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Effects of Zn Substitution in the Magnetic and Morphological Properties of Fe-Oxide Based Core-Shell Nanoparticles Produced in a Single Chemical Synthesis

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Abstract:

Magnetic, compositional and morphological properties of Zn-Fe-oxide core-shell bimagnetic nanoparticles (MNPs) were studied for three samples with 0.00, 0.06 and 0.10 Zn/Fe ratios, as obtained from Particle-Induced X-ray Emission analysis. The bimagnetic nanoparticles were produced in one step synthesis by the thermal decomposition of the respective acetylacetonates. The nanoparticles present an average particle size between 25 - 30 nm as inferred from Transmission Electron Microscopy (TEM). High-Resolution TEM images clearly show core-shell morphology for the particles in all samples. The core is composed by an antiferromagnetic (AFM) phase with Wüstite (Fe_{1-y}O) structure, while the shell is composed by $\text{Zn}_x\text{Fe}_{3-x}\text{O}_4$ ferrimagnetic (FiM) spinel phase. Despite the low solubility of the Zn in the Wüstite, Electron energy-loss spectroscopy (EELS) analysis indicates that the Zn is distributed almost homogeneously in the whole nanoparticle. This result gives information on the formation mechanisms of the particle, indicating that the Wüstite is formed firstly, and the superficial oxidation results in the FiM ferrite phase with similar Zn concentration than the core. Magnetization and in-field Mössbauer spectroscopy of the Zn-richest nanoparticles indicate that the AFM phase is strongly coupled to the FiM structure of the ferrite shell, resulting in a bias field (H_{EB}) appearing below $T_{\text{N}}^{\text{FeO}}$, with H_{EB} values that depends on the core-shell relative proportion. Magnetic characterization also indicates a strong magnetic frustration for the samples with higher Zn concentration, even at low temperatures.

INTRODUCTION

In the early 2000's, the groups of Sun^{1,2} and Hyeon^{3,4} reported the synthesis of magnetic ferrite nanoparticles by thermal decomposition at relatively high temperatures (about 540-640 K) of organo-metallic precursors in the presence of surfactants and a long-chain alcohol. This method becomes a hit since it produces magnetic nanoparticles with several advantages with respect to others⁵, such as a narrow size distribution and a high crystallinity, both critical for homogeneous magnetic properties, and simultaneously it allows an easy way to control stoichiometry². The success of this method is based in a well separately nucleation and growth/crystallization process that take place at different temperatures⁶, being extensively used to synthesize magnetic nanoparticles⁷⁻¹².

An interesting feature was reported on the fundamentals of the thermal decomposition method by Hao et al in 2007¹³: a massive reduction of iron ions in the synthesis solution, leading to the formation of nanoparticle with diameter larger than 14 nm containing the antiferromagnetic Wüstite phase (Fe_{1-y}O), which is followed by the oxidation and consequent formation of a ferrimagnetic (FiM) ferrite phase (Fe_3O_4 or $\gamma\text{-Fe}_2\text{O}_3$). The result was a bimagnetic core-shell nanoparticle produced in a one-pot reaction, with an antiferromagnetic (AFM) core (Fe_{1-y}O) below $T_N = 198 \text{ K}$ ¹⁴ and a FiM shell. In fact, in the work of Hyeon et al in 2004³, it is possible to observe remarkable features around 200 K in the ZFC Magnetization *vs.* Temperature curves for large particles ($d > 16 \text{ nm}$) that could be related to Néel temperature (T_N) of the AFM phase. When the formation of Fe_{1-y}O phase took place during the synthesis, the presence of the AFM phase is expected for larger particles since the oxidation to the FiM phase could not be complete, keeping an AFM core. Other works report on the formation of Fe_{1-y}O and a ferrite phase in a core-shell structure after oxidation by the thermal decomposition method¹⁵⁻¹⁹. As pointed out by Chen et al²⁰, the Wüstite phase is formed as consequence of reactions involving some commonly used hydrocarbon solvents and its formation contributes to significant cationic disorder. The later formation of the bimagnetic core-shell system due to the surface oxidation of Wüstite results in very distinct magnetic properties with respect to the monophasic nanoparticles²¹⁻²⁶. Therefore, it is very important determine the factors that govern the core/shell nanoparticles formation to establish the reproducibility of this synthesis method and the magnetic response of the system. Different magnetic ions present different degree of

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3 solubility in the Wüstite phase. While Chen et al²⁷ have determined a relatively
4 homogeneous distribution of Co ions along the entire particle, including core (Wüstite) and
5 shell (ferrite) and similar observation was found for the Mn ions²⁸ and also for the
6 diamagnetic Mg²⁺ ion²⁹⁻³¹; the Zn²⁺ ions present low solubility in the Wüstite with an upper
7 limit of approximately 10 %³². This result can lead to a stoichiometric heterogeneity when
8 the formation of Fe_{1-y}O takes place with important consequences in the magnetic response
9 of the system. As the production of iron oxides nanoparticles doped with Zn is of great
10 current interest due to the surface chemistry of the system as well its magnetic response for
11 bio-application and innovative technologies³³⁻³⁶, it is crucial determine the role of the
12 different parameter that establish the morphology and magnetic properties of this system.
13 With this premise, in this work we have produced by the thermal decomposition method
14 Zn-Fe_{1-y}O/Zn_xFe_{3-x}O₄ bimagnetic core-shell nanoparticles and study its structural,
15 morphology and magnetic properties. We have analyzed the effect of the low solubility of
16 the Zn in the Wüstite in the composition homogeneity and its effects on the nanoparticles
17 physical properties.
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32 **EXPERIMENTAL PROCEDURE**

33 *Sample Preparation*

34 Three sample with distinct Zn amounts in the composition were prepared by the thermal
35 decomposition of the Fe³⁺ acetylacetonate (Fe(C₅H₇O₂)₃ – Fe(acac)₃ - Aldrich 97%) and
36 Zn²⁺ acetylacetonate (Zn(C₅H₇O₂)₂·xH₂O – Zn(acac)₂ – Aldrich, with estimation of 20 %
37 wt. of water as obtained from DTA-TGA measurements). Samples were labeled according
38 with the Zn/Fe relation determined by Particle-Induced X-ray Emission (PIXE) analysis (z
39 = [Zn]_{PIXE}/[Fe]_{PIXE}): $z = 0.10$, $z = 0.06$, $z = 0.00$. For sample $z = 0.00$, a total of 5.7 mM of
40 Fe(acac)₃ were dispersed in 50 mL of octadecene (CH₃(CH₂)₁₅CH=CH₂ – Aldrich 90 %) with
41 56.6 mM of Oleic acid (CH₃(CH₂)₇CH=CH(CH₂)₇COOH - PanReac, PharmaGrade),
42 56.6 mM of oleylamine (CH₃(CH₂)₇CH=CH(CH₂)₇CH₂NH₂ – Aldrich 70%) and 5.7 mM
43 of 1,2-octanediol (CH₃(CH₂)₅CH(OH)CH₂OH – Aldrich 98%). Sample $z = 0.06$ was
44 prepared with 4.5 mM of Fe(acac)₃ and 0.5 mM of Zn(acac)₂ dispersed in 50 mL of
45 octadecene and adding 26 mM of Oleic acid, 26 mM of oleylamine and 3.5 mM of 1,2-
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3 octanediol. Sample z = 0.10 was prepared using 4.6 mM of Fe(acac)₃ and 1.1 mM of
4 Zn(acac)₂ dispersed in 50 mL of trioctylamine ([CH₃(CH₂)₇]₃N – Aldrich 98 %) and 22.5
5 mM of oleic acid. For all samples, a 3-neck flask was used with one neck for the N₂ (White
6 Martins - 5.0) flux, other for a mechanical stirring (the central one) and the third one for a
7 condenser. The synthesis solution of all samples were mechanically stirred at room
8 temperature, heated to 373-393 K during 30 minutes with a N₂ flux of 0.5 mL/min, without
9 the presence of condenser in order to remove humidity of the system. After that, the system
10 was closed with the condenser keeping the N₂ flux, and the solution was heated up to 473 K
11 with a controlled increment of temperature (about 3 K/min.). The solution was kept in this
12 temperature about 5 minutes and then it was heated up to the reflux condition with a
13 heating ramp of ~20 K/min: 610-615 K for octadecene and 630-635 K for trioctylamine.
14 After that, the heating source was removed in order to get a rapid cooling of the solution.
15 At 313 K, the flask is opened and 400 mL of Ethanol (96 % - Porta) is added in order to
16 precipitate the particles. Finally, a solution of Ethanol and Acetone (99.9 % - Sigma-
17 Aldrich) (about 10:1, respectively) was added on the particles followed by centrifugation at
18 14000 rpm for 30 minutes. A second cleaning only with acetone is performed, also
19 followed by centrifugation. The final product, the as-made sample, is a powder containing
20 hydrophobic nanoparticles coated with an organic layer of oleic acid.
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24 A differential thermal thermogravimetric analyzer (TGA) Shimadzu DTG-60H was used
25 to determinate the percentage of organic compound in the as-made nanoparticles to have
26 accuracy in the mass determination to normalize magnetization experimental data with the
27 oxide amount. For this, each as-made powder sample was heating up to 1173 K with a
28 heating rate of 3 K/min in Ar flux (100 mL/min) while the weight loss was measured,
29 which corresponds to the amount of organic compound in nanoparticles. This technique
30 also was used to determine the amount of water in the Zn(acac)₂, where the reagent was
31 heated up to 453 K with a heating rate of 3 K/min in Ar flux (100 mL/min).
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35 Structural characterization was carried out using a PANalytical Empyrean powder X-ray
36 diffractometer (CuK_α, λ=0.15418 nm). Samples were cleaning with acetone, to remove
37 organic material and a zero diffraction plate holder was used in measurement.
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41 Transmission Electron Microscopy (TEM) images were carried out in two microscopies.
42 For both microscopes, TEM specimens were prepared by dispersing the dried powdered
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3 samples in isopropanol and dropping the solution onto a carbon covered grid. For TEM and
4 High-Resolution TEM (HRTEM), the images were obtained with a Tecnai F20 G2
5 (ThermoFisher Scientific, formerly FEI) microscope operating at 200 kV at room
6 temperature. The mean particle size ($\langle D \rangle$) was determined after count ~ 300 particles and
7 fitting the corresponding histogram with a lognormal distribution. For the Scanning
8 Transmission Electron Microscopy - High Angle Annular Dark Field (STEM-HAADF), the
9 images were obtained in a Cs-probe-corrected Titan (ThermoFisher Scientific, formerly
10 FEI) at a working voltage of 300 KV, coupled with a HAADF detector (Fischione). The
11 chemical composition of the nanoparticles was analyzed by Electron Energy Loss
12 Spectroscopy (EELS), in order to obtain the concentration ratio Fe/Zn at the core and the
13 shell. Low-loss EEL spectra were obtained with a Tridiem Energy Filter (Gatan)
14 spectrometer. For every linescan, a 5 ms spectrum was obtained every 0.1 nm. The Energy
15 dispersion used for the analysis was 0.1 eV/pixel. The collection semi-angle (Beta) was
16 32.6 mrad for a camera length of 10 mm and a spectrometer entrance aperture of 1 μm . The
17 convergence semi-angle (Alpha) was 24.8 mrad for a condenser aperture of 70 μm . The
18 energy resolution, estimated from the full width at half maximum of the zero-loss peak, was
19 1.4 eV.

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32 Chemical composition of the nanoparticles was measured with Particle induced X-ray
33 emission spectroscopy (PIXE) technique³⁷, with using a 3 MeV H^+ beam in a NEC 5SDH
34 1.7 MV tandem accelerator with a NEC RC43 end-station³⁸. To prepare the nanoparticles
35 samples they were properly washed and dried and then placed on the sample holder with
36 carbon tape. PIXE spectra were analysed using GUPIX software³⁹.

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41 Magnetization measurements as function of temperature ($M(T)$) and isothermal curves ($T =$
42 5 K) as function of applied field ($M_{5\text{K}}(H)$) were measured in a commercial Quantum
43 Design MPMS-5S SQUID magnetometer. $M(T)$ curves were measured applied field of 50
44 Oe in Zero Field Cooling (ZFC) and Field Cooling (FC) protocols. $M_{5\text{K}}(H)$ curves were
45 measured in both conditions: after ZFC protocol and after a FC protocol. For the FC
46 protocol, the sample was cooled from $T = 300$ K to 5 K (starting above the $T_N \sim 200$ K of
47 Fe_{1-y}O AFM phase) with applied field of $H_{\text{FC}} = 50$ kOe. Isothermal curves at $T = 100$ K
48 ($M_{100\text{K}}(H)$) and 300 K ($M_{300\text{K}}(H)$) were measured in a LakeShore 7300 vibrating sample
49 magnetometer (VSM). $M_{100\text{K}}(H)$ measurements were also performed in both ZFC and FC
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3 protocols, with an applied field in the cooling procedure of $H_{FC} = 10$ kOe. $M_{300K}(H)$
4 measurements were used in order to determine saturation magnetization (M_S) at room
5 temperature. For all magnetic measurements, the samples were conditioned by dispersing
6 the nanoparticles in toluene and paraffin, followed by the toluene evaporation, in order to
7 avoid agglomeration and to immobilize them.

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11 In-field Mössbauer spectra (MS) of sample $z = 0.10$ were taken at 4.2 K in a liquid He flow
12 cryostat with a spectrometer in transmission geometry with using a $^{57}\text{Co}/\text{Rh}$ source. The
13 spectra were collected after a field cooling procedure with applied field of $H_{FC} = 120$ kOe
14 from 300 K (above T_N). Mössbauer samples were prepared by dispersing 20 mg of the
15 respective powder samples in boron nitride and pressing it between acrylic discs. The
16 sample was mounted in the bore of a 140 kOe superconducting magnet, in a vertical source-
17 sample-detector setup such that the direction of gamma-ray was parallel to the direction of
18 applied field. A sine-shaped velocity waveform was used to minimize mechanical noise.
19 All MS spectra were fitted by using Lorentzian line shapes with non-linear least-square
20 software, with calibrating the velocity scale with α -Fe foil at 300 K. Isomer shift values are
21 relative to α -Fe at 300 K. As a consequence of the experimental setup, the temperature of
22 the source during the low-temperature spectrum acquisition must be considered as an
23 important factor to the isomer shift (IS) values observed in order to compare with the
24 expected ones.
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38 RESULTS AND DISCUSSION

39 *Morphology and Composition*

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41 Figure 1 shows the measured PIXE spectra of all as-made samples and their corresponding
42 fittings performed with GUPIX software. The procedure of fabrication and sample
43 preparation involves the presence of extra elements, mainly carbon and oxygen due to the
44 imponderable contribution from different sources, like carbon tape, solvent remains and
45 contamination. Restricting us to the nanoparticles composition, just the Zn and Fe
46 quantities are taking into the account using the peaks associated to the K-shell transitions.
47 The value $z = [\text{Zn}]_{\text{PIXE}}/[\text{Fe}]_{\text{PIXE}}$ for each sample was used to label each sample in Table 1,
48 with the corresponding uncertain in the parenthesis. The amount of Zn in samples $z = 0.10$
49 (2) and 0.06 (2) are lower than the expected from the molar relation of precursors
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([Zn(acac)₂]/[Fe(acac)₃]) used in the synthesis procedure: 0.58 for sample $z = 0.10$ (2) and 0.30 for sample $z = 0.06$ (2). This result indicates that the incorporation of the Zn ions to the nanoparticles is limited in the synthesis procedure, probably as consequence of the low Zn solubility in the Wüstite phase.

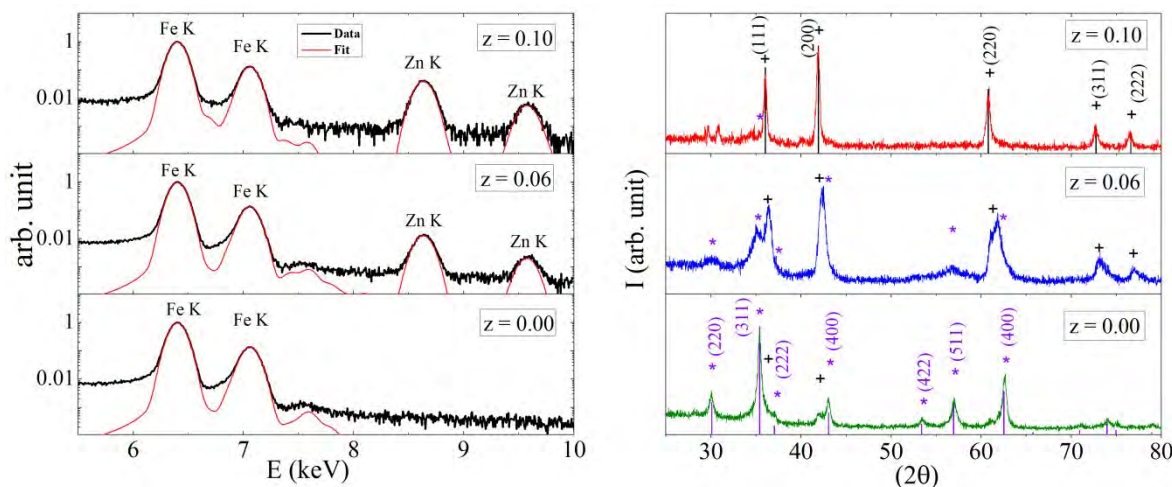


Figure 1 – Left Panel: PIXE spectra of all samples showing the pertinent energy range of Fe and Zn K-peaks. Black lines correspond to experimental data, and red lines are the respective fitting to obtain the $z = [\text{Zn}]/[\text{Fe}]$ ratio for each system. Right panel: XRD profiles of three samples. All peaks are indexed with the Fe_{1-y}O (+) and ferrite (*) phases where the bars correspond to the diffraction peaks according to the JCPDS card number: for Wüstite 00-006-0615 (upper panel) and magnetite 00-019-0629 (bottom panel).

Figure 1 also presents the XRD patterns of samples $z = 0.10$, 0.06 and 0.00 . The diffraction pattern of sample with $z = 0.10$ mainly corresponds to the phase Fe_{1-y}O, despite some peaks of ferrite phase are observed with very low intensity. The peaks (111), (200) and (220) of Fe_{1-y}O were fitted with a pseudo-voigt function in order to obtain the full width at half maximum (FWHM), then the crystallite size (D_{XRD}) was obtained by using the Scherrer's equation which results in the 20 nm-25 nm range. Instead the XRD profile of sample $z = 0.06$ clearly evidences the presence of ferrite structure together with the Fe_{1-y}O phase, in almost similar amounts when considering the peak areas. From the Scherrer's equation the D_{XRD} were calculated for these phases, obtaining 7 nm for the reflection (311) of ferrite phase and 15 nm for the reflection (111) of Fe_{1-y}O phase. Moreover, for sample $z = 0.00$, both phases are also observed, with a relation of Ferrite/Fe_{1-y}O $\approx 2/1$ obtained from

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3 the respective peak's area. The values of $D_{\text{XRD}} \approx 12$ nm and 21 nm were calculated from the
4 reflection (200) of Fe_{1-y}O phase and from the reflections (311) and (400) of the ferrite,
5 respectively. The XRD analysis of the three samples shows that the proportion of Wüstite
6 phase increases as the Zn concentration in the particles increases: from a predominant
7 ferrite phase for sample $z = 0.00$ to the widely dominant Wüstite phase in sample $z = 0.10$,
8 in concordance with the increment in the corresponding crystallite size D_{XRD} of this phase.
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10 Figure 2 shows representative TEM images at lower magnification of nanoparticles in
11 samples $z = 0.10$, 0.06 and 0.00. The diameter histogram for each sample was build up by
12 measuring the diameter of more than 300 nanoparticles, also presented in figure 2. The
13 histograms were fitted with lognormal distribution and the mean diameters are 30 nm, 26
14 nm and 26 nm for samples $z = 0.10$, 0.06 and 0.00, respectively, with a dispersion of $\sigma \sim 0.2$
15 nm. A core/shell structure is observed in TEM images even at low magnification. From
16 TEM images at higher magnification, the shell thickness were measured and the relative
17 volume of each component were calculated, obtaining the volume proportion for the shell
18 ($\% V_{\text{shell}}$) of 55% for sample $z = 0.10$, 65% for $z = 0.06$ and 70% for $z = 0.00$, as given in
19 Table 1. Figure 3a shows a representative HRTEM image of nanoparticle in $z = 0.10$,
20 where the core/shell is clearly evidenced. The Fast Fourier Transform (FFT) of different
21 zones in the same particle gives the reciprocal d -spacing of each morphological component
22 (core and shell). The FFT (upper right panel) of shell zone (light Blue Square in HRTEM)
23 shows exclusively a d -spacing corresponding to the planes (220) of ferrite, which is marked
24 as a red solid line circle; any contribution of Wüstite is not expected in this d -spacing
25 range. Other interplanar distances appear for a large zone inside the core (orange square in
26 HRTEM): the reciprocal distances observed in this FFT image between the planes (220)
27 and (400) of ferrite corresponds to a d -spacing referent to the plane (200) of Wüstite, which
28 is indicated in the figure as the green dashed-line. Similar characteristics are observed for
29 the HRTEM of $z = 0.06$ (figure 3b). The FFT of shell zone also shows spots corresponding
30 to the (220) planes of ferrite (marked in figure as a red solid line), while the FFT of a large
31 area taking core and shell of the particle present other distances: the red circles for (311)
32 and (222) of ferrite and green dashed-line circle for (111) of Wüstite phase. Finally, similar
33 results are also obtained from the HRTEM of $z = 0.00$ (figure 3c). The FFT in shell zone
34 shows spots corresponding to (220) planes of ferrites (red on line) and the FFT of the core
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3 shows family of planes: (111), (220), (311) and (222) (not marked in the figure) for ferrites
4 and the reciprocal interplanar distance for (111) of Wüstite. Our result shows that sample z
5 = 0.10 has a bigger core volume than the other samples, presenting a systematic increase of
6 % V_{shell} with the Zn amount. Nevertheless, its increment in the core volume with the amount
7 of Zn not necessarily can be an exclusive consequence of the amount of Zn in the sample or
8 in the synthesis solution. Others small changes in the synthesis conditions, including the
9 redox potential could be related to this change.

10 Figure 4 shows a representative STEM-HAADF image of the nanoparticles (left panel)
11 corresponding to sample with $z = 0.10$. The linescan used for STEM SI is indicated. Two
12 spectra corresponding to the shell and the center of the core were extracted at the positions
13 indicated in the image and shown in the right panel. In the spectra the Fe-M edges are
14 visible at about 60 eV. The Zn-M edges can be observed as a weak bump at about 84 eV.
15 For the extraction of these edges an AE^{-r} power law background was subtracted. The width
16 of the energy window for typical background fit for Zn edge and the estimated background
17 are indicated in the figure 4. The background subtracted Zn-M edges are shown. The Fe/Zn
18 ratio obtained at both positions was the same for the range 8-9. This evaluation indicates
19 that Zn is present over the whole volume of the nanoparticle and that the doping
20 concentration is approximately the same in the core and in the shell.
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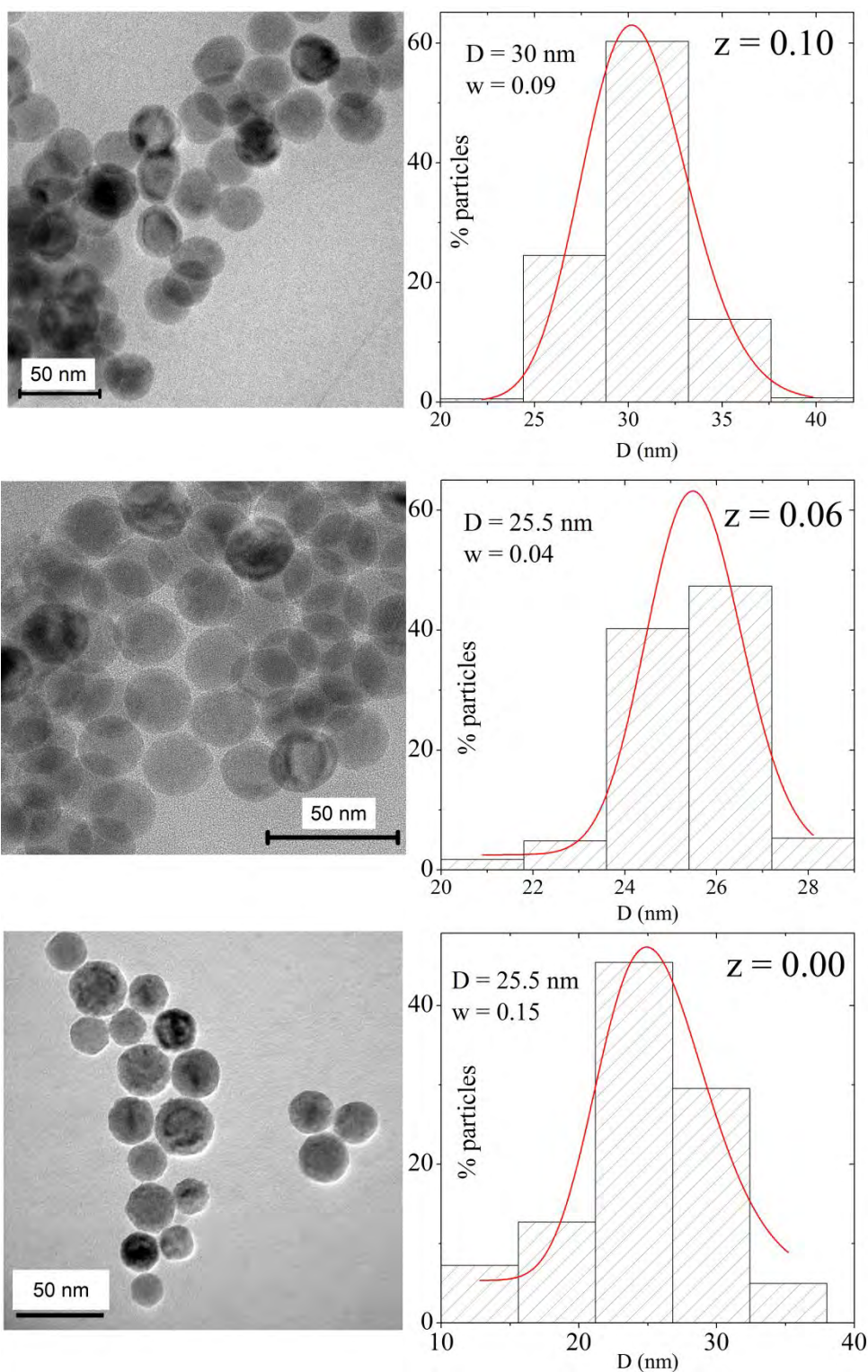


Figure 2 – Representative TEM image for each sample with the corresponding histogram of diameters obtained with measuring more than 300 particles. Each histogram is fitted with a lognormal distribution in order to obtain the mean diameter and the dispersion.

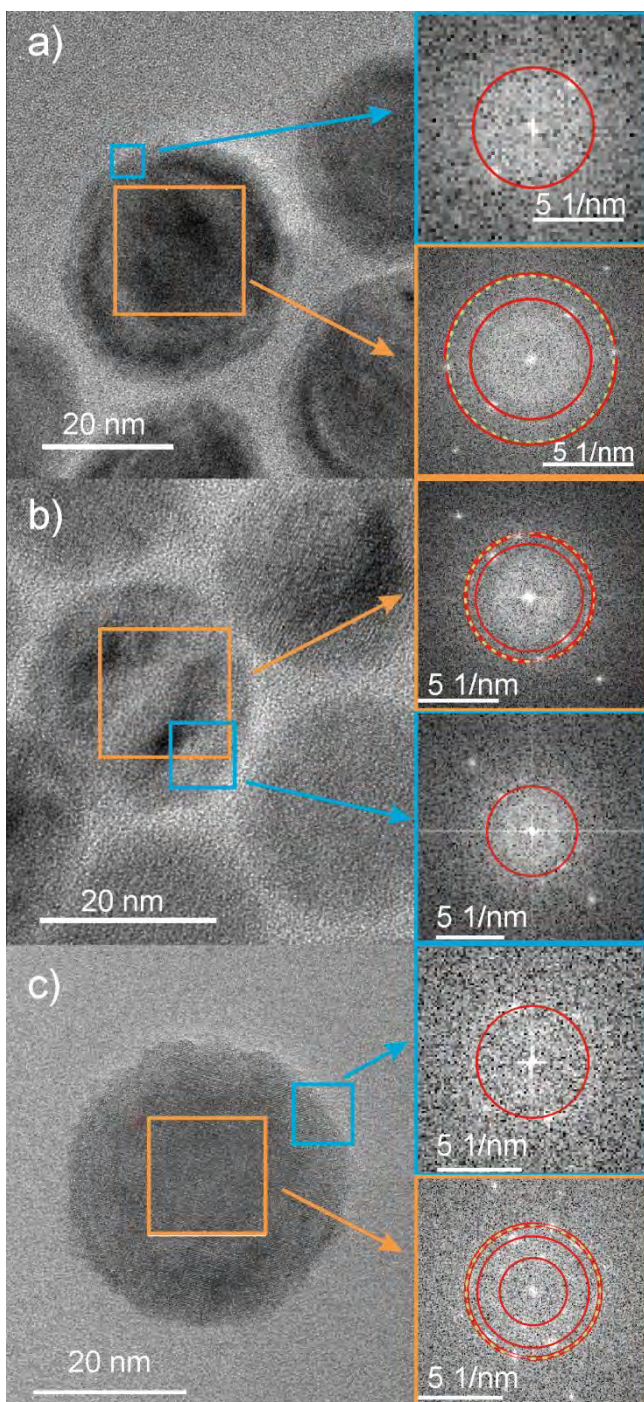


Figure 3 – Left panels show the HRTEM image of a representative nanoparticle of the sample with $z = 0.10$ (a), $z = 0.06$ (b) and $z = 0.00$ (c); right panels show the FFT from the zone marked with a square in the HRTEM image of the core (orange) and shell (blue), as indicated by the arrows. The FFT patterns of shell were indexed with the cubic spinel interplanar distances (solid red line), while the patterns of the core are indexed with cubic spinel and Wüstite (orange dashed line) interplanar distances.

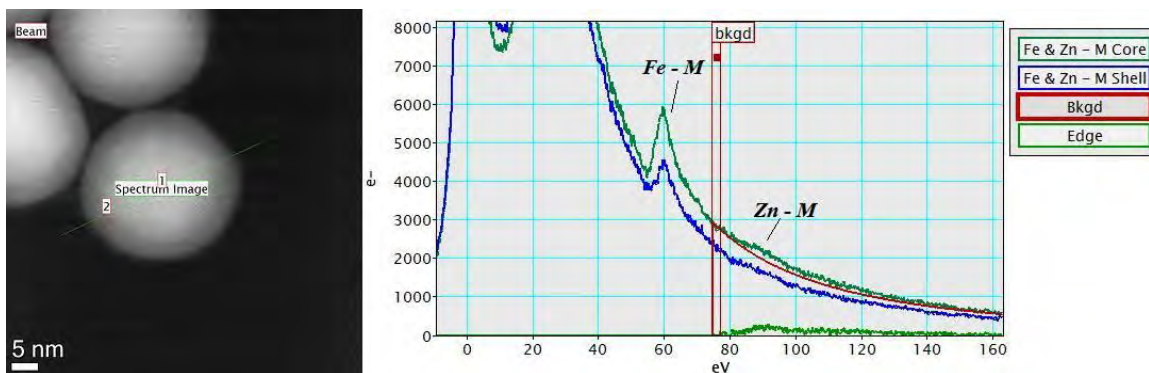


Figure 4 - STEM-HAADF image of nanoparticles of sample $z = 0.10$ (left) and low-loss EEL spectra (right) at different positions of the nanoparticle (indicated by 1 and 2 in the figure). The Fe-M and Zn-M edges are indicated. Background subtracted is indicated (red line). The background subtracted Zn-M edge is also shown.

Considering the morphological and compositional analyzes, we can obtain information about the growth mechanisms of the core-shell nanostructures. These results are consistent with the formation of the Fe_{1-y}O , at the initial stage of the synthesis followed by the formation of the ferrite phase as a consequence of the surface oxidation, which are in agreement with the PIXE analysis. This growth mechanism leads to the formation of well-defined interfaces, when they are compared with systems where the shell is growth in a second reaction^{22,40}, with the expectation of an improved magnetic exchange interaction between the phases. Other interesting feature is the reduction in the Zn incorporation in the formed material related to the nominal composition, which is probably related with the well-known low solubility of the Zn in the Wüstite structure mentioned before³².

Magnetic Properties

$M(T)$ curves of all samples measured in ZFC and FC protocols with $H = 50$ Oe are presented in figure 5. The characteristic FC magnetization curves of nanoparticles of FiM material below the blocking temperature is a slightly increasing curve with a not so pronounced slope⁴¹. For our samples, a marked decrease in the magnetization in $175 \text{ K} < T < 210 \text{ K}$ is observed for both ZFC and FC curves for all samples. This decrease is associated to the antiferromagnetic (AFM) ordering transition at the Néel temperature, T_N , of the Wüstite phase (T_N values of all samples are presented in Table 1). For these range of

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3 composition the AFM transition temperature does not show a significantly dependence with
4 the Zn concentration. It is noteworthy that the decrease in the magnetization at T_N is larger
5 with increasing the Zn content, in concordance with the increasing proportion of the AFM
6 Wüstite phase as revealed by the XRD and TEM analyzes. Notice that for sample $z = 0.10$,
7 where the Wüstite is the main phase, a sharper AFM transition at T_N is observed.
8 Interestingly, irreversibility between ZFC and FC curves is observed below T_N . This
9 irreversibility is ascribed to the magnetic behavior of the shell ferrimagnetic phase which is
10 strongly magnetic coupled to the AFM core below T_N . Therefore, as expect for strongly
11 FiM-soft/AFM-hard exchange coupled system, the blocking temperature shifts to the Néel
12 temperature of the hard magnetic component.^{21,23} Instead, for samples with $z = 0.06$ and z
13 $= 0.00$, the irreversibility temperature is above T_N , indicating that there is a distribution of
14 nanoparticles with blocking temperatures above T_N . This result could be the consequence
15 of the larger proportion of FiM spinel shell phase as determined by TEM and XRD
16 measurements and also to the lower concentration of Zn in the shell which lead to an
17 increasing of the shell magnetic anisotropy. As a consequence the AFM transition in these
18 sample is smoothed in comparison to that of sample $z = 0.10$.

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31 The saturation magnetization of the samples was determined from the $M(H)$ curves
32 measured at $T = 300$ K. As observed in Table 1, M_S decreases from 60 emu/g to 3 emu/g
33 when the concentration of Zn increases. The huge reduction in M_S for sample $z = 0.10$
34 comes mainly from the larger proportion of Wüstite phase, which is paramagnetic at room
35 temperature. Besides, the larger substitution of Fe by Zn, plus the surface disorder driven
36 by the thinner shell in this sample, could induce a strong spin-canting in the FiM shell,
37 which reduces M_S and the effective magnetic anisotropy K ⁴².

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43 The zinc ferrite normal spinel shows AFM ordering below ~ 10 K; instead, in our partially
44 substituted ferrite, FiM ordering is expected^{43,44}. Although the Zn^{2+} has tetrahedral site
45 preference in the spinel structure, at nanometric scale some cationic inversion degree is
46 observed. It is also well-known that the presence of Zn leads to spin-canting as resulting of
47 its diamagnetic nature and the consequent lack of exchange coupling for the Fe^{3+} and Fe^{2+}
48 in the neighborhood for Zn concentrations higher than 0.16⁴⁵. In addition, spin canting is
49 observed for systems as only consequence of the nanometric size^{43,44}, which should be
50 increased in our system as consequence of the thickness of the ferrite phase shell estimated
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of 5 nm for sample $z = 0.10$. Finally, a strong magnetic interaction between the FiM shell and the AFM core leads to the increment of uncompensated spins from the core and/or shell as well as to the increment of the magnetic canting in the shell phase^{21,46}. Thus, the low magnetization values for high field observed for sample $z = 0.10$ probably reflects the nanometric core-shell magnetically coupled nature of the nanoparticles together with the partial substitution of Fe ion by Zn ions in the shell.

The hysteresis curves at $T = 5$ K and $T = 100$ K were measured by using two different protocols (Figure 5 right panel): after cooling from 300 K without (ZFC) and with an applied magnetic field of 50 kOe (FC). Remarkable, in the FC mode all samples shows large exchange bias fields (H_{bias}). The presence of H_{bias} indicates a strong magnetic coupling between the AFM and the FiM phases, supporting the same conclusion observed in $M(T)$ curves where an increases of the blocking temperature of the shell to the AFM ordering transition temperature of the core was observed. Sample $z = 0.10$ presents the larger exchange bias field where at 5 K results $H_{\text{bias}}^{5\text{K}} \sim 3420$ Oe. However, the FC cycle measured up to 5 T corresponds to a minor cycle; and larger $H_{\text{bias}}^{5\text{K}}$ is expected for higher applied fields. Sample $z = 0.06$ presents a totally saturated magnetization loop at 5 K, besides the M_s as well the H_C measured from the FC cycle are larger than the ones obtained from the ZFC measurement. This results is a consequence of the magnetic coupling at the AFM/FiM interface, where the FC protocol preferentially orient the AFM surface spins, and as a consequence enhance the coercivity and also the saturation magnetization. As expected the sample with $z = 0.00$ presents the lower $H_{\text{bias}}^{5\text{K}}$ (86 Oe), and the largest value of M_s , since it has the smallest AFM phase proportion.

Table1: Nanoparticles characteristic obtained from TEM, $M(T)$ and $M(H)$ measurements: $\%V_{\text{shell}}$ is the volume percentage of shell as calculated from the shell thickness and core diameter from TEM images, M_s is the saturation magnetization at room temperature, T_N is the Néel temperature, $H_{\text{bias}}^{5\text{K}}$ and $H_{\text{bias}}^{100\text{K}}$ are the exchange bias fields and $H_C^{5\text{K}}$ and $H_C^{100\text{K}}$ are the coercive fields, obtained from $M_{\text{FC}}(H)$ curves at 5 K and 100 K, respectively.

Sample	$\%V_{\text{shell}}$ (%)	M_s (emu/g)	T_N (K)	$H_{\text{bias}}^{5\text{K}}$ (Oe)	$H_C^{5\text{K}}$ (Oe)	$H_{\text{bias}}^{100\text{K}}$ (Oe)	$H_C^{100\text{K}}$ (Oe)
$z=0.00$	70	59.6	196	86	30	30	454

$z=0.06$ (2)	65	16.8	180	2270	3580	435	673
$z=0.10$ (2)	55	3.2	177	3420	6150	429	276

Although in ideal AFM/FiM exchange coupled system it is predicted that the exchange bias properties should be present up to T_N ; in core/shell nanoparticles, as a consequence of finite size effect and surface imperfections, it is usually observed that the exchange bias field vanish at much lower temperature.²² In the present system the measurements performed at 100 K also present exchange bias field ($H_{\text{bias}}^{100\text{K}}$), indicating a good magnetic coupling through the interface; in fact, the exchange bias was observed up to 150 K for sample $z = 0.10$, close to the ordering temperature of the AFM phase. Figure 6 compares the variation of $H_{\text{bias}}^{5\text{K}}$ and $H_{\text{bias}}^{100\text{K}}$ with the Zn concentration. As expected, $H_{\text{bias}}^{5\text{K}}$ and also $H_C^{5\text{K}}$ increases with the Zn amount because the diameter of the AFM core grows and consequently the magnetic anisotropy of the AFM phase increases relatively to the other concentration. Instead, at higher temperature $H_{\text{bias}}^{100\text{K}}$ and $H_C^{100\text{K}}$ do not follow a monotonous dependence with the concentration. We ascribed this behavior to the thermal fluctuations, which are more important when the anisotropy is reduced as a consequence of the Zn incorporation in the structure. In summary, the overall magnetic measurements indicates that the nanoparticles cores are composed by the AFM phase with Wüstite structure while the shell presents a FiM ordering, both phase are strongly coupled below T_N even at temperature higher than 100 K and finally the properties of the ferrite phase in the shell differ from those of magnetite or maghemita as a consequence of the coupling and the surface disorder.

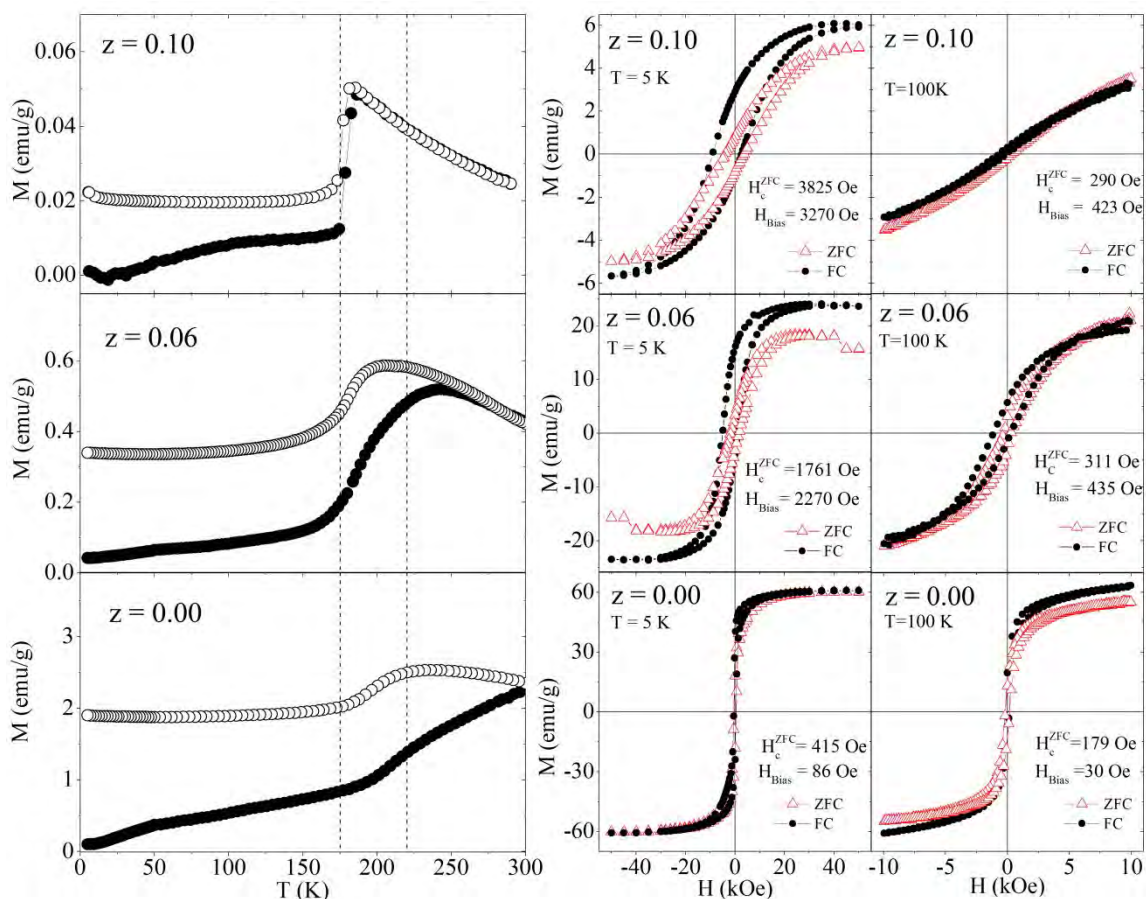


Figure 5 – Left panel: $M(T)$ data of all samples measured in both ZFC (solid symbols) and FC (open symbols) protocols with $H_{FC} = 50$ Oe. Right panel: ZFC (open symbols) and FC (solid symbols) hysteresis loops measured at 5 K (left panels) and 100 K (right panels). The values of coercive fields after ZFC (H_c^{ZFC}) and bias fields after FC defined as $H_{Bias} = (H_c^+ - H_c^-)/2$ are also displayed.

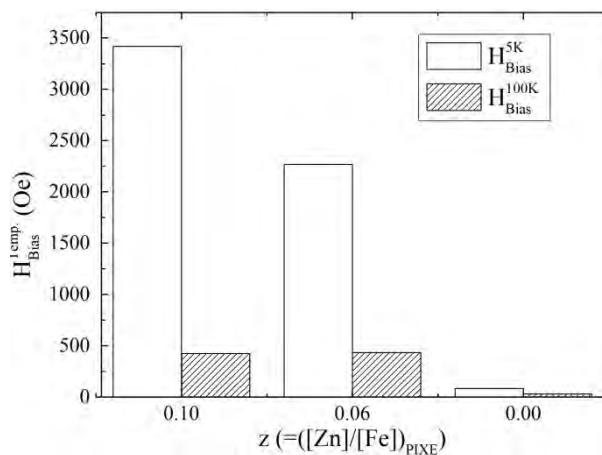


Figure 6 - Exchange bias field at 5 K ($H_{\text{bias}}^{5\text{K}}$) and 100 K ($H_{\text{bias}}^{100\text{K}}$) as a function of the Zn concentration of the sample obtained from PIXE analysis.

In-Field Mössbauer Spectroscopy

In order to get local information about the magnetic arrangement with the presence of the Zn in the bi-magnetic core-shell structure, in-field Mössbauer Spectroscopy was performed for sample $z = 0.10$. The spectra are given in figure 7 and they were collected with different applied fields. The spectra were fitted with two crystalline sextets, labeled sextets 1 and 2, a singlet probably related to Fe^{3+} ions in the tetrahedral interstitial sites of Wüstite structure^{18,47,48} and sextet with broad lines and lower effective field ($B_{\text{eff}} = \overrightarrow{B_{\text{hyp}}} + \overrightarrow{B_{\text{app}}}$, and $\overrightarrow{B_{\text{hyp}}}$ is the hyperfine field), labeled sextet 3. The linewidth was hold between the widths of the 6 lines of sextet: $\Gamma_{3,4} = W + 0.02526 * dW$, $\Gamma_{2,5} = W + 0.09344 * dW$, $\Gamma_{1,6} = W + 0.16147 * dW$. This can be shown to result from a relatively narrow effective field distribution, such as can often be seen in ferrites. All hyperfine parameters obtained from the fitting procedure are given in Table 2. For applied field from 0 to 120 kOe, the hyperfine field of the sextets 1 and 2, those with larger B_{eff} increase and decrease, respectively, with the applied field as expected for a strong antiferromagnetic alignment, specifically for a FiM ordering as the subspectra areas are different with a relative relation of 1:1.6. These sextets could be easily referred to the two AFM aligned magnetic sublattices of the ferrite^{49,50} consistent with the absence of the Verwey transition and in an opposition to the five sextets of the magnetite below the Verwey transition⁴⁷. Nevertheless, as the Zn ions are known to occupy preferentially the A site in the ferrite structure, we expect an increase in this relation from 1.67 or 2 (maghemite or magnetite), not the observed decrease. In another way, if we consider a strong coupling among both sites, with the AFM ordering in the phase with structure of Wüstite aligned with the A and B magnetic lattice of the ferrite, and considering presence of 10 % of Zn in the A site of the ferrite structure and the amounts of each phase as obtained from the volume estimated from HRTEM images, then a relation of areas about 1.60 between both sites is expected, very similar to that obtained from the fitting procedure. In this way, our In-Field Mössbauer data probably indicates a strongly correlation among the magnetic phases, with the sublattices of ferrite and those of AFM phase being strongly coupled and mostly aligned parallel and

antiparallel to the applied field direction. Another interesting feature concerning the dependence of the B_{eff} of these sites with the applied field is observed, with a non-linear dependence. Specifically, for $H_{\text{app}} < 120$ kOe, both sextets present a linear dependence but with angular coefficient of 0.8, lower than the expected for a perfect alignment (about 1), with B_{eff} at $H_{\text{app}} = 120$ kOe being the expected for a perfect alignment. Therefore, this effect probably indicates a canting of these lattices probably related to a spin-glass-like in the AFM phase or in the interface layer of the core/shell nanostructure⁴⁴. Moreover, the angles between the Fe magnetic moment and the external magnetic field (θ , displayed in Table 2) can be calculated from the D12 parameter⁵⁰, and these angles corroborate with this argument. The third sextet with lower area and B_{eff} is probably also related to a spin-glass-like magnetic structure in the AFM phase or in the interface between the magnetic phases.

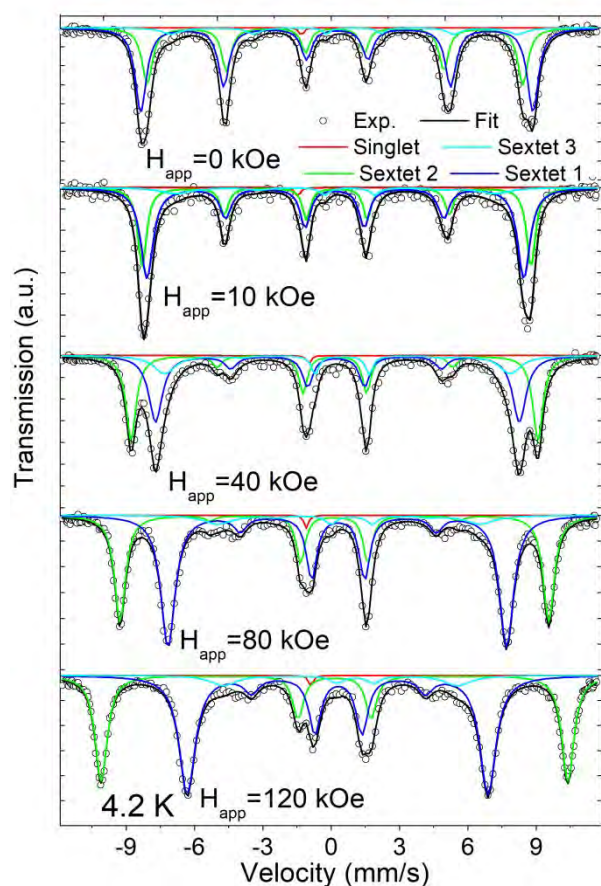


Figure 7 – In-Field Mossbauer Spectra of sample $z=0.10$ collected at 4.2 K with applied field up to 120 kOe. All the spectra were fitted with three sextets (Sextet 1 – blue, Sextet 2 – green and Sextet 3 – light blue) and a singlet (red).

Table 2 – Hyperfine Parameters obtained from the fitting procedure of the In-field Mössbauer spectra at 4.2 K: H_{app} is the applied field, B_{eff} is the effective field, IS is the Isomer Shift, QS is the quadrupolar splitting, W is the line width, dW is a parameter related with the width of each line of the sextets, D21 is the relation between the intensities of the lines 1 and 2, and the Area is the relative absorption area of each component.

H_{app} (kOe)	Subspectra	B_{eff} (T)	IS (mm/s)	QS (mm/s)	W (mm/s)	dW	D21	Θ degree	Area (%)
120	Sextet 1	63.4	0.28	-0.01	0.55	0.59	0	0	35.4
	Sextet 2	40.9	0.43	-0.04	0.60	1.60	0.12	156	56.7
	Singlet	-	-0.89	-	0.23	-	-	-	0.5
	Sextet 3	31.8	0.87	-0.64	0.59	5.00	0	0	7.4
80	Sextet 1	58.3	0.27	0.01	0.44	0.93	0.04	14	35.1
	Sextet 2	46.0	0.42	-0.04	0.52	1.39	0.10	158	56.2
	Singlet	-	-0.97	-	0.23	-	-	-	0.6
	Sextet 3	35.7	0.97	-0.11	0.38	7.07	0.12	24	8.1
40	Sextet 1	55.3	0.29	-0.02	0.41	0.81	0.12	24	30.5
	Sextet 2	49.4	0.38	0.06	0.53	1.19	0.17	152	48.9
	Singlet	-	-0.85	-	0.23	-	-	-	0.6
	Sextet 3	47.0	0.51	-0.21	0.25	5.56	0.07	18	20.0
10	Sextet 1	52.8	0.37	0.00	0.40	0.37	0.33	39	35.1
	Sextet 2	51.2	0.30	0.03	0.52	0.84	0.30	143	56.2
	Singlet	-	-1.33	-	0.23	-	-	-	0.6
	Sextet 3	43.4	0.63	-0.65	0.25	7.90	0.21	31	8.1
0	Sextet 1	51.1	0.30	0.00	0.42	0.77	0.69	-	35.5
	Sextet 2	53.3	0.39	-0.01	0.56	0.15	0.69	-	56.7
	Singlet	-	0.19	-	0.7	-	-	-	0.5
	Sextet 3	47.1	0.92	-0.42	0.25	4.61	0.69	-	7.3

CONCLUSIONS

The synthesis of Zn-incorporated iron-oxide core-shell bimagnetic nanoparticles presented in this work represents a novel approach to the control of their magnetic properties by a chemically-tunable, magnetically coupled ferrimagnetic shell. The formation of an AFM Fe_{1-y}O core strongly coupled to the FiM ferrite shell can be simply obtained from a standard one-pot thermal decomposition synthesis from $\text{Fe}(\text{acac})_3$ and $\text{Zn}(\text{acac})_2$ precursors. Compositional analysis by PIXE indicates that the Zn incorporation in the particle is lower than the nominal value of the synthesis, at the same time the compositional analysis by EELS indicates that the Zn is incorporated almost homogeneously in the whole particle, including the core and the shell. This result is interesting concerning the formation of the nanostructure, indicating the low solubility of the Zn ion in the AFM Wüstite phase (core), followed by the surface oxidation leading to the formation of the FiM ferrite phase at the shell, keeping the Zn/Fe relation along the whole nanoparticle. The incorporation of the Zn allows modified the overall magnetic properties of the MNPs. The effective anisotropy of the shell decreases with increasing the Zn amount, at the same time the AFM ordering of the core leads to an increment in the anisotropy of the whole system below T_N . No significant change in the T_N of the AFM phase was observed with the limited incorporation of the Zn. Finally, a strong magnetic coupling between both phases is evidenced by the presence of the exchange bias field, with a canting of the spins observed with increasing the Zn amount in the system.

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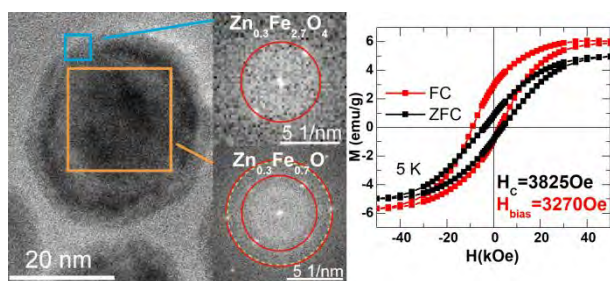
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TOC Graphic



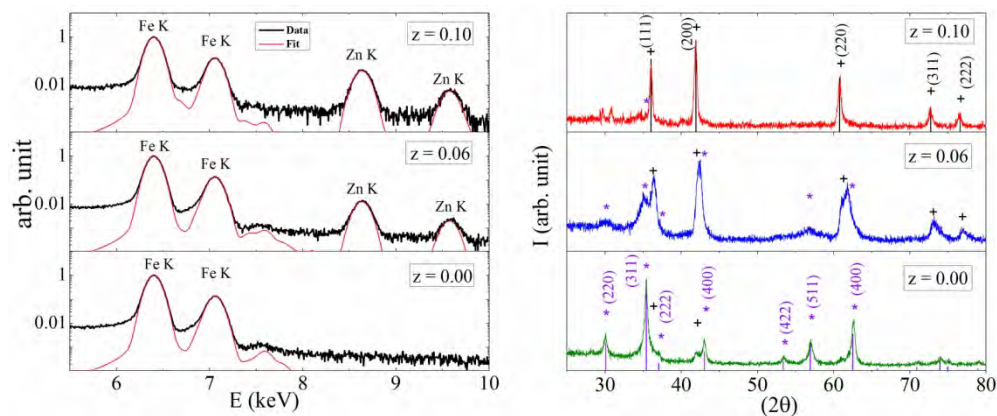


Figure 1 – Left Panel: PIXE spectra of all samples showing the pertinent energy range of Fe and Zn K-peaks. Black lines correspond to experimental data, and red lines are the respective fitting to obtain the $z = [\text{Zn}]/[\text{Fe}]$ ratio for each system. Right panel: XRD profiles of three samples. All peaks are indexed with the Fe_{1-y}O (+) and ferrite (*) phases where the bars correspond to the diffraction peaks according to the JCPDS card number: for Wüstite 00-006-0615 (upper panel) and magnetite 00-019-0629 (bottom panel).

159x66mm (300 x 300 DPI)

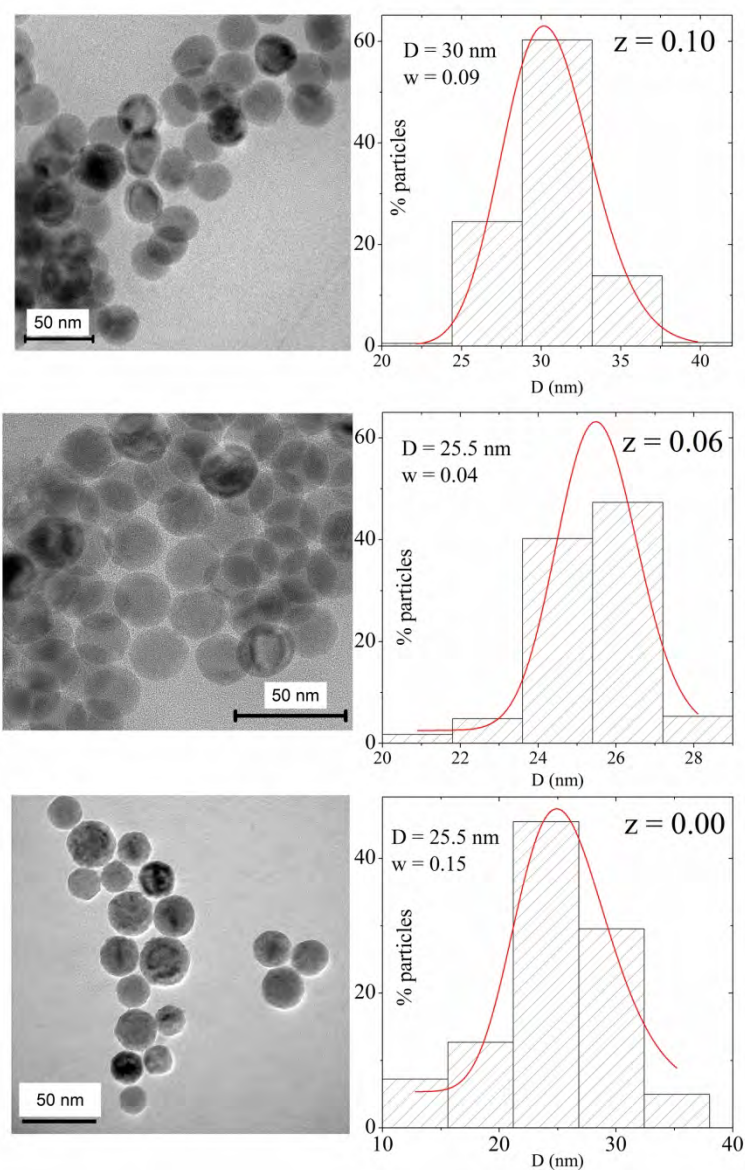


Figure 2 – Representative TEM image for each sample with the corresponding histogram of diameters obtained with measuring more than 300 particles. Each histogram is fitted with a lognormal distribution in order to obtain the mean diameter and the dispersion.

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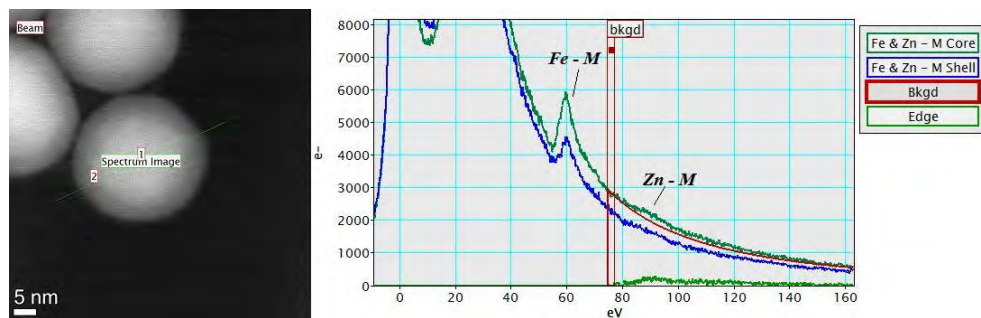


Figure 4 - STEM-HAADF image of nanoparticles of sample $z = 0.10$ (left) and low-loss EEL spectra (right) at different positions of the nanoparticle (indicated by 1 and 2 in the figure). The Fe-M and Zn-M edges are indicated. Background subtracted is indicated (red line). The background subtracted Zn-M edge is also shown.

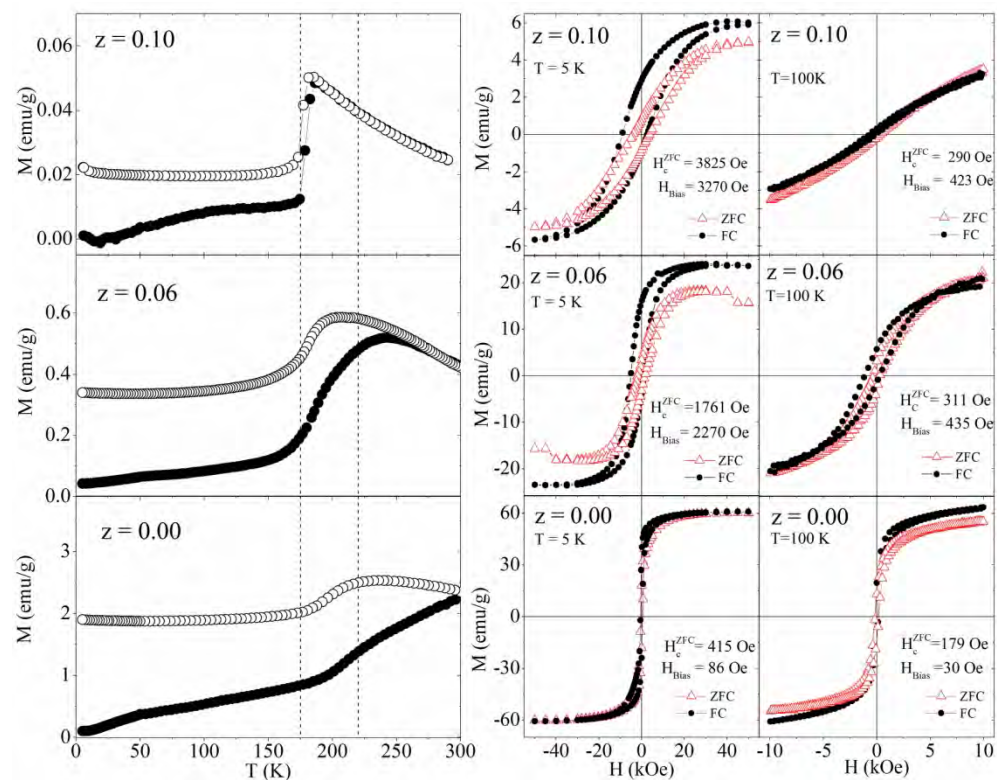


Figure 5 – Left panel: $M(T)$ data of all samples measured in both ZFC (solid symbols) and FC (open symbols) protocols with $H_{FC}=50$ Oe. Right panel: ZFC (open symbols) and FC (solid symbols) hysteresis loops measured at 5 K (left panels) and 100 K (right panels). The values of coercive fields after ZFC (H_c^{ZFC}) and bias fields after FC defined as $H_{bias}=(H_c^+ - H_c^-)/2$ are also displayed.

159x124mm (300 x 300 DPI)

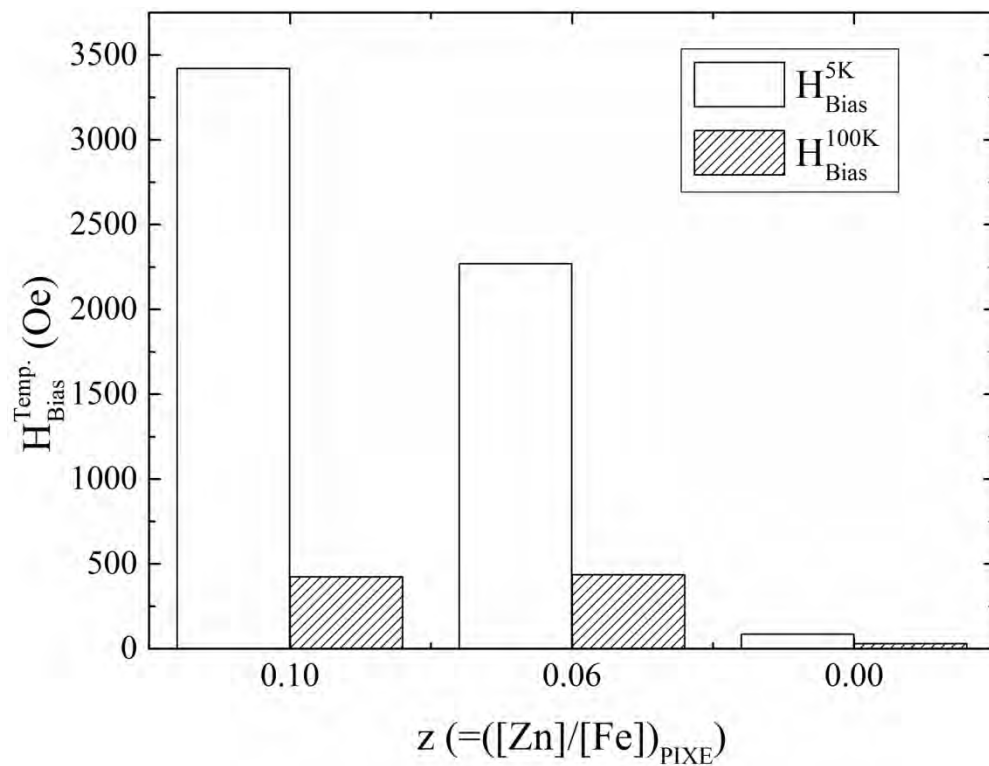
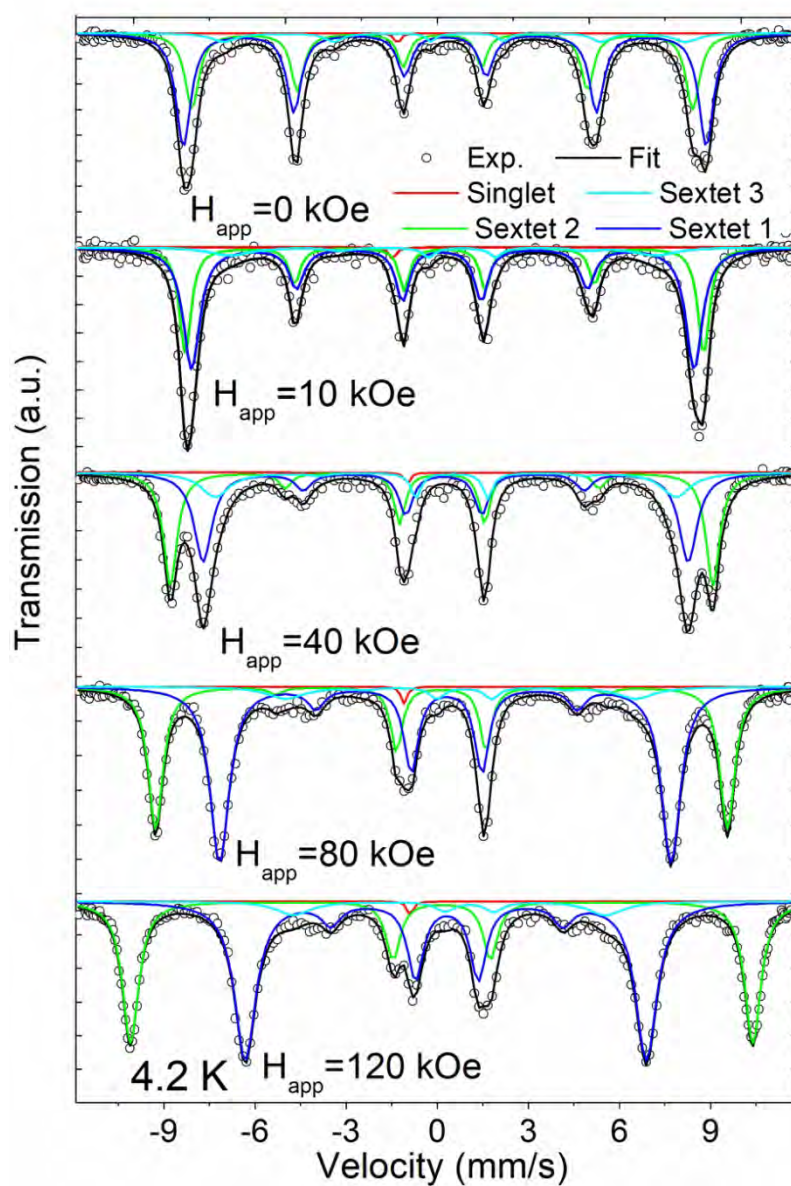


Figure 6 - Exchange bias field at 5 K ($H_{\text{Bias}}^{5\text{K}}$) and 100 K ($H_{\text{Bias}}^{100\text{K}}$) as a function of the Zn concentration of the sample obtained from PIXE analysis.

79x61mm (600 x 600 DPI)



45 Figure 7 – In-Field Mossbauer Spectra of sample $z=0.10$ collected at 4.2 K with applied field up to 120 kOe.
46 All the spectra were fitted with three sextets (Sextet 1 – blue, Sextet 2 - green and Sextet 3 – light blue)
47 and a singlet (red).

48 79x118mm (300 x 300 DPI)