 117 provoked by the filament sliding, and λ_e represents the deformation due to the
 118 cross-bridges elasticity.

119 3.1. Balance laws

120 Balance laws for the contracting skeletal muscle are derived by means of
 121 the principle of virtual power. This principle states that the external power
 122 equals the internal power, denoted by \mathcal{P} , plus the change of kinetic energy
 123 (Stålhand et al., 2008). By disregarding the temperature, the internal power
 124 can be written as the product of so-called forces and first time derivatives of the
 125 state variables. In this context, the principle of virtual power can be written as:

$$\mathcal{P} = F\dot{\lambda} + F_a\dot{\lambda}_a + F_{Ca}\dot{\beta} \quad (10)$$

126 F , F_a and F_{Ca} are the forces related to λ , λ_a and $\beta = [Ca^{2+}]$ the cal-
 127 cium concentration, respectively. Considering the Clausius-Planck inequality
 128 and ignoring thermal effects (Holzapfel, 2002):

$$\dot{\Psi} \leq \mathcal{P} \quad (11)$$

129 where Ψ is the free energy that coincides with the internal energy and is a
 130 function of the state and internal variables $\mathbf{s} = [R_{\text{off}}, D, A_1, A_2]^T$.

$$\Psi = \Psi(\lambda, \lambda_a, \mathbf{s}, \beta) \quad (12)$$

131 Introducing Eqs. (10) and (12) into Eq. (11):

$$\left(F - \frac{\partial \Psi}{\partial \lambda}\right)\dot{\lambda} + \left(F_a - \frac{\partial \Psi}{\partial \lambda_a}\right)\dot{\lambda}_a - \sum_{i=1}^4 \frac{\partial \Psi}{\partial s_i}\dot{s}_i + \left(F_{Ca} - \frac{\partial \Psi}{\partial \beta}\right)\dot{\beta} \geq 0 \quad (13)$$

132 To observe the thermodynamic compatibility, all evolutions of admissible
 133 state and internal variables should satisfy the above inequality. The simplest

134 and most common choice for constitutive equations which satisfy Eq. (13) is to
 135 consider the non-dissipative relation for first and third terms:

$$F = \frac{\partial \Psi}{\partial \lambda}, \quad F_{C_a} = \frac{\partial \Psi}{\partial \beta}. \quad (14)$$

136 The second term in Eq. (13) is forced to be non-negative as:

$$F_a - \frac{\partial \Psi}{\partial \lambda_a} = C \dot{\lambda}_a \quad (15)$$

137 where $C \geq 0$ is an arbitrary function. Following Stålhand et al. (2008), a
 138 thermodynamic force associated with the chemical states can be defined as:

$$S_i = -\frac{\partial \Psi}{\partial s_i} + a \quad (16)$$

139 with a as an arbitrary multiplier (Stålhand et al., 2008). Taking the con-
 140 straint $\sum_{i=1}^4 s_i = 1$ into account that is equivalent to $\sum_{i=1}^4 \dot{s}_i = 0$:

$$-\sum_{i=1}^4 \frac{\partial \Psi}{\partial s_i} \dot{s}_i = \sum_{i=1}^4 S_i \dot{s}_i \quad (17)$$

141 On the other hand, it is possible to introduce a linear relation for the ther-
 142 modynamic force S_i :

$$\sum_{j=1}^4 x_{ij} \dot{s}_j = S_i \quad (18)$$

143 Considering Eq. (17) and Eq. (18), inequality (13) will be satisfied for the
 144 third term if proposed matrix \mathbf{X} with components x_{ij} satisfies the subsequent
 145 condition:

$$\dot{\mathbf{s}} \cdot \mathbf{X} \dot{\mathbf{s}} \geq 0 \quad (19)$$

146 and the related constitutive equation can be derived combining Eqs. (18)
 147 and (17):

$$\sum_{j=1}^4 x_{ij} \dot{s}_j = -\frac{\partial \Psi}{\partial s_i} + a \quad (20)$$

148 Finally, by suggesting appropriate functions for F_a , F_{Ca} and matrix \mathbf{X} in
 149 terms of the state variables, it is possible to find β from Eq. (14).b, λ_a from Eq.
 150 (15) and the internal variables, s_i , from Eq. (20). Using Eq. (14).a the force F
 151 could be obtained for a given stretch λ .

152 4. Model specialization

153 As mentioned in the previous section, different functions have to be proposed
 154 to obtain the response of the tissue under different levels of activation induced
 155 by the calcium concentration. First, the following form of the strain energy is
 156 considered:

$$\Psi = \Psi_e(\lambda) + \mathcal{N}(\lambda_a)\Psi_a(\lambda_e, s_i) + \Psi_{XB}(s_i) + \Psi_c(\beta) \quad (21)$$

157 where the total energy has been decomposed into several terms related to
 158 the energy associated with the passive behavior of the muscle $\Psi_e(\lambda)$, the elastic
 159 energy in the cross-bridge attachment $\Psi_a(\lambda_e, s_i)$, the energy in the chemical
 160 phase $\Psi_{XB}(s_i)$ and the energy associated with the calcium concentration $\Psi_c(\beta)$.

161 A specific 1D form of the free energy in Grasa et al. (2016) is selected for
 162 the passive muscle behavior:

$$\Psi_e = c_1\left(\lambda^2 + \frac{2}{\lambda} - 3\right) + \Psi'_e \quad (22)$$

$$\Psi'_e = \begin{cases} 0 & \lambda < \bar{\lambda} \\ \frac{c_3}{c_4} \left(\exp [c_4(\lambda^2 - \bar{\lambda})] - c_4(\lambda^2 - \bar{\lambda}) - 1 \right) & \lambda > \bar{\lambda} \end{cases}$$

163 where, c_1 , c_3 and c_4 are material constants and $\bar{\lambda}$ defines the transition
 164 from the linear to the exponential behavior common in collagen fiber reinforced
 165 materials. The energy in the cross-bridge attachment is scaled in Eq. (21) by
 166 the overlapping function $0 \leq \mathcal{N}(\lambda_a) \leq 1$ defined as:

$$\mathcal{N} = \exp \left[\frac{-(\lambda_a^{opt} - \lambda_a)^2}{2\xi^2} \right] \quad (23)$$

167 here λ_a^{opt} represents the stretch of the tissue where the overlapping between
 168 the actin and myosin filaments is optimum, that is, muscle develops the max-
 169 imum force. The parameter ξ governs the band-width of the function. The
 170 elastic energy in cross-bridge attachment Ψ_a can be defined as:

$$\Psi_a = (E_1 s_3 + E_2 s_4) \frac{1}{2} (\lambda_e - 1)^2 \quad (24)$$

171 where E_1 and E_2 are material constants and $(E_1 s_3 + E_2 s_4)$ acts as stiffness
 172 of the cross-bridges.

173 Specific forms of F_a and C functions in Eq. (15) are suggested as follows:

$$F_a = -\nu s_4 \left(\frac{s_3}{\max(s_3)} \right) \mathcal{N}(\lambda_a) \quad (25)$$

$$C = \frac{1}{\mu} (f_1 s_3 + f_2 s_4) \mathcal{N}(\lambda_a) \quad (26)$$

174 where ν , μ , f_1 and f_2 are constants of the model. F_a can be interpreted
 175 as the mechanical force which is generated by the transition of chemical energy
 176 during the power stroke. Alternatively, F_a can be considered as the equivalent
 177 mechanical friction force between filaments (Gestrelus and Borgström, 1986;
 178 Stållhand et al., 2008). In both ways, more cross-bridge connections develop
 179 greater forces, hence $s_3/\max(s_3)$ is considered to scale the active force according
 180 to the number of connections.

181 Substituting Eqs. (21-26) into Eq. (15), the evolution law for $\dot{\lambda}_a$ is obtained:

$$\dot{\lambda}_a = \frac{\mu}{(f_1 s_3 + f_2 s_4)} \left[-\nu s_4 \left(\frac{s_3}{\max(s_3)} \right) + (E_1 s_3 + E_2 s_4) \left(\frac{\lambda_e}{\lambda_a} (\lambda_e - 1) - \frac{(\lambda_a - \lambda_a^{opt})}{2\xi^2} (\lambda_e - 1)^2 \right) \right] \quad (27)$$

182 The explicit statement of total stress in skeletal muscle is achieved by intro-
 183 ducing Eqs. (21) and (9) in Eq. (14).a as follow:

$$F = \frac{\partial \Psi_e}{\partial \lambda} + \frac{\mathcal{N}(\lambda_a)}{\lambda_a} \frac{\partial \Psi_a}{\partial \lambda_e} \quad (28)$$

184 Regarding the chemical phase, the third term in Eq. (21), $\Psi_{XB}(s_i)$ is associ-
 185 ated with cross-bridge free energy and can be additively decomposed assuming
 186 no coupling between the states (Stålhand et al., 2008). To find the chemical
 187 states s_i , in addition to assigning the $\Psi_{XB}(s_i)$ function, it would be necessary
 188 to find matrix \mathbf{X} which can be a function of λ_a , λ_c and β . Therefore, specifically
 189 designed experiments would be required to find the coefficients for $\Psi_{XB}(s_i)$ and
 190 components of \mathbf{X} . To the best of the authors' knowledge, no such investigation
 191 which measures the chemical and mechanical coupling in sarcomeres has been
 192 reported in the literature. However, Stålhand et al. (2008) proved that a system
 193 of differential equations like Eq. (2) is a specific form of Eq. (20). Therefore,
 194 in this study s_i have been obtained solving Eq. (2).

195 The last term of Eq. (21), $\Psi_C(\beta)$ is the free energy related to the calcium ion
 196 concentration in the skeletal muscle cell. Considering Eq. (14).b, the simplest
 197 function for this chemical free energy (Stålhand et al., 2008) is:

$$\Psi_C = \frac{1}{2}\beta^2 \quad (29)$$

198 The chemical force, F_{Ca} , can be evaluated for a non-dissipative chemical
 199 state by introducing Eqs. (29) and (21) in (14) giving the result:

$$F_{Ca} = \beta \quad (30)$$

200 With the information about Ca^{2+} , assigning initial values to the four cross-
 201 bridge states, the two myofilament distortions and assuming that at the be-
 202 ginning of force generating $\dot{\lambda}_a = 0$, it would be possible to solve Eqs (2), (7),
 203 (8) and (1). The initial values of $[R_{off}, D, A_1, A_2, x_1, x_2]$ have been chosen as
 204 $[0.999997, 10^{-5}, 10^{-5}, 10^{-5}, 0, 10^{-2}/1.6]$, respectively. With these levels, the ac-
 205 tive velocity of contraction $\dot{\lambda}_a$ can be updated using Eq. (27). Therefore, active

206 and elastic stretches can be computed using $\dot{\lambda}_a$ along with Eq. (9). Finally,
 207 the muscle force will be determined using Eq. (28). The solving process can be
 208 summarized with the equation system as follows:

$$\begin{aligned}
 \sum_{i=1}^4 s_i &= 1 \\
 [\dot{s}_i] &= A_{ij}(Ca, \mathbf{x}) s_j + B_i \quad i = 2 - 4 \\
 [\dot{x}_i] &= C_{ij}(Ca, \mathbf{s}) x_j + D_i \quad i = 1 - 2 \\
 \dot{\lambda} &= E(\mathbf{x}, \mathbf{s}, \lambda_a, \lambda_e) \\
 F &= F(\mathbf{s}, \lambda_a, \lambda_e)
 \end{aligned} \tag{31}$$

209 5. Model assessment and results

210 To study the ability of the model to predict muscle fatigue, different as-
 211 sumptions for the parameters that define its behavior were established due to
 212 the lack of experimental results for identical animal species. Initially, as shown
 213 in the following section, the unfatigued behavior of the muscle was analyzed to
 214 reproduce a single tetanic contraction and the force-velocity relationship. Then,
 215 the effect of the Ca^{2+} concentration was assessed in the force/stress response.
 216 Finally, the model was adjusted to reproduce the results reported by Sierra et al.
 217 (2017) for the rabbit *Extensor Digitorum Longus* (EDL) muscle, which provided
 218 the muscle force evolution during repeated isometric contractions.

219 5.1. Model configuration and validation

220 As mentioned, Ca^{2+} release from the sarcoplasmic reticulum is the input of
 221 the model which starts the cross-bridge cycling and the production of muscle
 222 force. In this work, the Ca^{2+} evolution has been taken from experimental
 223 data reported in the literature (Calderón, 2013) avoiding the definition of a
 224 mathematical expression for β and consequently for Ψ_C and F_{Ca} . However, in
 225 (Calderón, 2013) the maximum level of the steady-state calcium concentration
 226 for the EDL muscle was not reported. Therefore, a single Ca^{2+} transient signal
 227 has been scaled following data reported in Stephenson and Williams (1982).

228 Furthermore, other calcium properties ($[Ca_{50}^{2+}]$ and p in Eqs. (3) and (4))
229 have been obtained from the $F-pCa$ curve of the EDL muscle reported in
230 Stephenson and Williams (1982) ($pCa = -\log([Ca^{2+}])$ and $pCa_{50} = -\log([Ca_{50}^{2+}])$,
231 see Table 1). The calcium signal for maximal (tetanic) or submaximal contrac-
232 tions was obtained by the summation of the twitch signal according to the
233 stimulation time and frequency.

234
235 A Matlab nonlinear least-squares curve fitting program was used to adjust
236 the simulation result with the experimental isometric contraction of rabbit EDL
237 muscle (Sierra et al., 2017). This fitting allowed all the parameters involved in
238 the model to be obtained using as initial seeds those proposed by Razumova et al.
239 (1999) for the rates related to cross-bridge dynamics. The passive mechanical
240 behavior was taken from the literature (Calvo et al., 2010). A summary of
241 the final parameter values is presented in Table 1 distinguishing between those
242 related to cross-bridges, macroscopic mechanical behavior and calcium kinetics.
243 The simulated isometric stress for a 0.2 s contraction stimulated at 100 Hz
244 together with the experimental contraction of the unfatigued EDL muscle is
245 represented in Fig. 2. The isometric stress reaches the maximum value (0.433
246 MPa) at 0.14 s, remains at this level for a short time and starts to decrease with
247 slow and then fast rates.

248 The ability of the model to reproduce concentric contractions was also eval-
249 uated. A set of different load levels was introduced to obtain the well-known
250 force-velocity relationship. In Fig. 3 the simulation results are compared to the
251 experimental data of EDL rat muscle at $35^{\circ}C$ (Ranatunga, 1984). The muscle
252 response is presented as the shortening velocity against different percentages
253 of the maximum isometric stress. With increasing the resistance tension stress
254 from 5% to 50% of maximum isometric stress, the shortening velocity was de-
255 creased from 125.2 to 29.9 mm/s.

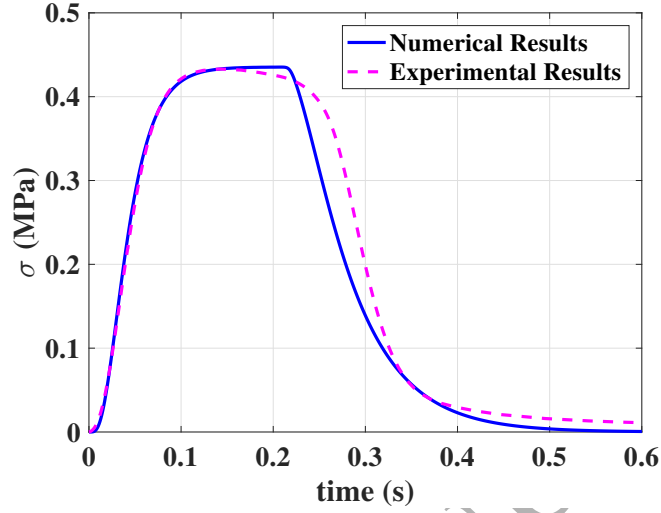


Figure 2: Isometric contraction for unfatigued EDL muscle of rabbit (the dashed line shows the experimental data (Sierra et al., 2017) and the solid line the model simulated result)

Cross Bridge Parameters

k_{on}^0 (1/s)	0	f_0 (1/s)	123.623	g_0 (1/s)	1.001
k_{on}^{Ca} (1/s)	51.937	f' (1/s)	477.470	ϱ (-)	1
k_{off}^0 (1/s)	12.115	h (1/s)	8.213	v (-)	1.001
k_{off}^{Ca} (1/s)	174.778	h' (1/s)	50.626	x_0 (-)	0.006

Mechanical Parameters

E_1 (N)	143.815	ν (N)	51	c_1 (N)	0.001
E_2 (N)	567.842	μ (-)	2.229	c_3 (N)	0.054
f_1 (N·s)	9.113	ξ (-)	0.164	c_4 (N)	0.783
f_2 (N·s)	13.058	λ_{opt} (-)	1	$\bar{\lambda}$ (-)	1.254

Calcium Properties

pCa_{max}	4.975	pCa_{50}	5.816	p	4
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Table 1: Model parameters determined for the EDL rabbit muscle. The parameters were obtained fitting the passive and active behavior characterized experimentally.

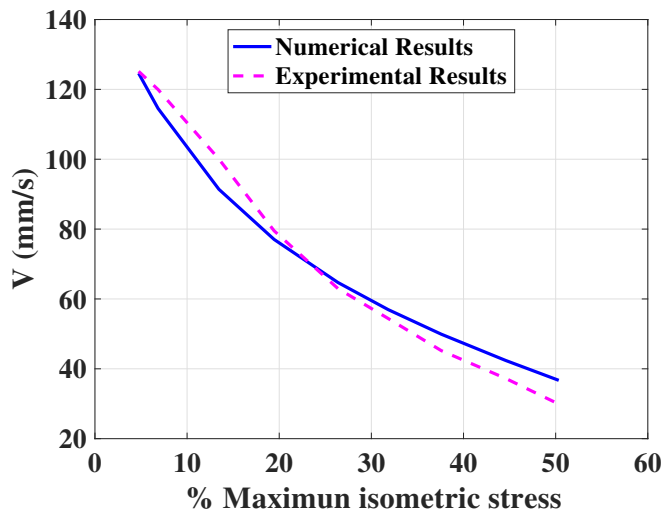


Figure 3: Experimental (Ranatunga, 1984) and simulated force-velocity relationship for rat EDL muscle.

256 5.2. Fatigue simulation

257 Once the potential of the model has been assessed for isometric and concentric contractions, in this section the role of calcium is evaluated during fatigue.
 258
 259 As mentioned, this phenomenon has a distinctive effect on the mechanical performance of muscles and is associated with a depression in both the $[Ca^{2+}]$ level
 260 and the myofibrillar sensitivity. However, the $[Ca^{2+}]$ binding cooperatively remains the same (Chin and Allen, 1998; Debold et al., 2006; Fitts, 2008). In fact,
 261 an increase in the level of $[Ca_{50}^{2+}]$, requires more $[Ca^{2+}]$ to produce the same level of force. Experimental results of the $[Ca_{50}^{2+}]$ increment in fast skeletal muscle
 262 fiber reported by Chin and Allen (1998) have been extrapolated according to the intensity of the fatigue test (stimulating the muscle with 100 Hz every 10
 263 second for one hour (Sierra et al., 2017)). It has also been assumed that $[Ca_{50}^{2+}]$
 264 will change linearly at different levels of fatigue as well.
 265
 266
 267
 268

269 Fig. 4 represents the experimental evolution of maximal repetitive isometric contractions of rabbit EDL muscle (Sierra et al., 2017). The tissue was
 270 stimulated under tetanic conditions for 0.2 s every 10 s during one hour. As
 271

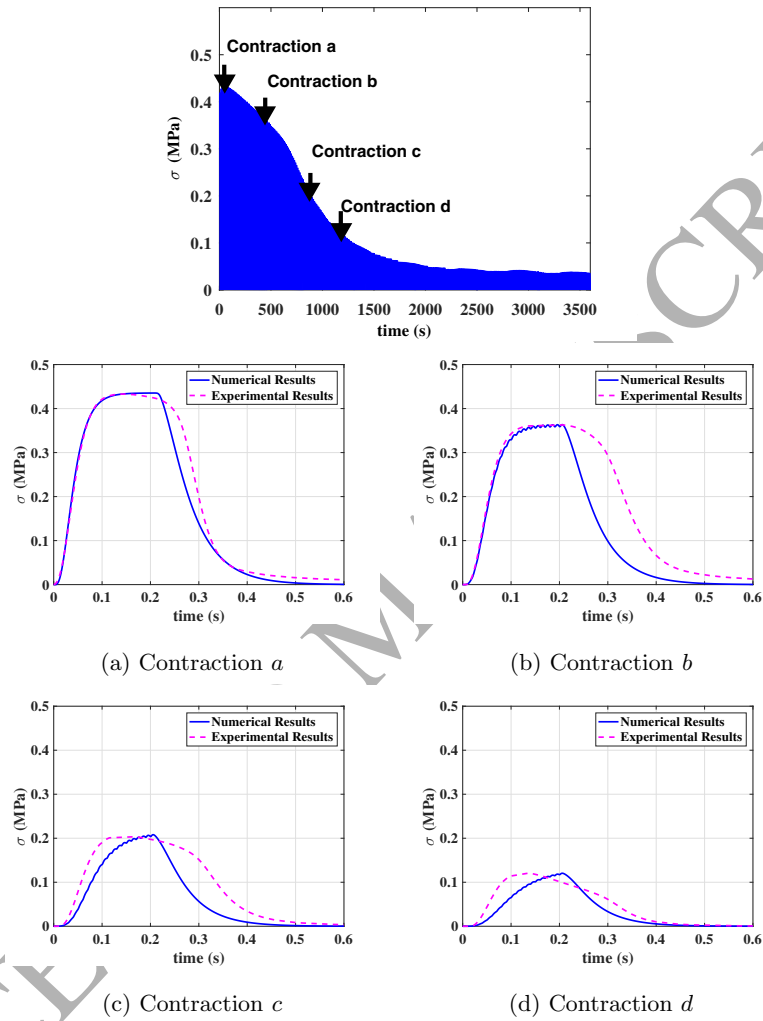


Figure 4: Four experimental and the equivalent simulated isometric contractions of rabbit EDL muscle fatigue tests. (Top panel) Continuous tetanic contraction; each vertical line shows a tetanus which is simulated for 0.2 s with 100 Hz after 10 s rest. (Bottom panel) Four extracted contraction-relaxation cycles of rabbit EDL in (a) unfatigued muscle, (b) 16%, (c) 53% and (d) 72% reduction of maximum force due to fatigued conditions.

	pCa	pCa_{50}
Contraction <i>a</i>	4.9748	5.8158
Contraction <i>b</i>	5.7017	5.8111
Contraction <i>c</i>	5.9111	5.8004
Contraction <i>d</i>	5.9807	5.7949

Table 2: Ca^{2+} concentration levels for the four simulated contractions (at 30 s, 400 s, 800 s and 1200 s) considered in the fatigue processes.

indicated in the figure (vertical arrows) four contractions have been considered to fit with the model. The first contraction (a) was selected at the beginning of the test when the maximum isometric stress is reached. In the second selected contraction, the muscle developed approximately 84% (b), in the third 47% (c) and in the last 28% of its maximum stress. The fitting results for these four contractions can be observed in the same figure with the experimental contractions. Estimated levels of pCa and pCa_{50} for these simulated contractions are presented in Table 2, varying from $pCa = 4.9748$ and $pCa_{50} = 5.8158$ in the unfatigued contraction to $pCa = 5.9807$ and $pCa_{50} = 5.7949$ for the highest fatigue level. As can be observed, $[Ca^{2+}]$ has to decrease to fit the maximum level of force.

Using the $[Ca^{2+}]$ and $[Ca_{50}^{2+}]$ levels determined for the four selected fatigued contractions in Fig. 4, the force-velocity relationship has been reproduced for a set of external loads from 5% to 80% of the maximum generated force in each condition (Fig. 5). In Fig. 5, the predicted velocity has been normalized by the maximum velocity of the unfatigued condition. As can be observed, the results of the model show that along with the increase in the fatigue level, the contractile velocity of the muscle decreases, that is, the most fatigued condition has the lowest shortening velocity. Moreover, the velocity reduction has a non-linear behavior looking at fixed levels of the resistance force reached.

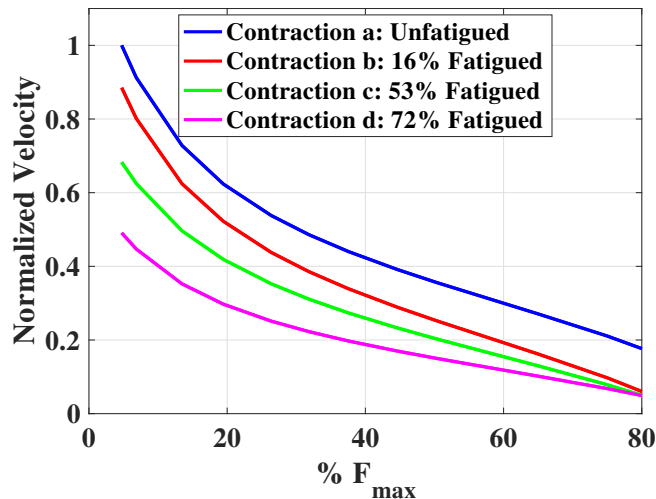


Figure 5: Simulated force-velocity curve in unfatigued and fatigued conditions. Ca properties in each condition are reported in Table 2. Velocity data has been normalized to the maximum value of the unfatigued contraction.

292 6. Discussion

293 The multi-scale model developed in this work, using the $[Ca^{2+}]$ level as an
 294 input governing parameter of the skeletal muscle contraction, is able to repro-
 295 duce the most important characteristic features of fatigue such as force deficit,
 296 reduction in shortening velocity and alterations in the contraction-relaxation
 297 cycle. In previous studies related to muscle fatigue, it has been accepted that
 298 this force decreasing phenomenon is a consequence of a lower number of connec-
 299 tions in post- and pre-power stroke states (Fitts, 2008; Place et al., 2009; Jones,
 300 2010). Although the variation in other chemical factors could influence the
 301 characteristics of the Ca^{2+} signal and functional features of the muscle (Allen
 302 and Westerblad, 2001; Godt and Nosek, 1989; Debold et al., 2016; Allen et al.,
 303 2008), they have not been considered in this model. This limitation could be
 304 responsible of the differences in the fitting of the relaxation phase and in the
 305 fatigued states in Figs. 2 and 4.

306 Cross bridge cycling starts after an increase in the Ca^{2+} concentration level,

307 and the proportion of connections in the force bearing states (A_1 and A_2) are
 308 determined according to different transition rates. The available experimental
 309 information for these rates is limited to just two states (attached or detached)
 310 of the actin and myosin filaments (Brenner, 1988; Metzger et al., 1989; Sweeney
 311 and Stull, 1990; Rome et al., 1999). The initial data in the model has been
 312 taken from Razumova et al. (1999) which corresponds to a proposed feasible
 313 group of rate coefficients and those in Campbell et al. (2001) and Rice et al.
 314 (2008) relating to cardiac muscle fibers. The final rates found in the present
 315 model by fitting the experimental rabbit EDL results are slightly different from
 316 those mentioned and in general very similar to other muscles and species. For
 317 example, the detachment rate ($g = 1.0004 \text{ s}^{-1}$) is in the range of previously
 318 reported experimental measurements for rabbit fast-twitch psoas (Rome et al.,
 319 1999).

320 It has also been shown that these rates depend on various factors such as
 321 calcium concentration (Metzger and Moss, 1990), myofilaments distortion, co-
 322 operation effects (Razumova et al., 2000) and temperature (Stephenson and
 323 Williams, 1981; Rice et al., 2008). Although the dependency on temperature
 324 has been neglected, the effects of other parameters have been considered, even
 325 if in a simple form. According to Eqs. (3) and (4), the k_{on} and k_{off} rates change
 326 their rest value (k_{on}^0 and k_{off}^0) depending on the Ca^{2+} concentration. As the XB
 327 kinematics cycle starts according to the k_{on} and k_{off} values, all four states of
 328 the XB cycle will be affected by $[Ca^{2+}]$ (Razumova et al., 2000). As mentioned
 329 before, the v parameter in Eq. (5) models the effect of the XB neighboring con-
 330 tribution on f rate. Consequently, f_0 in Eq. (5) is equivalent to the f rate when
 331 none of the neighboring XBs are in force bearing condition (pre- and post-power
 332 stroke - A_1 and A_2) ($v = 1 \Rightarrow f = f_0$). On the other hand, myofilaments distor-
 333 tion will change the space between contiguous XBs and accordingly their level
 334 of contribution to the force bearing status. Therefore, in the f rate the nearest
 335 neighbor influence has been scaled with respect to the myofilament distortion
 336 (x_1, x_2) (Razumova et al., 1999). Moreover, the detachment rate g depends on
 337 the myofilament distortion in the post-power stroke condition (x_2). Accord-

338 ing to Eq. (6), g_0 can be interpreted as the detachment of XBs in isometric
 339 condition (Razumova et al., 1999).

340 The derivation of constitutive equations from balance laws (Stålhand et al.,
 341 2008; Hernández-Gascón et al., 2013) has resulted in a good performance in nu-
 342 merical simulations of living tissues (Stålhand et al., 2011; Grasa et al., 2016).
 343 The incorporation of cross-bridge cycling at the micro level to the elastic energy
 344 function has required some additional approximations. Although Razumova
 345 et al. (1999) considered a fixed value of stiffness for the cross-bridges attached
 346 states (A_1 and A_2), the model assumed different stiffnesses (E_1 and E_2 , re-
 347 spectively). These macroscopic stiffnesses are actually the product of the cor-
 348 responding state stiffness in the sarcomere level by the total number of related
 349 attachments in the whole muscle. The larger value of E_2 compared to E_1 is re-
 350 lated to the hypothesis that the post-power stroke attachment is stronger than
 351 the pre-power stroke (Jones, 2010; Fitts, 2008). Moreover, there is some ex-
 352 perimental evidence that during fatigue condition the stiffness per cross-bridge
 353 decreases (Fitts, 2008; Place et al., 2009; Nocella et al., 2017). However, this
 354 variation has been ignored.

355 Data recorded from calcium distribution signals in single fibers of skeletal
 356 muscle under fatigue conditions have confirmed that $[Ca^{2+}]$ declines along with
 357 the fatigue phenomenon (Westerblad and Allen, 1993, 1991). Moreover, under
 358 fatigue conditions, the resting threshold of Ca^{2+} rises as a consequence of a
 359 decrease in calcium sensitivity. This circumstance is accompanied by a $[Ca_{50}^{2+}]$
 360 increase. It has been shown that with decreasing $[Ca^{2+}]$ and increasing $[Ca_{50}^{2+}]$,
 361 the presented model is able to predict different levels of fatigue. However, exper-
 362 imental studies of fatigue have demonstrated that in addition to the intracellular
 363 Ca^{2+} variation, the rate of sarcoplasmic reticulum (SR) Ca^{2+} uptake also varies
 364 (Westerblad and Allen, 1993). This SR environment and its chemical kinetics
 365 has not been considered in the present model. Furthermore, according to the
 366 literature (Westerblad and Allen, 1993; Rome, 2006), one of the effective factor
 367 in skeletal muscle relaxation is Ca^{2+} uptake. Differences between the experi-
 368 mental data and simulated results in the relaxation part of the contraction cycle

369 could be due to neglecting these effects.

370 The shortening velocity against constant resistive forces in concentric con-
371 ditions has also been simulated after fatigue (Fig. 5). Similar to experimental
372 findings (Haan et al., 1989; Curtin and Edman, 1994; Ruiter et al., 2000), the
373 model is able to predict the decreasing effect on the velocity of contraction along
374 with the increase in fatigue level.

375 7. Conclusions

376 A multi-scale chemo-mechanical model for predicting the active behavior
377 of skeletal muscle under both unfatigued and fatigued conditions has been pre-
378 sented. The constitutive framework of the model adheres to the thermodynamic
379 laws so that it is suitable to be applied in 3D simulations. A comparison of the
380 computational results with experimental data demonstrates the ability of the
381 model to simulate isometric and concentric contractions and the force-velocity
382 relationship in skeletal muscles. Due to the lack of experimental results to deter-
383 mine the whole set of parameters for the same animal species, slightly differences
384 counteracting effects in the proposed values could provide the same model out-
385 come. Thereby, the selection was made carefully always in the range of the
386 parameters available in the literature.

387 The chemical features of the model make it appropriate for investigating
388 chemically dependent phenomena such as fatigue. Good agreement was found
389 between the simulated and experimental results in the force development of
390 fatigued muscle. Moreover, it has been shown that the model can predict the
391 reduction of contractile velocity as a consequence of fatigue inefficiency. As the
392 model is activated by the calcium ion transient signal, it would be possible to
393 predict different levels of muscle activation (twitch summation) based on the
394 calcium signal accumulation. Therefore, when defining the relation between
395 electrical stimulation and released intracellular calcium ion, the model could be
396 useful for simulating techniques such as Functional Electrical Stimulation.

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